## VIII-KK Theoretical Studies on Electronic Structure and Dynamics of Electronically Excited States

VIII-KK-1 Millimeter-Wave Spectroscopy of the Internal-Rotation Band of the X-HCN Complex (X = He, Ne, and Ar) and the Intermolecular Potential Energy Surface

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Millimeter-wave absorption spectroscopy combined with a pulsed-jet expansion technique was applied to the measurement of the internal-rotation band of the X-HCN (X = He, Ne, and Ar) in the frequency region of 95-125 GHz. The observed transition frequencies were analyzed including their hyperfine splitting to yield an intermolecular potential energy surface, as improved from the one given by a coupled-cluster single double (triple) ab initio calculation. Regarding the He complex, the surface obtained has a global minimum in the linear configuration (He···H–C–N) with a well depth of 30.2cm<sup>-1</sup>, and the saddle point located in the anti-linear configuration (H–C–N···He) is higher by 8.174 cm<sup>-1</sup> in energy than the global minimum. The distance  $R_{\rm m}$ between the He atom and the center of mass of HCN along the minimum energy path shows a strong angular dependence;  $R_{\rm m} = 4.169$  Å and 4.040 Å in the linear and anti-linear forms, respectively, while it is 3.528 Å in a T-shaped configuration. The energy level diagram is consistent with the millimeter-wave observation. And now, we are trying to measure the stretching motion located just below the dissociation limits.

#### VIII-KK-2 Theoretical Study of Vibrational Spectra of *p*-Tert-Butylcalix[4]crown-6-Ether Complexed with Ethyl Ammonium Cations

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#### [Chem. Phys. Lett. 374, 572-576 (2003)]

Theoretical infrared (IR) absorption spectra were calculated for *p-tert*-butylcalix[4] crown-6-ether (1) in the cone conformer and its ethyl ammonium complex. The IR spectra were obtained by restricted Hartree-Fock (RHF) calculations with the 6–31 G basis set. For the purpose of an absorption band assignment on the host molecule, guest molecule, and complexes thereof, we discussed a way to distinguish a specific molecule by comparing the calculated vibrational spectra. The theoretical result for the host molecule and its ethyl ammonium complex were preliminarily compared with the experimental result, and found that the calculated result agrees well with the features of the experimental spectra.

VIII-KK-3 He\*(2<sup>3</sup>S) Penning Ionization of H<sub>2</sub>S I. Theoretical Franck-Condon Factors for the H<sub>2</sub>S(X<sup>1</sup>A<sub>1</sub>, v' = 0)  $\rightarrow$  H<sub>2</sub>S<sup>+</sup> (X<sup>2</sup>B<sub>1</sub>, A<sup>2</sup>A<sub>1</sub>) Ionization and the H<sub>2</sub>S<sup>+</sup> (A-X) Transitions

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In order to elucidate the ionization dynamics, in particular the vibrational distribution, of  $H_2S^+(A)$ produced by the Penning ionization of H<sub>2</sub>S with He\*(2<sup>3</sup>S) atoms, the Franck-Condon factors (FCFs) were presented for the  $H_2S(\tilde{X}) \rightarrow H_2S^+(\tilde{X}, \tilde{A})$  ionization and the H<sub>2</sub>S<sup>+</sup>( $\tilde{A}-\tilde{X}$ ) transition, and Einstein's A coefficients were presented for the latter transition. The FCFs were obtained by quantum vibrational calculations using the global potential energy surfaces (PESs) of  $H_2S(\tilde{X}^1A_1)$  and  $H_2S^+(\tilde{X}^2B_1, \tilde{A}^2A_1, \tilde{B}^2B_2)$  electronic states. The global PESs were determined by the multireference configuration interaction calculations with the Davidson correction (MRCI+Q) and the interpolant moving least squares (IMLS) method combined with the Shepard interpolation. The obtained FCFs exhibit that the  $H_2S^+(\tilde{X})$  state primarily populates the vibrational ground state since its equilibrium geometry is almost equal to that of  $H_2S(\tilde{X})$ , while the bending mode  $(v_2)$  is strongly enhanced for the  $H_2S^+(\tilde{A})$  state; the maximum in the population is around  $v_2 = 6-7$ . In the same manner, the bending progressions are expected to consist of the great part of the H<sub>2</sub>S<sup>+</sup>( $\tilde{A}$ - $\tilde{X}$ ) emission. A detailed comparison with the experimental study for this system is reported in the accompanying paper, Paper II.

#### VIII-KK-4 He\*( $2^3$ S) Penning Ionization of H<sub>2</sub>S II. Formation of SH<sup>+</sup>( $A^3\Pi$ ) and H<sub>2</sub>S<sup>+</sup> ( $A^2A_1$ ) Ions

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Emissions in the 200–750 nm region produced by the collision of He(2<sup>3</sup>S) with H<sub>2</sub>S were studied under single-collision conditions. The hydrogen Balmer lines and the SH<sup>+</sup>( $A^{3}\Pi$ – $X^{3}\Sigma^{-}$ ) and H<sub>2</sub>S<sup>+</sup>( $\tilde{A}^{2}A_{1}-\tilde{X}^{2}B_{1}$ ) bands were assigned. The total emission cross section ( $\sigma_{em}$ ) was evaluated to be (1.7 ± 0.3) × 10<sup>-20</sup> m<sup>2</sup> at a collision energy of 150 meV. The  $\sigma_{em}$ 's of the SH<sup>+</sup>(A–X) and H<sub>2</sub>S<sup>+</sup>( $\tilde{A}-\tilde{X}$ ) bands decreased with increase in the collision energy in the 115–200 meV range, indicating that attractive forces are effective for the incident channels with regard to the formation of these species. The rotational distribution of SH<sup>+</sup>( $A^{3}\Pi$ , v' = 0) is represented by a Boltzmann temperature of 870 ± 80 K. The H<sub>2</sub>S<sup>+</sup>( $\tilde{A}^{2}A_{1}-\tilde{X}^{2}B_{1}$ ) emission, which was assigned

for the first time in the Penning ionization of H<sub>2</sub>S, primarily consists of the bending progressions. The internal populations of  $H_2S^+(\tilde{A})$  were analyzed using the vibrational energies and Einstein's A coefficients calculated in this study. The details of the calculation and derived spectroscopic constants are reported in the accompanying paper, Paper I. The populations obtained for the bending vibration  $(v_2')$  of  $H_2S^+(\tilde{A})$  show an inverted distribution with a peak at  $v_2' = 3$ . This distribution is shifted lower compared that with a peak at  $v_2' = 4-5$  observed by He(2<sup>3</sup>S) Penning ionization electron spectroscopy and that with a peak at  $v_2' = 6-7$ predicted by the theoretical Franck-Condon factors (FCFs) for the  $H_2S(\tilde{X})$ – $H_2S^+(\tilde{A})$  ionization. The origin of the difference is discussed concerning the formation mechanism of  $H_2S^+(A^2A_1)$ .

#### VIII-KK-5 Theoretical Analysis of the Oxygen Insertion Process in the Oxidation Reactions of $H_2O + H/Si(100)$ and $2H + H_2O/Si(100)$ ; Calculation of an Ab Initio Molecular Orbital Method and an Analysis of the Tunneling Reaction

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The reaction paths were analyzed, by an ab initio molecular orbital method, for the surface reaction systems,  $2H + H_2O/Si(100)-(2\times1)$  and  $H_2O + H/Si(100)-(2\times1)$ , in which SiH<sub>2</sub> species with one or two oxygen atom-inserted back bonds have been observed as new stable reaction products. It was found that common metastable states exist in both systems, and the initial energy is sufficiently higher than all transition state energies in the former system, while in the latter system, the energy of the highest transition state is much higher than the initial energy, and thus a tunneling effect plays an important role.

#### VIII-KK-6 Optimal Control of Quantum Chaotic Dynamics

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We numerically investigate how chaos affects controllability of wavepacket dynamics using the monotonically convergent method of optimal control theory introduced by Zhu, Bonita, and Rabitz [J. Chem. Phys. **108**, 1953 (1998)]. We find that a quantized kicked rotor in a bounded phase space is successfully controlled with the method, *i.e.*, an initial Gaussian wavepacket can be steered to a target Gaussian wavepacket with a high probability, even in strongly chaotic regions when the target time T is larger than a minimal time  $T_{min}$ . We also find that  $T_{min}$  saturates to a certain value as the system becomes a random matrix system.