## **Time-dependent Hartree-Fock study**

## on photoinduced dynamics in two-dimensional organic conductors

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Photoinduced melting of charge order (CO) in quasi-two-dimensional organic (ET=BEDT-TTF) and  $\theta$ -(ET)<sub>2</sub>RbZn(SCN)<sub>4</sub> is investigated conductors  $\alpha$  -(ET)<sub>2</sub>I<sub>3</sub> theoretically.  $\alpha$  -(ET)<sub>2</sub>I<sub>3</sub> and  $\theta$  -(ET)<sub>2</sub>RbZn(SCN)<sub>4</sub> are the typical compounds with a quarter filled  $\pi$ -band which are known to exhibit CO. The photoinduced melting of CO in these systems has been observed by using the femtosecond spectroscopy recently[1]. It indicates that a semimacroscopic metallic domain is generated for  $\alpha$  -(ET)<sub>2</sub>I<sub>3</sub>, whereas only local melting of CO is created for  $\theta$  -(ET)<sub>2</sub>RbZn(SCN)<sub>4</sub>. In particular, the dynamics in  $\alpha$  -(ET)<sub>2</sub>I<sub>3</sub> shows critical slowing down and strong dependence on the excitation intensity and temperature. This clear difference is considered to originate from the roles of lattice distortion in stabilizing COs. In fact, the transition is of first-order with a large structural distortion and large discontinuity in  $\theta$  -(ET)<sub>2</sub>RbZn(SCN)<sub>4</sub>, whereas the lattice distortion in  $\alpha$  -(ET)<sub>2</sub>I<sub>3</sub> is relatively small. These observations are consistent with the theoretical calculations with Hartree-Fock approximation[2,3]. Therefore, it is of great interest to investigate the photoinduced melting of CO in these systems and compare their dynamics.

In order to study the photoinduced dynamics in each compound, we solve the timedependent Schrödinger equation numerically within the Hartree-Fock approximation for an extended Peierls-Hubbard model. The obtained charge, spin and lattice dynamics are considered to reflect the different natures of charge ordered states in these systems. In particular, the melting of CO needs more energy for  $\theta$  -(ET)<sub>2</sub>RbZn(SCN)<sub>4</sub> than for  $\alpha$  -(ET)<sub>2</sub>I<sub>3</sub>, which is a consequence of large lattice distortion and the essential role of electronphonon coupling in stabilizing the CO in  $\theta$  -(ET)<sub>2</sub>RbZn(SCN)<sub>4</sub>.

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