Exploring Quantum-Classical Boundary

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Education

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Professional Employment

1992 Research Associate. Tohoku University Associate Professor, Tohoku University Professor, Institute for Molecular Science

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2009 Visiting Professor, The University of Tokyo (-2011)

2012 Visiting Professor, University of Heidelberg

2014 Visiting Professor, University of Strasbourg (-2016)

Awards

1998 Award by Research Foundation for Opto-Science and Technology

2007 JSPS Prize

2007 Japan Academy Medal

Fellow of the American Physical Society

2012 Humboldt Research Award

2017 Hiroshi Takuma Memorial Prize of Matsuo Foundation

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It is observed in a double-slit experiment by Tonomura and coworkers that single electrons recorded as dots on a detector screen build up to show an interference pattern, which is delocalized over the screen. 1) This observation indicates that a delocalized wave function of an isolated electron interacts with the screen, which is composed of many nuclei and electrons interacting with each other, and becomes localized in space. This change, referred to as "collapse" in quantum theory, is often accepted as a discontinuous change, but a basic question arises: When and how the delocalized wave function becomes localized? Our objective is uncovering this mystery by observing the spatiotemporal evolution of a wave function delocalized over many particles interacting with each other. Having this objective in mind, we have developed coherent control with precisions on the picometer spatial and attosecond temporal scales. Now we apply this ultrafast and ultrahigh-precision coherent control to delocalized wave functions of macroscopic many-particle systems such as an array of ultracold rubidium (Rb) Rydberg atoms in an optical lattice, as depicted schematically in Figure 1, envisaging the quantum-classical boundary connected smoothly.

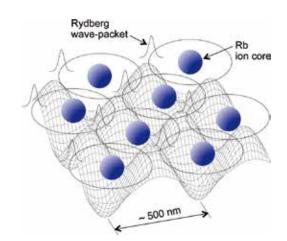


Figure 1. Schematic of the many-body system of ultracold Rydberg atoms.2

Selected Publications

- H. Katsuki et al., "Visualizing Picometric Quantum Ripples of Ultrafast Wave-Packet Interference," Science 311, 1589-1592
- · H. Katsuki et al., "Actively Tailored Spatiotemporal Images of Quantum Interference on the Picometer and Femtosecond Scales," Phys. Rev. Lett. 102, 103602 (2009).
- · K. Hosaka et al., "Ultrafast Fourier Transform with a Femtosecond-Laser-Driven Molecule," Phys. Rev. Lett. 104, 180501 (2010).
- · H. Goto et al., "Strong-Laser-Induced Quantum Interference," Nat.

- Phys. 7, 383-385 (2011).
- · H. Katsuki et al., "All-Optical Control and Visualization of Ultrafast Two-Dimensional Atomic Motions in a Single Crystal of Bismuth," Nat. Commun. 4, 2801 (2013).
- · N. Takei et al., "Direct Observation of Ultrafast Many-Body Electron Dynamics in an Ultracold Rydberg Gas," Nat. Commun. 7, 13449 (2016).
- C. Liu et al., "Attosecond Control of Restoration of Electronic Structure Symmetry," Phys. Rev. Lett. 121, 173201 (2018).

1. Development of an "Ultrafast Quantum Simulator" by Optical Control with Precisions on the Attosecond Temporal and Submicron Spatial Scales^{3–5)}

Quantum many-body problems are at the heart of a variety of physical functionalities including superconductivity and magnetism in solid materials. It is extremely hard, however, to solve such quantum many-body problems. In solving the Hubbard model with 1000 particles, for example, the diagonalization would take 10 to the power of 573 years even with the world's fastest supercomputers. In this project, we develop a novel quantum simulator that can simulate quantum many-body dynamics for more than 1000 particles within one nanosecond, combining our two unique experimental resources: "coherent control with attosecond precision" and "a strongly-correlated ultracold Rydberg gas."

We have completed a standard hardware of this ultrafast quantum simulator composed of an array of ultracold Rb atoms trapped in an optical lattice and excited to Rydberg levels with a coherent picosecond laser pulse, as schematically illustrated in Figure 2.^{3,4}) The broad bandwidth of the ps laser pulse has allowed us to excite the atoms in the neighboring lattice sites to Rydberg levels simultaneously for the first time. We expect that the overlap between their Rydberg orbitals could give rise to an exotic metal-like phase, where the electron could delocalize over the lattice.

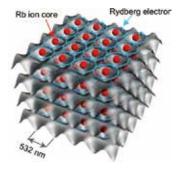


Figure 2. Schematic of the hardware of the ultrafast quantum simulator, 3,4)

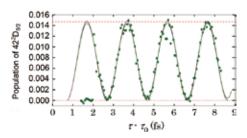


Figure 3. Time domain Ramsey interferometry of ultracold ⁸⁷Rb atoms with attosecond precision to be used as a readout interface of the ultrafast quantum simulator. Population of the $42^2D_{5/2}$ Rydberg state is plotted as a function of the delay τ between two laser pulses, where $\tau_0 \sim 50$ ps. Adopted from Ref. 5).

We have also completed a readout interface of our ultrafast quantum simulator, which is the time domain Ramsey interferometry of ultracold Rydberg atoms with attosecond precision, whose contrast is almost 100% as shown in Figure 3.⁵⁾ The phase and visibility of this Ramsey interferogram are highly sensitive to the nature and strength of many-body interactions among the Rydberg atoms.

2. Application of an "Ultrafast Quantum Simulator" to Quantum Computing³⁾

We are developing a cold-atom based quantum annealer with the hardware of the ultrafast quantum simulator.⁸⁾ The cold-atom annealer has advantages against the one with the superconducting qubits. Those advantages include scalability and efficiency. All to all connections among physical bits necessary for quantum annealing could also be easier with cold atoms than superconducting qubits.

So far we have developed arbitrary two dimensional optical trap arrays for cold atoms, which are necessary for quantum annealing, 8) in tight collaborations with Hamamatsu Photonics K. K.³⁾ Their examples are shown in Figure 4, where we used a convex lens for clear visualization of the trap arrays. We have recently replaced this convex lens by an objective lens to realize the world's smallest arbitrary trap array with its nearest neighbor distance now only 1.06 micron, which used to be typically ~4 micron in previous works.⁹⁾



Figure 4. Examples of arbitrary optical trap arrays generated with a low NA convex lens ($L = 8.3 \, \mu \text{m}$; wavelength = 633 nm; NA = 0.0446), corresponding to $L = 640 \, \text{nm}$ with NA = 0.75 at the wavelength 820 nm.³⁾

References

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- White Paper 2018 on Manufacturing Industries published by Ministry of Economy Trade and Industry, JAPAN.
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