# **RESEARCH ACTIVITIES**

# Theoretical and Computational Molecular Science

It is our goal to develop new theoretical and computational methods based on quantum mechanics, statistical mechanics, and molecular simulation in order to predict and understand the structures, reactions, and functions of molecules in gas, solution, and condensed phases as well as in nano- and bio-systems prior to or in cooperation with experiments.

# Theoretical Study and Design of Functional Molecules: New Bonding, Structures, and Reactions

# Department of Theoretical and Computational Molecular Science Division of Theoretical Molecular Science I



NAGASE, Shigeru LUO, Gangfu KATOUDA, Michio GUO, Jing-Doing KARTHIKEYAN, Subramanlan ZHOU, Xin SAKAKI, Shigeyoshi YAMADA, Mariko KONDO, Naoko Professor ( -March, 2012)\* IMS Fellow<sup>†</sup> Post-Doctoral Fellow<sup>‡</sup> Post-Doctoral Fellow<sup>§</sup> Post-Doctoral Fellow<sup>§</sup> Post-Doctoral Fellow<sup>||</sup> Visiting Scientist\* Secretary Secretary

In theoretical and computational chemistry, it is an important goal to develop functional molecules prior to or in cooperation with experiment. Thus, new bonds and structures provided by heavier atoms are investigated together with the reactivities. In addition, chemical modification and properties of nanocarbons are investigated to develop functional nanomolecular systems. Efficient computational methods are also developed to perform reliable quantum chemistry calculations for small and large molecular systems.

# 1. Quest for Stable Multiple Bonds between Lead Atoms

The heavier analogues of alkynes, REER (E = Si, Ge, Sn, Pb), have attracted special interest in main-group element chemistry. Accordingly, all the heavier analogues have been synthesized and isolated uo to now. The X-ray crystal structure of the heaviest analogue, Ar\*PbPbAr\* (Ar\* =  $C_6H_3$ -2,6- $(C_6H_2$ -2,4,6- $iPr_3$ )<sub>2</sub>), has shown that the Pb–Pb bond distance is much longer than Pb–Pb single bond distances. Theoretical calculations have revealed that Ar\*PbPbAr\* has no  $\pi$  bond between Pb atoms. Therefore, it has been widely accepted that the heaviest Pb analogues of alkynes take a singly bonded structure, unlike the Si, Ge, and Sn cases. However, we have pointed out that Ar\*PbPbAr\* has a triply bonded structure in solution and crystal structures are not very helpful for compounds with bulky groups.<sup>1</sup>

The Pb–Pb bond distance of 3.071 Å calculated for the triply bonded structure of Ar\*Pb–PbAr\* is considerably longer than the Pb–Pb single bond distance of 2.844 Å in Ph<sub>3</sub>PbPbPh<sub>3</sub>. Since we have found that electropositive silyl groups decreases the Pb–Pb bond distance, several bulky silyl groups were tested using density functional calculations.<sup>2)</sup> The optimized structure of R<sup>Si</sup>PbPbR<sup>Si</sup> (R<sup>Si</sup> = Si<sup>i</sup>Pr{CH(SiMe<sub>3</sub>)<sub>2</sub>}<sub>2</sub>) is depiched in Figure 1, the two Pb atoms being triply bonded. The Pb–Pb bond distance of 2.696 Å is remarkably shorter

than the Pb–Pb single bond distance of 2.844 Å in Ph<sub>3</sub>Pb-PbPh<sub>3</sub>, and it is considerably shorter than the shortest Pb–Pb double bond distance of 2.903 Å known to date. However, the triply bonded structure is 4.3 kcal/mol less stable than the singly bonded structure. We are still searching for good substituent groups that stabilize a triply bonded structure with a sufficiently short Pb–Pb bond.



Figure 1. Triply bonded structure of R<sup>Si</sup>PbPbR<sup>Si</sup>.

As a new type of double bonds, diatomic molecules stabilized by the coordination of dative N-heterocyclic carbenes are of considerable interest currently. This stabilization has been performed for the synthesis and isolation of L: $\rightarrow$ Si=Si $\leftarrow$ :L and L: $\rightarrow$ Ge=Ge $\leftarrow$ :L (L = :C[N(2,6-<sup>*i*</sup>Pr<sub>2</sub>-C<sub>6</sub>H<sub>3</sub>)CH]<sub>2</sub>). The Si– Si and Ge–Ge bond distances compare well with typical double bond distances of R<sub>2</sub>Si=SiR<sub>2</sub> and R<sub>2</sub>Ge=GeR<sub>2</sub>. Because the heaviest Pb case remains unknown, we have investigated L: $\rightarrow$ Pb=Pb $\leftarrow$ :L.<sup>2)</sup> The optimized structure is shown in Figure 2. Molecular orbital analysis confirms that the two Pb atoms are doubly bonded. It is notable that the Pb–Pb double bond distance of 2.833 Å is considerably shorter than the shortest Pb–Pb double bond distance of 2.903 Å known to date and differs little from the Pb–Pb bond distance of Pb<sub>2</sub>. It is expected that  $L:\rightarrow Pb=Pb\leftarrow:L$  is an interesting synthetic target.



Figure 2. Optimized structure of L: $\rightarrow$ Pb=Pb $\leftarrow$ :L.

### 2. Short Bonds between Transition Metals

As a novel transition metal complex with an unsupported two-coordinate Fe atom,  $R^*FeFe(\eta^5-C_5H_5)(CO)_2$  ( $R^* = C_6H^2$ , 2, 5, -*i*Pr<sub>3</sub>) has recently 2, 6-Ar<sub>2</sub>-3, 5-*i*Pr<sub>2</sub> where Ar =  $C_6H_2$ -2, 5, 6-*i*Pr<sub>3</sub>) has recently been synthesized and investigated.<sup>3)</sup> X-ray crystal analysis shows that the Fe–Fe bond distances of 2.393 Å is much shorter than the known unsupported Fe–Fe bond distances (2.687–3.138 Å), suggesting that  $R^*FeFe(\eta^5-C_5H_5)(CO)_2$  has the shortest unsupported Fe–Fe bond. Interestingly, calculations show that the Fe–Fe bond is greatly shortened in the bicyclic four-membered ring presented in Figure 3a.<sup>2</sup>) Figure 3. Bicyclic four-membered ring complexes.



It is remarkable that the Fe–Fe bond distance of 2.045 Å is much shorter than the shortest Fe–Fe bond distance of 2.127 Å reported to date. This bond shortening assisted by two Sn atoms is very recently realized for the bicyclic four-membered Ru<sub>2</sub>Sn<sub>2</sub> ring complex presented in Figure 3b.<sup>4)</sup> The Ru–Ru bond distance of 2.343 Å determined through X-ray crustal analysis agrees well with the calculated value of 2.363 Å. It is shorter than the Ru–Ru distance of 2.449–2.469 Å in the related bridge complexes, for which it is assumed that the two Ru atoms are triply bonded. Localized molecular orbital analysis reveals that one clear  $\sigma$  bond exists between the Ru atoms, while two three-centered  $\sigma$  orbitals are delocalized over each of the three-membered Ru<sub>2</sub>Sn rings, which make an important contribution to Ru–Ru bonding. As a result, the Ru–Ru bond has a somewhat multiple-bond character.

It is also predicted theoretically that the Ru–Ru bond distance in the bicyclic four-membered Ru<sub>2</sub>Sn<sub>2</sub> ring is shortened further in the bicyclic six-membered Ru<sub>2</sub>Sn<sub>4</sub> ring. Experimental confirmation is in progress.

### 3. Nano-Carbon Systems

We have performed theoretical calculations for (a) polarized nonresonant Raman spectra of graphene nanoriboins<sup>5)</sup> and (b) transport properties of transition metal intercalated graphene.<sup>6)</sup> In collaboration with experiment, we have also performed calculations for (c) stable radical anions inside fullerene cages,<sup>7)</sup> (d) the core-shell interplay in carbide cluster metallofullerenes,<sup>8)</sup> (e) the cocrystal of La@C<sub>82</sub> and nickel porphyrin with high electron mobility,<sup>9)</sup> and (f) chemical understanding of carbide cluster metallofullerenes.<sup>10)</sup>

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<sup>\*</sup> Present Address; Fukui Institute for Fundamental Chemistry, Kyoto University, Kyoto 606-8103

<sup>†</sup> Present Address; Department of Physics, Peking University, Beijing 100871, P. R. China

<sup>&</sup>lt;sup>‡</sup> Present Address; Riken Advanced Institute for Computational Science, Kobe 650-0047

<sup>§</sup> Present Address; Department of Chemistry, Sungkyunkwan University, Suwon 440-746, Korea

<sup>||</sup> Present Address; Harbin Institute of Technology, Harbin 150080, P. R. China

# Electron and Electromagnetic Field Dynamics in Nanostructures

Department of Theoretical and Computational Molecular Science Division of Theoretical Molecular Science I



NOBUSADA, Katsuyuki YASUIKE, Tomokazu IIDA, Kenji NODA, Masashi YAMADA, Mariko Associate Professor Assistant Professor IMS Fellow Post-Doctoral Fellow Secretary

We have developed theoretical methods to calculate photoinduced electron dynamics in nanostructured materials such as nanoparticles, quantum-dot arrays, and adsorbate-surface systems. Specifically, we have formulated generalized theory of a light-matter interaction beyond a dipole approximation with the aim of understanding the near-field excitation of nanostructures. Furthermore, a highly efficient computational program of massively parallel calculations for electron dynamics has been developed to investigate optical response of nanostructures more than ten-nanometers. Structural and electronic properties of gold-thiolate clusters have also been elucidated in collaboration with an experimental group.

## 1. Massively-Parallel TDDFT Calculations Based on Finite Difference Method in Real-Time and Real-Space

A highly efficient computational program of massively parallel calculations for electron dynamics has been developed in an effort to apply the method to optical response of nanostructures more than ten-nanometers. The approach is based on time-dependent density functional theory calculations in real-time and real-space. The computational code is implemented by using very simple algorithms with a finite difference method in space derivative and Taylor expansion in timepropagation. Since the computational program is free from the algorithms of eigenvalue problems and fast-fourier-transformation, which are usually implemented in conventional quantum chemistry or band structure calculations, the program is highly suitable for massively parallel calculations. The method is applied to optical response of nanostructures constructed from  $C_{60}$  as benchmark systems. We achieved 8.15% peak performance on the K computer with 1920 nodes (15360 cores) and 3.25% peak performance with 12288 nodes (98304 cores). The peak performance decreases with increasing the nodes because of the network communications due to summing up electron density. The computed absorption spectrum of a face-centered cubic unit of solid  $C_{60}$  well reproduces the experimental result.



**Figure 1.** (a) Absorption spectrum of a face-centered cubic unit associated with a solid thin film of  $C_{60}$  in comparison with (b) the experimental observation.

# 2. Enhanced Raman Spectrum of Pyrazine with the Aid of Resonant Electron Dynamics in a Nearby Cluster<sup>1)</sup>

We have investigated the electron dynamics relevant to the

Raman enhancement of pyrazine by a nearby Na<sub>4</sub> cluster. The present time-dependent analysis allows us to illustrate that the electronic excitation in Na4 is closely associated with surfaceenhanced Raman scattering in pyrazine. More specifically, it was clearly demonstrated that some specific enhanced vibrational modes strongly couple with the plasmonic electron motion of Na<sub>4</sub>. The displacement along the enhanced mode leads to a change of the electron dynamics in the entire region of the system, whereas the non-enhanced mode causes only a localized change in the junction area between pyrazine and Na<sub>4</sub>. The strong obstructive mode to the plasmonic electron motion is strongly enhanced. In the present system, all the obstructive modes are readily understood by the characteristics of the atoms in-phase oscillating in the x-direction at the side of Na<sub>4</sub>. This picture gives a clear explanation for the enhancement. The present results show the potential ability of clusterenhanced Raman scattering. A microcluster selectively adsorbed by a specific molecule can be designed using recently developed techniques for the chemical synthesis of clusters. Such artificially designed clusters significantly enhance the Raman scattering of analyte molecules and therefore provide a new way of utilizing enhanced Raman scattering in various research fields.



**Figure 2.** (a) Time-dependent induced dipole moments (*x*-component) in the C<sub>4</sub>H<sub>4</sub>N<sub>2</sub> (red) and Na<sub>4</sub> (blue) sides of C<sub>4</sub>H<sub>4</sub>N<sub>2</sub>–Na<sub>4</sub> under *x*-polarized laser excitation at  $\omega = 1.23$  eV. The dipole moment of isolated C<sub>4</sub>H<sub>4</sub>N<sub>2</sub> under the same excitation condition is also shown (black). Fourier component (1.23 eV) of the time-dependent induced density of (b) Na<sub>4</sub> and (c) C<sub>4</sub>N<sub>2</sub>H<sub>4</sub>–Na<sub>4</sub>.

## 3. Palladium Doping of Magic Gold Cluster Au<sub>38</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub>: Formation of Pd<sub>2</sub>Au<sub>36</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub> with Higher Stability than Au<sub>38</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub><sup>2)</sup>

Gold clusters protected by thiolates have attracted considerable attention as building blocks for new functional materials because they exhibit size-specific physical and chemical properties. Among these, Au<sub>25</sub>(SR)<sub>18</sub>, Au<sub>38</sub>(SR)<sub>24</sub>, Au<sub>68</sub>(SR)<sub>34</sub>, Au<sub>102</sub>(SR)<sub>44</sub>, and Au<sub>144</sub>(SR)<sub>60</sub> are promising, because they exhibit higher thermodynamic and chemical stabilities than clusters of other sizes. Many studies have been conducted on the isolation, size-selective synthesis, stabilities, structures, chemical and physical properties, and applications of these stable clusters. In this study, a phenylethanethiolateprotected Pd<sub>2</sub>Au<sub>36</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub> cluster, which is a two-Pd atom-doped cluster of the well studied magic gold cluster Au<sub>38</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub>, was synthesized in high purity and its stability was investigated. The experimental and theoretical results demonstrate that Pd<sub>2</sub>Au<sub>36</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub> is more stable than Au<sub>38</sub>(SC<sub>2</sub>H<sub>4</sub>Ph)<sub>24</sub> against degradation in solution and core etching by thiols.

### 4. Effect of Copper Doping on Electronic Structure, Geometric Structure, and Stability of Thiolate-Protected Au<sub>25</sub> Nanoclusters<sup>3)</sup>

Several recent studies have attempted to impart [Au<sub>25</sub> (SR)<sub>18</sub>]<sup>-</sup> with new properties by doping with foreign atoms. In this study, we investigated the effect of copper doping on the electronic structure, geometric structure, and stability of  $[Au_{25}(SR)_{18}]^{-}$  with the aim of investigating the effect of foreign atom doping of [Au<sub>25</sub>(SR)<sub>18</sub>]<sup>-</sup>. Cu<sub>n</sub>Au<sub>25-n</sub>(SC<sub>2</sub>H<sub>4</sub> Ph)<sub>18</sub> was synthesized by reducing complexes formed by the reaction between metal salts (copper and gold salts) and PhC<sub>2</sub> H<sub>4</sub>SH with NaBH<sub>4</sub>. Mass analysis revealed that the products contained  $Cu_nAu_{25-n}(SC_2H_4Ph)_{18}$  (n = 1-5) in high purity. Experimental and theoretical analysis of the synthesized clusters revealed that copper doping alters the optical properties and redox potentials of the cluster, greatly distorts its geometric structure, and reduces the cluster stability in solution. These findings are expected to be useful for developing design guidelines for functionalizing [Au<sub>25</sub> (SR)<sub>18</sub>]<sup>-</sup> through doping with foreign atoms.

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# Advanced Electronic Structure Theory in Quantum Chemistry

# Department of Theoretical and Computational Molecular Science Division of Theoretical Molecular Science I



YANAI, Takeshi KURASHIGE, Yuki CHALUPSKY, Jakub TRAN, Lan SAITOW, Masaaki YAMADA, Mariko Associate Professor Assistant Professor IMS Fellow Graduate Student Graduate Student\* Secretary

Aiming at predictive computational modelings of molecular electronic structures with ab initio quantum chemistry calculations, our scientific exploration is to establish a cuttingedge theoretical methodology that allows one to compute accurately and efficiently the complex electronic structures, in which strongly-interacting electrons play a crucial role to characterize the nature of molecules. The complicated electronic structures can be handled accurately with the multireference theory, which deals with multiple important electronic configurations on equal footing. However, with the standard multireference methods such as the complete active space self-consistent field (CASSCF), the tractable size of the reference space is limited to small active space because the complexity of the calculations grows exponentially with the reference size. The existing multireference methods are nevertheless usefully applied to chemical theory problems such as exploring chemical reactions of bonding, dissociation and isomerization along the reaction coordinates, electronically excited states, unstable electronic structures of radical systems, and multiple covalent bindings in molecular metal complexes, etc. Our resultant works to be reported here are (1) to develop a type of the multireference correlation model named Canonical Transformation (CT) theory, which can efficiently describe short-range dynamic correlation on top of the multi-configurational staring wave function, (2) to construct the extensive complete active space self-consistent field (CASSCF) method combined with ab initio density matrix renormalization group (DMRG) method for making unprecedentedly larger active spaces available for the CASSCF calculations, and (3) to develop an efficient second-order perturbation theory that can use large active space with the DMRG-SCF reference wavefunction.

# 1. Canonical Transcorrelated Theory with Projected Slater-Type Geminals<sup>1)</sup>

An effective Hamiltonian perturbed with explicit inter-

Canonical transcorrelated theory



Figure 1. Canonical Transcorrelation method with F12 correlation factor.

electronic correlation is derived from similarity transformation of Hamiltonian using a unitary operator with Slater-type geminals. The Slater-type geminal is projected onto the excitation (and de-excitation) component as in the F12 theory. Simplification is made by truncating higher-body operators, resulting in a correlated Hamiltonian which is Hermitian and has exactly the same complexity as the original Hamiltonian in the second quantized form. It can thus be easily combined with arbitrary correlation models proposed to date. The present approach constructs a singularity-free Hamiltonian a priori, similarly to the so-called transcorrelated theory, while the use of the canonical transformation<sup>2)</sup> assures that the effective Hamiltonian is two-body and Hermite. Our theory is naturally extensible to multireference calculations on the basis of the generalized normal ordering. The construction of the effective Hamiltonian is non-iterative. The numerical assessments demonstrate that the present scheme improves the basis set convergence of the post-mean-field calculations at a similar rate to the explicitly correlated methods proposed by others

that couple geminals and conventional excitations.

In this study, we propose a canonical transcorrelated Hamiltonian with the F12 operator:

$$\hat{H}^{F12} \equiv \hat{H} + [\hat{H}, \hat{A}^{F12}]_{12} + \frac{1}{2}[[\hat{F}, \hat{A}^{F12}]_{12}, \hat{A}^{F12}]_{12}$$

which is derived in two ways of formal approximations: (i) terminating the expansion at the second order and (ii) replacing the uncorrelated Hamiltonian  $\underline{H}$  at the second order term (*i.e.*, the double commutator) by the Fock operator F. Note that F is the effective one-particle approximation to  $\underline{H}$ . This truncation of the infinite expansion is correct through the second order in perturbation.

We use an anti-Hermitian generator with projected geminal functions to make the transcorrelated *H* Hermite:

$$\begin{split} \hat{A}^{F12} &= \frac{1}{2} G_{ij}^{\alpha\beta} (\hat{E}_{ij}^{\alpha\beta} - \hat{E}_{\alpha\beta}^{ij}), \\ G_{ij}^{\alpha\beta} &= \frac{3}{8} \langle \alpha\beta | \hat{Q}_{12}F_{12} | i j \rangle + \frac{1}{8} \langle \alpha\beta | \hat{Q}_{12}F_{12} | j t \rangle, \\ F_{12} &= -\gamma^{-1} \exp(-\gamma r_{12}), \end{split}$$

in which we have fixed the amplitudes by those determined by the first-order cusp condition. The explicit electron correlation with the *projected* Slater-type geminals is built into a Hamiltonian through the canonical transformation. The present approach provides a formulation to effectively remove high-energy orbital components by using the F12 factor as a regulator for renormalizing them into the smaller-size orbital space. The features of the canonical transcorrelated theory are:

- The resulting effective Hamiltonian is already perturbed with a considerable amount of the dynamic correlation associated with the interelectronic Coulomb singularity;
- It remains Hermitian and has exactly the same size, dimension, and quartic complexity as the bare Hamiltonian;



Figure 2. Determination of oxidation state of  $Mn_4CaO_5$  cluster in photosystem II from a multireference wavefunction theory based on Denisty Matrix Renormalization Group method.

- (3) There is no adjustable parameter in the geminal excitations with Ten-no's fixed amplitude ansatz, since the F12 amplitudes for the present transformation are predetermined and calculated in a non-iterative manner;
- (4) The theory is extensible to multireference models on the basis of the generalized normal ordering of Mukherjee and Kutzelnigg;
- (5) In contrast to the standard F12 theories which couple the F12 and conventional excitations in the amplitude equations or Valeev's *a posteori* F12 corrections, we have introduced an *a priori* F12 transformed Hamiltonian that can be readily used in conjunction with arbitrary correlation models to describe the remaining orbital correlation.

We have demonstrated its applications to various solvers in quantum chemical methods, such as CCSD(T), QCISD, MP4, CCSDT, CCSDTQ, and so forth. The benchmarks on small molecules have revealed that the numerical performance of our explicit correlation scheme is comparable to that of other F12 theories. In our method, the F12 correction and the orbital correlation are treated separately at different steps, and thus it is indicated that they are more or less additively separable. (**Figure 1**)

### 2. Oxidation State of Mn<sub>4</sub>Ca Cluster in Photosystem II: A Quantum-Chemical Density-Matrix Renormalization Group Study<sup>3)</sup>

The X-ray diffraction (XRD) structure of photosystem II at a resolution of 1.9 Å was recently reported, entailing atomic-details of the  $Mn_4Ca$  cluster. Meanwhile, there is an earlier study suggesting that the high-valent  $Mn^{III-IV}$  ions of the cluster are potentially damaged by X-rays and reduced towards the lower-valent  $Mn^{II}$ , involving structural deformation. Thus, the record-resolution XRD measurement used a low-level X-ray dose to avoid such damage. We report a theoretical analysis identifying the oxidation states of the Mn ions as fingerprints to be compared with the widely-accepted oxidation state. Super high-dimensional multireference wavefunctions were calculated to account for nonperturbative interactions arising from four Mn 3d and five  $\mu$ -oxo 2p shells. (**Figure 2**)

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# Developing the Statistical Mechanics Theory of Liquids in Chemistry and Biophysics

Department of Theoretical and Computational Molecular Science Division of Theoretical Molecular Science II



HIRATA, Fumio YOSHIDA, Norio MARUYAMA, Yutaka PHONGPHANPHANEE, Saree SINDHIKARA, Daniel J. KIYOTA, Yasuomi SUETAKE, Yasumi KONDO, Naoko YAMADA, Mariko

Professor ( -March, 2012)\* Assistant Professor<sup>†</sup> Post-Doctoral Fellow Post-Doctoral Fellow Post-Doctoral Fellow Secretary Secretary Secretary

"Molecular recognition" is an essential elementary process for protein to function. The process is a thermodynamic process which is characterized with the free energy difference between two states of a host-guest system, namely, associated and dissociated states. It is readily understood that the structural fluctuation of protein gives a big effect on the free energy barrier. In that respect, the "molecular recognition" is a thermodynamic process which is conjugated with the structural fluctuation of protein.

We have been developing a new theory concerning the molecular recognition, based on the 3D-RISM/RISM theory which is a statistical mechanics of liquids. The theory has successfully "probed" small ligands such as water molecules and ions bound in a small cavity of protein.<sup>1–3)</sup>

# 1. Elucidating the Molecular Origin of Hydrolysis Energy of Pyrophosphate in Water<sup>4)</sup>

The molecular origin of the energy produced by the ATP hydrolysis has been one of the long-standing fundamental issues. A classical view is that the negative hydrolysis free energy of ATP originates from intra-molecular effects connected with the backbone P-O bond, so called "high-energy bond." On the other hand, it has also been recognized that solvation effects are essential in determining the hydrolysis free energy. Here, using the 3D-RISM-SCF (three-dimensional reference interaction site model self-consistent field) theory that integrates the ab initio quantum chemistry method and the statistical mechanical theory of liquids, we investigate the molecular origin of hydrolysis free energy of pyrophosphate, an ATP analog, in water. We demonstrate that our theory quantitatively reproduces the experimental results without the use of empirical parameters. We clarify the crucial role of water in converting the hydrolysis free energy in the gas phase determined solely by intra-molecular effects, which ranges from endothermic, thermoneutral to highly exothermic depending on the charged state of pyrophosphate, into moderately

exothermic in the aqueous phase irrespective of the charged state as observed in experimental data. We elucidate that this is brought about by different natures of solute–water interactions depending on the charged state of solute species: the hydration free energy of low-charged state is mainly subjected to shortrange hydrogen-bonds, while that of high-charged state is dominated by long-range electrostatic interactions. We thus provide unambiguous evidence on the critical role of water in determining the ATP hydrolysis free energy.

A HO-P-O-P-OH + H<sub>2</sub>O 
$$\longrightarrow$$
 2 HO-P-OH  
 $\dot{O}H$   $\dot{O}H$  + H<sub>2</sub>O  $\longrightarrow$  2 HO-P-OH  
 $\dot{O}H$   $\dot{O}H$   
B  $\begin{array}{c} 0 & 0 \\ O-P-O-P-OH + H_2O & \longrightarrow HO-P-OH + HO-P-O
 $\dot{O}H$   $\dot{O}H$   $\dot{O}H$   
C  $\begin{array}{c} 0 & 0 \\ O-P-O-P-O & + H_2O & \longrightarrow HO-P-O
 $\dot{O}H & \dot{O}H \end{array}$   $\dot{O}H$   
D  $\begin{array}{c} 0 & 0 \\ O-P-O-P-O & + H_2O & \longrightarrow HO-P-O
 $\dot{O}H & \dot{O}H \end{array}$   $\dot{O}H$   $\dot{O}H$   $\dot{O}H$$$$ 

**Scheme 1.** Schematic description of the hydrolysis reaction of pyrophosphate for the four possible charged states.

**Table 1.** The reaction free energies in the gas phase and in the aqueous phase at 298.15 K and 1.0 atm computed by DFT at the B3LYP/6-31+(d) level and the 3D-RISM-SCF theory. Units are in kcal/mol.

Reaction	gas phase	aqueous phase	Exp.
	(DFT)	(3D-RISM/RISM)	
А	-1.7	-8.9	-9.5
В	21.3	-6.2	-7.5
С	-56.6	-8.1	-7.7
D	-119.0	-7.7	-7.1

# 2. Placevent: An Algorithm for Predicting of Explicit Solvent Atom Distribution— Application to HIV-1 Protease and F-ATP Synthase<sup>5)</sup>

Location of water and ions in native structure of protein is of essential importance for its stability and for its functions. However, determination of the position of those species in protein is not an easy task for any experimetal methods currently available, X-ray, NMR, neutron diffraction, and the molecular simulation.

We have created a simple algorithm for automatically predicting the explicit solvent atom distribution of biomolecules. The explicit distribution is coerced from the 3D continuous distribution resulting from a 3D-RISM calculation. This procedure predicts optimal location of solvent molecules and ions given a rigid biomolecular structure and the solvent composition. We show examples of predicting water molecules near the KNI-272 bound form of HIV-1 protease and predicting both sodium ions and water molecules near the rotor ring of F-ATP synthase. Our results give excellent agreement with experimental structure with an average prediction error of 0.45–0.65 Å. Further, unlike experimental methods, this method does not suffer from the partial occupancy limit. Our method can be performed directly on 3D-RISM output within minutes. It is extremely useful for examining multiple specific solventsolute interactions, as a convenient method for generating initial solvent structures for MD calculations, and may assist in refinement of experimental structures.



Figure 1. Water molecules near KNI-275. KNI-275 is shown as a translucent surface, HIV-1 protease as a cartoon. Crystal water molecules near KNI-275 are shown in orange, elsewhere in red. Water molecules placed by this method near KNI-275 are shown in green, elsewhere in pink.

### 3. Modified Andersen Method for Accelerating 3D-RISM Calculations Using Graphics Processing Unit<sup>6)</sup>

Increasing attention has been paid to the 3D-RISM theory due mainly to its capability of treating "solvation" of bio-

\* Present Address; College of Life Sciences, Ritsumeikan University † Present Address; Graduate School of Sciences, Kyushu University molecules such as protein and DNA without using any adjustable parameters, which is the case in the continuum model. The method was highlighed in several symposiums in the latest ACS meeting held in Philadelphia. Superiority of 3D-RISM to the continuum models, the Poisson-Boltzmann and generalized Born equations, was addressed unambiguously by several talks in the symposiums as far as physical soundness, amount of information produced, applicability to drug design, and so on, are concerned. However, there still remains one point which makes people stick to the continuum models. That is the computation cost. The cost to perform the 3D-RISM calculation is far higher than that of the continuum models. We have proposed a fast algorithm to solve the 3D-RISM equation on a graphics processing unit (GPU). It was the large memory space required for convergence of iteration that banned 3D-RISM from GPU. In order to overcome the difficulty, we replaced the conventional MDIIS algorithm by Anderson's method with some modification. Using this method on a Tesla C2070 GPU, we reduced the total computational time by a factor of eight, 1.4 times by the modified Andersen method and 5.7 times by GPU, compared to calculations on an Intel Xeon machine (8 cores, 3.33 GHz) with the conventional method.



**Figure 2.** Root mean square residual against the number of iteration steps for the calculation of 3D site distribution profiles of water around a DNA molecule by 3D-RISM. Solid, dotted, and dashed lines are for MDIIS, Anderson, and Modified Anderson algorithms, respectively.

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# **Theory of Photoinduced Phase Transitions**

Department of Theoretical and Computational Molecular Science Division of Theoretical Molecular Science II



YONEMITSU, Kenji TANAKA, Yasuhiro NISHIOKA, Keita MIYAZAKI, Mitake KONDO, Naoko Associate Professor ( –March, 2012)\* Assistant Professor IMS Fellow<sup>†</sup> Visiting Scientist<sup>‡</sup> Secretary

Photoirradiation of materials usually creates electrons and holes, which are often accompanied by local structural deformation. With the help of cooperativity, the electronic and/or structural deformation can proliferate to change the physical property such as conductivity, permittivity, and magnetic susceptibility. The resultant nonequilibrium phase may not be reached by changing temperature or pressure because the energy of a photon is much higher than thermal energies. Our theoretical researches are focused on the mechanisms and dynamics of photoinduced phase transitions, how they are controlled, and how the photoinduced electron–lattice states are different from those which are realized in thermal equilibrium.<sup>1</sup>

### 1. Normal Mode Analysis for Intra- and Inter-Molecular Electron–Phonon Coupled Systems

The ground-state properties of a molecular compound  $Et_2Me_2Sb[Pd(dmit)_2]_2$  are theoretically studied, which shows a dimer-Mott character in Pd(dmit)\_2 layers at high temperature and charge order mainly stabilized by electron–phonon interactions at low temperature. An effective extended Peierls-Hubbard model is constructed with intra- and inter-molecular electronic and phonon degrees of freedom. Using a mean-field approximation, the energies and optimized structures are calculated for isolated neutral, monovalent, divalent Pd(dmit)\_2 dimers and their two-dimensional crystallized states. The optical conductivities of the latter are calculated by a single configuration interaction method. Through these numerical calculations, model parameters have been evaluated by comparing the theoretical and experimental ground-state properties.

Then, the normal mode analysis is performed for the intradimer C=C stretching vibrations and intradimer intermonomer stretching vibrations. Molecular vibrations with different symmetries are coupled to different combinations of electrons and holes within a molecule, so that their frequencies depend on the molecular charge in different manners. This



**Figure 1.** C=C stretching phonon mode B (infrared active) schematically represented in the upper panel. Its time evolution (middle panel) and time-frequency spectrogram (lower panel) on the initially neutral dimer are shown when the charge-separated state is photoexcited.

information is useful to analyze photoinduced transient states, where the relation between their frequencies and the molecular charge is modified from the equilibrium counterpart.



**Figure 2.** C=C stretching phonon mode C (infrared active) schematically represented in the upper panel. Its time evolution (middle panel) and time-frequency spectrogram (lower panel) on the initially neutral dimer are shown when the charge-separated state is photoexcited.

### 2. Photoinduced Dynamics in Intra- and Inter-Molecular Electron-Phonon Coupled Systems<sup>2)</sup>

For the photoinduced dynamics in  $Et_2Me_2Sb[Pd(dmit)_2]_2$ , the time-dependent Schrödinger equation is numerically solved for the Hartree-Fock wave function on the cluster of sixteen dimers (eight sites in each dimer, 128 sites in total). The classical equation of motion is solved for phonons. Photoexcitation is introduced through the Peierls phase in the transfer integrals.

We focus on intradimer C=C stretching vibrations that are infrared active and experimentally observed so far, which are schematically shown in the upper panels of Figures 1 and 2. The middle panels show the time evolution of their amplitudes on the initially neutral dimer when the charge-separated state is photoexcited. The lower panels show their time-frequency spectrograms.

The initially neutral and divalent dimers become monovalent at about 0.5 ps after photoexcitation. After 0.5 or 1 ps, their oscillations are noticeable and sinusoidal. The frequency of the B mode is actually close to that in a monovalent dimer in the equilibrium state. The B mode is coupled with the antibonding-LUMO (electron)–antibonding-HOMO (hole) excitation within the dimer. Their correlation is recovered and becomes close to that in equilibrium after 1 ps.

However, the frequency of the C mode after 1 ps is not close to that in a neutral dimer or that in a monovalent dimer, but it is rather close to the bare frequency (*i.e.*, the frequency if the couplings with electron–hole excitations are turned off) of  $1331 \text{ cm}^{-1}$ . The C mode is coupled with the antibonding-HOMO (electron)–bonding-HOMO (hole) excitation and the antibonding-LUMO (electron)–bonding-LUMO (hole) excitation within the dimer, so that it is largely softened in equilibrium. Their correlations are largely modified and far from that in equilibrium after 1 ps. This phonon mode appears to oscillate almost independently of electron–hole excitations.

Thus, a usual picture based on the adiabatic potential fails even if phonons are treated classically. The correlation between a phonon and its relevant electron–hole excitation(s) depends on the mode or on the symmetries of the vibrational pattern and the wave functions of electron(s) and hole(s). The analysis of these correlations at an early stage of photoinduced chargeorder melting is essential for future manipulation of nonequilibrium phases.

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- <sup>†</sup> Present Address; Department of Physics, Faculty of Science and Engineering, Chuo University, 1-13-27 Kasuga, Bunkyo-ku, Tokyo 112-8551
- ‡ from Hakodate National College of Technology

<sup>\*</sup> Present Position; Professor, Department of Physics, Faculty of Science and Engineering, Chuo University, 1-13-27 Kasuga, Bunkyo-ku, Tokyo 112-8551

# Theoretical Studies on Condensed Phase Dynamics

Department of Theoretical and Computational Molecular Science Division of Computational Molecular Science



SAITO, Shinji KIM, Kang HIGASHI, Masahiro YAGASAKI, Takuma IMOTO, Sho KAWAGUCHI, Ritsuko Professor Assistant Professor JSPS Post-Doctoral Fellow Post-Doctoral Fellow Graduate Student Secretary

Liquids and biological systems show complicated dynamics because of their structural flexibility and dynamical hierarchy. Understanding these complicated dynamics is indispensable to elucidate chemical reactions and relaxation in solutions and functions of proteins. In this year, we investigated ultrafast proton transfer in solution<sup>1)</sup> and dynamics of liquid and supercooled states.<sup>2–5)</sup> In particular, we examined complicated dynamics in terms of nonlinear response functions and multi-time correlation functions.

## 1. Direct Simulation of Excited-State Intramolecular Proton Transfer and Vibrational Coherence of 10-Hydroxybenzo[*h*]quinoline in Solution<sup>1)</sup>

We investigate an ultrafast excited-state intramolecular proton transfer (ESIPT) reaction and the subsequent coherent vibrational motion of 10-hydroxybenzo-[h]quinoline in cyclohexane by the electronically embedded multiconfiguration Shepard interpolation method, which enables us to generate the potential energy surface of the reaction effectively and thus carry out a direct excited-state dynamics simulation with low computational costs. The calculated time scale of the ESIPT and the frequencies and lifetimes of coherent motions are in good agreement with the experimental results. The present study reveals that the coherent motions are caused by not only the proton transfer itself but also the backbone displacement induced by the ESIPT. We also discuss the effects of the solvent on the dynamics of the coherent vibrational modes.

## 2. Insights in Quantum Dynamical Effects in the Infrared Spectroscopy of Liquid Water from a Semiclassical Study with an *Ab Initio*-Based Flexible and Polarizable Force Field<sup>2)</sup>

The dynamical properties of liquid water play an important role in many processes in Nature. In this paper we focus on the infrared (IR) absorption spectrum of liquid water based on the linearized semiclassical initial value representation (LSC-IVR) with the local Gaussian approximation (LGA) [Liu and Miller, J. Chem. Phys. 131, 074113 (2009)] and an ab initio based, flexible, polarizable Thole-type model (TTM3-F) [Fanourgakis and Xantheas, J. Chem. Phys. 128, 074506 (2008)]. Although the LSC-IVR (LGA) gives the exact result for the isolated 3-dimensional shifted harmonic stretching model, it yields a blue-shifted peak position for the more realistic anharmonic stretching potential. By using the short time information of the LSCIVR correlation function, however, it is shown how one can obtain more accurate results for the position of the stretching peak. Due to the physical decay in the condensed phase system, the LSC-IVR (LGA) is a good and practical approximate quantum approach for the IR spectrum of liquid water. The present results offer valuable insight into future attempts to improve the accuracy of the TTM3-F potential or other ab intio-based models in reproducing the IR spectrum of liquid water.

## 3. Energy Relaxation of Intermolecular Motions in Supercooled Water and Ice: A Molecular Dynamics Study<sup>3)</sup>

We investigate the energy relaxation of intermolecular motions in liquid water at temperatures ranging from 220 K to 300 K and in ice at 220 K using molecular dynamics simulations. We employ the recently developed frequency resolved transient kinetic energy analysis, which provides detailed information on energy relaxation in condensed phases like two-color pump-probe spectroscopy. It is shown that the energy cascading in liquid water is characterized by four processes. The temperature dependences of the earlier three processes, the rotational-rotational, rotational-translational, and translational-translational energy transfers, are explained in terms of the density of states of the intermolecular motions. The last process is the slow energy transfer arising from the transitions between potential energy basins caused by the excitation of the low frequency translational motion. This process is absent in ice because the hydrogen bond network rearrangement, which accompanies the interbasin transitions in liquid water, cannot take place in the solid phase. We find that the last process in supercooled water is well approximated by a stretched exponential function. The stretching parameter,  $\beta$ , decreases from 1 to 0.72 with decreasing temperature. This result indicates that the dynamics of liquid water becomes heterogeneous at lower temperatures.

# 4. Fluctuations and Dynamics of Liquid Water Revealed by Nonlinear Spectroscopy<sup>4)</sup>

Many efforts have been devoted to elucidate the intra- and intermolecular dynamics in liquid water because of its important roles in many fields of science and engineering. Multidimensional nonlinear spectroscopy is a powerful tool to investigate the dynamics. Since nonlinear response functions are described by more than one time variable, it is possible to analyze static and dynamic mode couplings. Here, we review the intra- and intermolecular dynamics of liquid water revealed by recent nonlinear spectroscopic experiments and computer simulations. In particular, we discuss the anharmonic coupling, population relaxation, anisotropy decay, and spectral diffusion of intra- and intermolecular motions of water and their temperature dependence, which play important role in ultrafast dynamics and relaxations in water.

# 5. Anomalous Temperature Dependence of Isobaric Heat Capacity of Water below 0 °C<sup>5)</sup>

When pure liquid water is cooled below its freezing/ melting temperature, it exhibits a number of striking anomalies. Most remarkable among these anomalies is the temperature dependence of the isobaric heat capacity,  $C_{\rm P}$ , that exhibits first a rise, and then a fall, on lowering temperature substantially below 0 °C.<sup>1-4)</sup> In contrast, the isochoric heat capacity,  $C_{\rm V}$ , remains weakly temperature dependent and displays no such anomaly. The reason for this surprisingly large difference is not well understood. To understand this and other anomalies of low temperature water, we examine both wave number and frequency dependent temperature fluctuation by long molecular dynamics simulations. Significant differences between constant pressure and constant volume conditions appear below 240 K in the spatio-temporal correlation of temperature fluctuation. Shell-wise decomposition of relative contribution to the temperature fluctuation reveals an increase in contribution from the distant regions, extending even up to the *fifth* hydration shell, at low temperatures, more significant under isobaric than under isochoric conditions. While the temperature fluctuation time correlation function (TFCF) exhibits the expected slow-down with lowering temperature, it shows a rather surprisingly sharp crossover from a markedly fragile to a weakly fragile liquid around 220 K. We establish that this crossover of TFCF (and the related anomalies) arises from a percolation transition in the population of clusters made of liquid-like molecules, defined by coordination number (and consistent with local volumes obtained from Voronoi polyhedra). The disappearance of large liquid-like clusters below 220 K display characteristic features consistent with theory<sup>5)</sup> of percolation. As temperature is further lowered, TFCF exhibits a power law decay and the relaxation time, when fitted to Vogel-Fulcher-Tammann law, reveal a dynamic transition around 160-170 K. Our computed two-dimensional IR and Raman spectra both also signal a dynamical transition around 170 K and additionally carry signatures of the percolation transition at 220 K that could be measured experimentally.

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# Theoretical Study on Molecular Excited States and Chemical Reactions

Department of Theoretical and Computational Molecular Science Division of Computational Molecular Science



EHARA, Masahiro FUKUDA, Ryoichi TASHIRO, Motomichi BOBUATONG, Karan CHITHONG, Rungtiwa MAITARAD, Phornphimon MEEPRASERT, Jittima HORIKAWA, Takenori PROMKATKAEW, Malinee KAWAGUCHI, Ritsuko Professor Assistant Professor IMS Research Assistant Professor Post-Doctral Fellow\* Visiting Scientist Visiting Scientist Visiting Scientist Graduate Student Graduate Student<sup>†</sup> Secretary

Molecules in the excited states show characteristic photophysical properties and reactivity. We investigate the molecular excited states and chemical reactions which are relevant in chemistry, physics, and chemical biology with developing the highly accurate electronic structure theory. We are also interested in the excited-state dynamics and energy relaxation so that we also develop the methodology of large-scale quantum dynamics. In this report, we present our recent studies on the development of the CAP/SAC-CI method,<sup>1)</sup> electronic spectra of annulated dinuclear free-base phthalocyanine,<sup>2)</sup> and aerobic oxidation of methanol to formic acid on  $Au_{20}^{-}$  cluster.<sup>3)</sup>

### 1. Development of CAP/SAC-CI Method for Calculating Resonance States of Metastable Anions<sup>1)</sup>

A "resonance" is an electronically metastable state, that is, a state of an (N+1)-electron system that lies energetically above the ground state of the associated *N*-electron system and can consequently decay by electron autodetachment. Resonances are intermediates in electron-induced processes and electron-catalyzed reactions. Resonances are part of the continuum and are represented by non-square-integrable (non- $L^2$ ) wavefunctions. Computational methods for resonances are thus necessarily combinations of a method to address the continuum nature of the state and a method to address its many-body nature. It is this combination that renders computing the resonance parameters,  $E_r$  and  $\tau$ , of a many-electron system a challenging task.

In this work, we have developed the complex absorbing potential (CAP)/SAC-CI method to investigate resonance states of metastable anions. The method has been implemented in the projected scheme and applied to the  $\pi^*$  resonance state of formaldehyde. The dependence on both valence and diffuse basis sets up to *g*-function, the number of SAC-CI states in the projection, and the effect of perturbation selection are examined. The potential energy curve and decay width are calculated in the C–O stretching coordinate (Figure 1), and the Franck-Condon factors for transitions from neutral to resonance state

are evaluated to interpret the electron transmission (ET) spectrum. (Figure 2)







**Figure 2.** Vibrational structure in  $\pi^*$  resonance of H<sub>2</sub>CO by the projected CAP/SAC-CI method compared with the electron transmission spectrum.

# 2. Electronic Spectra of Annulated Dinuclear Free-Base Phthalocyanines<sup>2)</sup>

There are many applications for organic dyes with strong photoabsorption in the near-IR region, including photo-energy conversion, molecular sensors and devices, and biological applications. A reliable theory that can treat large conjugated system is urgently required for molecular designs and analysis of the electronic structure of near-IR absorbing materials. The direct SAC-CI method is accurate and efficient for studying large conjugated molecules.

The electronic excited states and electronic absorption spectra of annulated dinuclear free-base phthalocyanine (C<sub>58</sub>H<sub>30</sub>N<sub>16</sub>) are studied through quantum chemical calculations using the SAC-CI method. Three tautomers are possible with respect to the position of the pyrrole protons; therefore, the SAC-CI calculations for these tautomers were performed. The lower energy shift of the Q-bands because of dimerization is explained by the decrease in the HOMO-LUMO gaps resulting from the bonding and antibonding interactions between the monomer units. The relative energies of these tautomers are examined using DFT calculations for several peripheral substituents. The relative energies of these tautomers significantly depend on the substituents, and therefore, the abundance ratios of the three tautomers were affected by the substituents. The absorption spectra were simulated from the SAC-CI results weighted by the Boltzmann factors obtained from the DFT calculations. The SAC-CI spectra reproduce the experimental findings well. The thermal-averaged SAC-CI spectra could explain the observed substituent effect on the structure of the Q-bands in terms of the relative stabilities and the abundance ratios of the tautomers.



**Figure 3.** Molecular structure of annulated dinuclear free-base phthalocyanine and its electronic absorption spectra.

# 3. Aerobic Oxidation of Methanol to Formic Acid on $Au_{20}^{-}$ Cluster<sup>3)</sup>

During the past two decades, the chemical transformation of hydrocarbons catalyzed by nanometer-sized gold clusters has been of significant interest in industrial and academic research because of its remarkable potential for green chemistry and economic significance Useful and practical reactions with gold catalysts have been extensively developed since the pioneering work of Haruta and co-workers. Among these reactions, aerobic oxidation (oxidation with molecular oxygen) on gold clusters has received special attention because it enables us to perform such reactions at ambient conditions and low temperatures, providing high selectivity to the desired products.

In this work, we investigated the aerobic oxidation of methanol to formic acid catalyzed by Au<sub>20</sub><sup>-</sup> using density functional theory with the M06 functional. Possible reaction pathways are examined taking account of full structure relaxation of the Au<sub>20</sub><sup>-</sup> cluster. The proposed reaction mechanism consists of three elementary steps (Figure 4): (1) formation of formaldehyde from methoxy species activated by a superoxolike anion on the gold cluster; (2) nucleophilic addition by the hydroxyl group of a hydroperoxyl-like complex to formaldehyde resulting in a hemiacetal intermediate; and (3) formation of formic acid by hydrogen transfer from the hemiacetal intermediate to atomic oxygen attached to the gold cluster. A comparison of the computed energetics of various elementary steps indicates that C-H bond dissociation of the methoxy species leading to formation of formaldehyde is the rate-determining step. A possible reaction pathway involving single-step hydrogen abstraction, a concerted mechanism, is also discussed. The stabilities of reactants, intermediates and transition state structures are governed by the coordination number of the gold atoms, charge distribution, cooperative effect and structural distortion, which are the key parameters for understanding the relationship between the structure of the gold cluster and catalytic activity in the aerobic oxidation of alcohol.



**Figure 4.** Energy diagram of reaction pathway I in the aerobic oxidation of methanol to formic acid on  $Au_{20}^{-}$  (kcal/mol).

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<sup>†</sup> carrying out graduate research on Cooperative Education Program of IMS with Kasetsart University

# Development of New Algorithms for Molecular Dynamics Simulation and Its Application to Biomolecular Systems

Department of Theoretical and Computational Molecular Science Division of Computational Molecular Science



OKUMURA, Hisashi ITOH, G. Satoru MORI, Yoshiharu NOMURA, Hitomi KAWAGUCHI, Ritsuko Associate Professor Assistant Professor IMS Research Assistant Professor Graduate Student\* Secretary

In the conventional canonical-ensemble simulations, it is difficult to realize efficient samplings in proteins because the simulations tend to get trapped in a few of many local-minimum states. To overcome these difficulties, we have proposed new generalized-ensemble algorithms, such as multibaricmultithermal algorithm, partial multicanonical algorithm, van der Waals replica exchange method, and Coulomb replica exchange method. It is important to realize efficient samplings in the conformational space and to predict the native structures of proteins. We apply these methods to proteins and peptides.

### 1. Temperature and Pressure Denaturation of Chignolin: Folding and Unfolding Simulation by Multibaric-Multithermal Molecular Dynamics Method

We performed a multibaric-multithermal molecular dynamics (MD) simulation of a 10-residue protein, chignolin and discussed its folding thermodynamics last year. We further investigated the denaturation mechanisms this year.<sup>1)</sup>

All-atom model for the protein with Amber parm99SB force field were employed in explicit TIP3P water. This MD simulation covered wide ranges of temperature between  $T = 260 \times 560$  K and pressure between  $P = 0.1 \times 600$  MPa and sampled many conformations without getting trapped in local-minimum free-energy states.

Radial distribution function g(r) between the chignolin heavy atoms and the water oxygen atoms, the first peak position of g(r),  $r_1$ , the number of hydration water molecules, and the number of hydrophobic contacts were calculated. As the temperature increases,  $r_1$  increases and the number of water molecules around chignolin decreases. It represents that chignolin gains more space to move around and is transferred to a high entropy state, even if it losses potential energy by breaking hydrogen bonds between the protein atoms or between the protein and water. Entropy of the transition state and the unfolded state is higher than the native state. This is the reason why the probabilities of not only the unfolded state such as the extended structure but also the transition state increase as the temperature increases.

On the other hand, the number of hydration water molecules increases with the increasing pressure. There are fewer water molecules near the hydrophobic residues when the protein is folded at the room pressure. As the pressure increases, more water molecules come closer to the hydrophobic residues and the number of hydrophobic contacts decreases. The hydrophobic residues then get separated from one another and the protein is unfolded.

### 2. Coulomb Replica-Exchange Method: Handling Electrostatic Attractive and Repulsive Forces for Biomolecules

We propose a new type of the Hamiltonian replica-exchange method for molecular dynamics (MD) and Monte Carlo simulations, which we refer to as the Coulomb replica-exchange method. In this method, electrostatic charge parameters in the Coulomb interactions are exchanged among replicas while temperatures are exchanged in the usual replica-exchange method. By varying the atom charges, the Coulomb replicaexchange method overcomes free-energy barriers and realizes more efficient sampling in the conformational space than the replica-exchange method. Furthermore, this method requires only a smaller number of replicas because only the atom charges of solute molecules are employed as exchanged parameters.

We performed Coulomb replica-exchange MD simulations of an alanine dipeptide in explicit water solvent and compared the results with those of the conventional canonical, replicaexchange, and van der Waals replica-exchange methods. Two force fields of AMBER parm99 and AMBER parm99SB were employed. As a result, the Coulomb replica-exchange method sampled all local-minimum free-energy states more frequently than the other methods for both force fields. Moreover, the Coulomb, van der Waals, and usual replica-exchange methods were applied to a fragment of an amyloid- $\beta$  peptide (A $\beta$ ) in explicit water solvent to compare the sampling efficiency of these methods for a larger system. The Coulomb replicaexchange method sampled structures of the A $\beta$  fragment more efficiently than the other methods. We obtained  $\beta$ -helix,  $\alpha$ -helix, 3<sub>10</sub>-helix,  $\beta$ -hairpin, and  $\beta$ -sheet structures as stable structures and revealed pathways of conformational transitions among these structures from a free-energy landscape, as shown in Figure 1.



**Figure 1.** Folding pathways of amyloid- $\beta$  peptide obtained from Coulomb replica-exchange MD simulations.

## 3. Replica Exchange Molecular Dynamics Simulation of Chitosan for Drug Delivery System Based on Carbon Nanotube

Chitosan is an important biopolymer in the medical applications because of its excellent biocompatibility. It has been recently highlighted in the targeted drug delivery system (DDS) by improvement of the carbon nanotube (CNT) solubility. To investigate the effect of chitosan length, the two targeted DDSs with 30 and 60 chitosan monomers were performed by replica-exchange molecular dynamics simulations at temperatures in the range of 300-455 K. Each DDS model contains the epidermal growth factor (EGF), chitosan (CS) of 30 (30CS) and 60 (60CS) monomers, single-wall CNT (SWCNT) and gemcitabine (Gemzar) as the model payload anticancer drug, called EGF/30CS/ SWCNT/Gemzar and EGF/60CS/SWCNT/Gemzar, respectively. The SWCNT confines gemcitabine inside its cavity, while the outer surface is wrapped by chitosan in which one end is linked to the EGF. The results showed that in the EGF/30CS/SWCNT/Gemzar DDS the 30CS chain was not long enough to wrap around the SWCNT, and consequently the EGF was located so close to the tube as to potentially cause steric inhibition of the binding of EGF to its receptor (EGFR), which is highly expressed on the surface of cancer cells. On the other hand, this phenomenon is not observed in the EGF/60CS/SWCNT/Gemzar DDS in which the 60CS was found to completely wrap over the CNT outer surface using only 50 chitosan units. Although an increase in the temperature is likely to influence the overall DDS structure, and especially the orbit of helical chitosan on the SWCNT and the EGF conformation, gemcitabine is still encapsulated inside the tube.

### 4. Monte Carlo Simulation for Isotropic– Nematic Phase Transition of Infinitely Thin Liquid Crystal Molecules

We are also interested in isotropic-nematic phase transition of liquid crystal molecules. We gave a criterion to test a non-biaxial behavior of infinitely thin hard platelets based upon the components of three order parameter tensors. We investigated the nematic behavior of monodisperse infinitely thin rectangular hard platelet systems by this criterion. Starting with a square platelet system, and we compared it with rectangular platelet systems of various aspect ratios. For each system, we performed equilibration runs by using isobaric Monte Carlo simulations. Each system did not show a biaxial nematic behavior but a uniaxial nematic one, despite of the shape anisotropy of those platelets. The relationship between effective diameters by simulations and theoretical effective diameters of these systems was also determined.

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# Ultimate Quantum Measurements for Quantum Dynamics

# Department of Theoretical and Computational Molecular Science Division of Theoretical and Computational Molecular Science



SHIKANO, Yutaka NAKANE, Junko

Research Associate Professor (Feb., 2012–) Secretary

Due to great development on experimental technologies, it is possible to capture quantum dynamics in some physical and chemical systems. On the other hand, all experiments are in principle open and dissipative systems. Up to now, the wellexplained experiments are approximated to the equilibrium situation. However, by recent technological development, some experiments reach to a transition from equilibrium to non-equilibrium situations. While there are the well-known tools on the non-equilibrium situations; the linear response theory and the Keldysh Green function method, this analysis cannot basically catch dynamical situations. Our goal is to construct the time-resolved theoretical models included the non-equilibrium situations. However, the quantum measurement theory is needed on measuring quantum dynamics, especially considering the measurement backaction. Our current activities are to resolve how sensitive (quantum) measurement can we carry out in principle, to build up some toy models on quantum dynamic, and to explain photoluminescence phenomenon in nitrogen vacancy center in diamond.

# 1. Quantum Measurement Sensitivity without Squeezing Technique<sup>1)</sup>

As alluded before, our aim is to capture quantum dynamical phenomena. Capturing some phenomena needs to carry out the measurement. The conventional quantum measurement technique has huge measurement backaction. The measurement backaction prevents us chasing quantum dynamics like the classical trajectory. On the other hand, reducing the measurement backaction needs the tiny coupling between the target and probe quantum systems. However, under this situation, the signal in the probe system is also tiny small, that is, it is difficult to capture information. To resolve this problem, the squeezing technique was proposed and was experimentally implemented. However, this technique is practically difficult to be implemented. Our proposal is to use the weak measurement initiated by Aharonov, Albert, and Vaidman without squeezing technique. The profound meaning and interpretation of the weak measurement is seen in the review paper.<sup>2)</sup> The key of this method is to take the post-selection of the target system. Due to this effect, tiny probe signal can be amplified. In the original proposal by Aharonov, Albert, and Vaidman, the amplification factor is infinite by the approximation method. However, the effect measurement backaction is simultaneously amplified. When the probe state is Gaussian to be used in the original proposal, we have analytically shown the upper bound of the amplification factor. We have analytically derived the probe state to maximally amplify the signal by the variational method.<sup>1)</sup> By this optimal probe state, the amplification factor has no upper bound. Our result tells us the infinitely amplified single under the known coupling between the target and the probe. However, this result is ignored to consider the noise. We will extend our result to the signal-to-noise ratio optimization. Also, we are collaborating with the experiments to amplify the signal and consider the applications; the gravitational wave detector.

# 2. Discrete Time Quantum Walk as Quantum Dynamical Simulator<sup>3)</sup>

The discrete time quantum walk is defined as a quantum mechanical analogue of the classical random walk but is not the quantization of the classical random walk. This mathematical description is very simple but leads to many quantum dynamical phenomena. This is a toy model to better understand the quantum dynamics. Also, this has recently been various experimental demonstrations in the ultracold atoms in the optical lattice, trapped ions, and optical systems. We have analytically shown that the one- and two-dimensional discrete time quantum walks can be taken as the quantum dynamical simulator,<sup>3)</sup> which concept is to emulate some classes of the differential equations, for example, the Dirac equation. Our approximation is used from the discrete lattice to continuous line for the large time steps of the discrete time quantum walk. This mathematical treatment is so powerful like the relation-ship between the cellular automaton and the integrable system.

# 3. Photoluminescence Phenomenon from Solid-State System<sup>4)</sup>

A method to measure some physical properties by light is widely used in physical, chemical, and biological systems. Therefore, laser science has been developed along with our demands from science and technology. A single photon source is expected as the low power laser source and the quantum communication tool. A nitrogen vacancy center in diamond and a quantum dot in a semiconductor system are the promising candidate of the single photon source. Especially, a nitrogen vacancy center in diamond has been attracted since this is run at room temperature. For an application as the highly controlled photon source, the photoluminescence process in the nitrogen vacancy center in diamond needs to be well understood. This system has the S = 1 electronic spin with the hyperfine structure 2.87 GHz. Furthermore, inserting a magnetic field, the different magnetic states are separated due to the Zeeman shift. Since this electronic spin is highly

localized, the local magnetic field evaluation is needed. We have shown the method to evaluate the local magnetic field from the conventional confocal microscopy.<sup>4)</sup> Also, the nitrogen vacancy center in diamond is a candidate of the quantum memory. Since a lifetime of the nuclear spin of a <sup>13</sup>C atom nearly located in the nitrogen vacancy spot is long (~ sec order), the perfect quantum state transfer is needed. However, it is difficult to evaluate the location of the <sup>13</sup>C atom. On the other word, it is difficult to evaluate the coupling constant between the electronic spin of the nitrogen vacancy center and the nuclear spin of the <sup>13</sup>C atom. We have proposed the simple scheme to evaluate the coupling constant between the two spins system under the dissipative situation.<sup>5)</sup> As the next step, we will study the photoluminescence process from the nitrogen vacancy center in diamond.

As the current activities of our laboratory, we are studying the photoluminescence processes of the quantum dots and the exciton-polariton Bose-Einstein condensations in the two dimensional electronic gas of the semiconductor. These materials are expected to be used as the classical optical devices; the optical switching, collaborated with the various experimentalists.

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# Theoretical Studies of Chemical Dynamics in Condensed and Biomolecular Systems

Department of Theoretical and Computational Molecular Science Division of Theoretical and Computational Molecular Science



ISHIZAKI, Akihito YAMADA, Mariko Research Associate Professor (March, 2012–) Secretary

Photosynthesis provides the energy source for essentially all living things on Earth, and its functionality has been one of the most fascinating mysteries of life. Photosynthetic conversion of the energy of sunlight into its chemical form suitable for cellular processes involves a variety of physicochemical mechanisms. The conversion starts with the absorption of a photon of sunlight by one of the light-harvesting pigments, followed by transfer of electronic excitation energy to the reaction center, where charge separation is initiated. At low light intensities, surprisingly, the quantum efficiency of the transfer is near unity. A longstanding question in photosynthesis has been the following: How does light harvesting deliver such high efficiency in the presence of disordered and fluctuating dissipative environments? Why does not energy get lost? At high light intensities, on the other hand, the reaction center is protected by regulation mechanisms that lead to quenching of excess excitation energy in light harvesting proteins. The precise mechanisms of these initial steps of photosynthesis are not yet fully elucidated from the standpoint of molecular science.

# 1. Electronic Excitation Transfer Dynamics in Light Harvesting Systems

Recently, the technique of two-dimensional electronic spectroscopy has been applied to explore photosynthetic light harvesting complexes. The observation of long-lived quantum superposition or coherence between eigenstates of electronic excitations (excitons) in a pigment-protein complex stimulated a huge burst of activity among experimentalists and theorists. Much of the interest arose because the finding of electronic quantum coherence is a "warm, wet, and noisy" biological system was considered very surprising. The initial experiments were carried out at 77 K, but more recent studies have detected coherence lasting at least 300 fs at physiological temperatures. These observations have led to the suggestion that electronic quantum coherence might play a role in achieving the remarkable efficiency of photosynthetic light harvesting. At the same time, the observations have raised questions regarding the role of the surrounding protein in protecting the coherence. In order to elucidate origins of the long-lived electronic coherence and its interplay with the protein environment, we have investigated appropriate theoretical frameworks and concepts with the use of statistical mechanics<sup>1)</sup> and in cooperation with experiment.<sup>2)</sup>

In particular, we have discussed significance of the finite timescale site-dependent reorganization dynamics of the protein environment surrounding (bacterio)chlorophyll molecules. The electronic coupling  $\hbar J$  between pigments and the excitation-environment coupling characterized by the reorganization energy  $\hbar\lambda$  are two fundamental interaction mechanisms determining the nature of electronic energy transfer (EET) in photosynthetic complexes. Ordinarily, photosynthetic EET is discussed only in terms of the mutual relation between magnitudes of the two couplings, as just described. However, the nature of EET also depends on the mutual relation between two timescales, the characteristic timescale of the environmental reorganization,  $\tau$  (Figure 1), and the inverse of the electronic coupling,  $J^{-1}$ , that is the time the excitation needs to move from one pigment to another neglecting any additional perturbations. We showed that even in the region of reorganization energy much larger than that of a typical situation of photosynthetic light harvesting the sluggish reorganization dynamics allows the excitation to stay above an energy barrier separating two local minima, which correspond to the two sites in the adiabatic potential surface, for a prolonged time (Figure 2). This is an essential origin of the experimentally observed long-lived electronic coherence.

Further, we showed this energy barrier is much small in comparison with the thermal energy  $k_{\rm B}T$  (almost activationless) in the parameter regime corresponding to natural light harvesting systems. It can be considered that natural photosynthetic light harvesting systems avoid local trapping of electronic excitations, which yields a situation that electronic excitation energy gets lost in in the presence of disordered and fluctuating dissipative environments. By extending this argument, we are now tackling the above questions regarding the remarkable quantum efficiency and smart regulation mechanism that photosynthetic light harvesting systems exhibit.



**Figure 1.** Schematic illustration of the *m*th pigment embedded in a protein (*a*) and the electronic ground and excited states of the *m*th pigment,  $|\varphi_{mg}\rangle$  and  $|\varphi_{me}\rangle$ , affected by nuclear motion of the protein environment (*b*). After electronic excitation in accordance with the vertical Franck-Condon transition, reorganization takes place from the equilibrium nuclear configuration with respect to the electronic ground state  $|\varphi_{mg}\rangle$  to the actual equilibrium configuration in the excited state  $|\varphi_{me}\rangle$  with dissipation of the reorganization energy,  $\hbar\lambda_m$ . This reorganization of the protein environment proceeds on a characteristic timescale,  $\tau_m$ .



associated with pigment 1 Figure 2. Schematic illustration of adiabatic excitonic potential surface of coupled two pigments. The point of origin corresponds to the Franck-Condon state. The energy of the point is higher than the barrier between the minima; therefore, we find that the electronic excited state is delocalized just after the excitation. As time increases,

the dissipation of reorganization energy proceeds and the excitation will fall off into one of the minima and become localized. Namely, sluggish dissipation of reorganization energy increases the time an electronic excitation stays above an energy barrier separating pigments and thus prolongs delocalization over the pigments.

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#### Award

ISHIZAKI, Akihito; Short-term Fellowship at Wissenschaftskolleg zu Berlin, 2012/2013.

# Theory and Computation of Reactions and Properties in Solutions and Liquids

Department of Theoretical and Computational Molecular Science Division of Computational Molecular Science



ISHIDA, Tateki

Assistant Professor

We focus on the projects both on ultrafast photoinduced electron energy transfer in the excited state in solution and on ionic liquids (ILs). The project on photoinduced electron energy transfer processes in the excited state in solution focuses on the development of a theoretical method to describe electron energy transfer including solvent motion and dynamics with the theoretical treatment we have developed. On the other hand, ILs' projects concentrate the study of specific interionic dynamics in ILs and the investigation of a new perspective on the physically and chemically unique characteristics of ILs.

### 1. The Theoretical Study of Photoinduced Electron Energy Transfer Processes in the Excited State in Solution

We have developed a procedure for describing the timedependent evolution of the electronic structure of a solute molecule in solution, coupling an electronic structure theory with solvent motion in the formalism of an interaction site model. We extend this prescription for studying electron energy transfer processes in the excited state in solution. It is indicated that the coupling between solvation dynamics and a fast intramolecular electron energy transfer is likely to play an important role in the emergence of photoinduced unique functionalities in biochemical and metal complex systems.

### 2. The Unique Physical and Chemical Properties of Ionic Liquids through Interionic Interactions: Theoretical Investigation with Molecular Dynamics Simulations<sup>1)</sup>

Ionic liquids (ILs) have been found to possess a wide potential variety of interesting physical and chemical properties. We consider the unique properties of ILs which owe to the specific interionic interaction between ionic species. In particular, we study the importance of both the cross-correlation between cation and anion species, and the effect of polarization in ILs. On the basis of recent theoretical studies on ILs employing molecular dynamics simulations, how the collective dynamics through interionic interactions cause the unique physical and chemical properties of ILs and how electronic polarizability effects modify interionic dynamics have been investigated. Those include the investigation of the contribution of ionic motions with velocity cross-correlation functions and the study of many-body polarization effects on the cage effect in ILs.

# 3. Investigations of New Perspectives on the Characteristics of Ionic Liquids: Microscopic Aspects in Dicationic Ionic Liquids<sup>2)</sup>

The interionic vibrations in imidazolium-based dicationic ionic liquids (ILs) containing the bis(trifluoromethylsulfonyl) -amide ([NTf2]-) counteranion were investigated using femtosecond optical-heterodyne-detected Raman-induced Kerr effect spectroscopy. The microscopic nature of the dicationic ILs  $([C_n(MIm)_2][NTf_2]_2$ , where n = 6, 10, and 12; MIm = Nmethylimidazolium) was compared with that of the corresponding monocationic ILs ( $[C_nMIm][NTf_2]$ , where n = 3, 5, and 6) used as reference samples. Low-frequency Kerr spectra within the frequency range 0-200 cm<sup>-1</sup> of the ILs revealed that the spectral profile of the dicationic ILs as well as that of the corresponding monocationic ILs is bimodal. The distinguished line-shape of the low-frequency Kerr spectrum of [C<sub>3</sub>MIm] [NTf<sub>2</sub>] from the other ILs can be accounted for by the homogeneous nature in the microstructure of the IL, but the other ILs indicate microsegregation structures due to the longer nonpolar alkylene linker or alkyl group in the cations.

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# **Visiting Professors**



## Visiting Associate Professor HASEGAWA, Jun-ya (from Kyoto University)

Quantum Chemistry for the Excited States of Functional Molecules in Proteins and Solutions Molecular interactions between chromophore and environment are the essential to furnish a protein with the photo-functionality. I am interested in the machinery of the photo-functions such as photosynthesis, vision, and bioluminescence. To understand the mechanism and to develop chemical concept behind the photo-functions, we develop electronic structure theories for excited state, analytical method for

excitation-energy transfer pathway, and a hybrid quantum-mechanics/molecular mechanics method. In recent studies, we have clarified color-tuning mechanism of photo-functional proteins and excitation transfer mechanism of bridge-mediated donor-acceptor systems. We are also interested in developing a configuration interaction picture for the solvatochromic response of the molecular environment.



Visiting Associate Professor ANDO, Koji (from Kyoto University)

### Quantum Transfer Processes in Chemical and Biological Systems

At the core of chemistry and biochemistry are reduction-oxidation and acid-base reactions, whose elementary processes are electron and proton transfers. Our research group has been working on theoretical and computational modeling of these inherently quantum dynamical processes. One recent achievement is a development of electron transfer (ET) pathway analysis method with use of fragment molecular orbital

calculations, by which the ETs in a bacterial photosynthetic reaction center have been studied. Another is a development of nuclear wave packet molecular dynamics simulation method and its applications to hydrogen-bond exchange dynamics in water. It has been also discovered that the latter can be extended to electron wave packet simulations by exploiting the non-orthogonal valence-bond theory, which anticipates a non-Born-Oppenheimer electron–nuclear dynamical treatment.



### Visiting Associate Professor **MORISHITA, Tetsuya** (from AIST)

#### First-Principles Molecular-Dynamics Simulations of Liquids and Glassy Materials

I have been interested in structural and dynamical properties of non-crystalline materials including nanoscale materials. In 2011, I have found that liquid silicon exhibits compressed exponential relaxation over a wide temperature range including the supercooled regime, in contrast to water and other glass-forming liquids, using first-principles molecular-dynamics simulations.

I have also developed a new method for free-energy calculation, Logarithmic Mean-Force Dynamics (LogMFD). This method was successfully applied to reconstruction of the free-energy profile of a dipeptide molecule, showing that LogMFD considerably outperforms conventional free-energy calculation methods such as thermodynamics integration.

# **RESEARCH ACTIVITIES**

# **RESEARCH ACTIVITIES** Photo-Molecular Science

We study the interaction between molecules and optical fields with its possible applications to active control of molecular functionality and reactivity. We also develop novel light sources to promote those studies. Two research facilities, the Laser Research Center for Molecular Science and the UVSOR, closely collaborate with the Department.

The core topics of the Department include ultrahigh-precision coherent control of gas- and condensed-phase molecules, high-resolution optical microscopy applied to nanomaterials, synchrotron-based spectroscopy of core-excited molecules and solid-state materials, vacuum-UV photochemistry, and the development of novel laser- and synchrotron-radiation sources.

# Development of Advanced Near-Field Spectroscopy and Application to Nanometric Systems

Department of Photo-Molecular Science Division of Photo-Molecular Science I



OKAMOTO, Hiromi NARUSHIMA, Tetsuya NISHIYAMA, Yoshio KOWAKA, Yasuyuki WU, Huijun

HASHIYADA, Shun OCHIAI, Takao ISHIKAWA, Akiko NOMURA, Emiko YAMASAKI, Yumi Professor Assistant Professor IMS Research Assistant Professor Post-Doctoral Fellow Graduate Student ( -March, 2012) Technical Fellow (April, 2012- ) Graduate Student Graduate Student\* Technical Fellow Secretary Secretary

There is much demand for the study of local optical properties of molecular assemblies and materials, to understand nanoscale physical and chemical phenomena and/or to construct nanoscale optoelectronic devices. Scanning nearfield optical microscopy (SNOM) is an imaging method that enables spatial resolution beyond the diffraction limit of light. Combination of this technique with various advanced spectroscopic methods may offer a direct probe for dynamics in nanomaterials and nanoscale functionalities. It may provide essential and basic knowledge to analyze origins of characteristic features of the nanomaterial systems. We have constructed apparatuses of near-field spectroscopy and microscopy for excited-state studies of nanomaterials, with the feasibilities of nonlinear and time-resolved measurements. The developed apparatuses enable near-field measurements of twophoton induced emission and femtosecond time-resolved signals, in addition to conventional transmission, emission, and Raman scattering. Based on these methods, we are investigating the characteristic spatiotemporal behaviors of various metal-nanoparticle systems and molecular assemblies.

### 1. Visualization of Localized Optical Fields and Plasmon Wavefunctions in Metal Nanostructures

We recently reported that wavefunctions of localized plasmon resonances of chemically synthesized metal (Au and Ag) nanoparticles are visualized by near-field transmission or two-photon excitation measurements.<sup>1)</sup> The same methods were also applied to Au nanoparticle assemblies to visualize confined optical fields.<sup>1)</sup> It was revealed for the Au nanoparticle dimers that highly localized optical field is generated at the interstitial sites between the particles. In many-particle assemblies, the localized fields were especially intensified at the rim parts of the assemblies, and such a characteristic field distribution has been attributed to interaction between plasmon excitations induced on the particles. We also observed confined optical fields in gaps between circular apertures opened on thin gold films.

We are extending the studies to various two-dimensional metal nanostructures manufactured by the electron-beam lithography technique, in part as collaboration with researchers of other institution, with which structures that are difficult to obtain with the chemical methods can be available. The nanostructures now under study are circular Au nanodisks, assembled Au nanodisks with designed arrangements, elongated rectangular nanoapertures opened on Au thin films, and so forth. Such a study is essential as a basis for designing unique optical properties and functions of metal nanostructures.

### 2. Studies of Metal-Nanostructure Modified Photovoltaic Cells by Near-Field Excited Site-Specific Photocurrent Detection

Metal nanoparticles and their assemblies collect photon energies to yield confined and enhanced optical fields in the vicinities of the particles due to plasmon resonances. Recently, it has been reported that efficiencies of photoenergy conversion can be improved by the use of noble metal nanostructures. The photoenergy conversion system ranges from wet-type and solid-state photo-current conversion cells to photo-chemical conversion systems. To reveal the mechanism of the enhanced photoenergy conversion process, studies of detailed nanostructures and site-dependent photocurrent measurements are essential.

We applied SNOM to clarify effects of surface plasmon resonance on photo-current conversion in GaAs semiconductor cells modified with Au nanoparticles, by photocurrent imaging measurements with localized near-field photoirradiation (Figure 1).<sup>2)</sup> Isolated nanospheres caused local photocurrent sup-



**Figure 1.** Near-field photocurrent images for GaAs photodiode modified with Au spherical nanoparticles. (a) Phorocurrent image at 532 nm for isolated particles (top) and line profiles along the dashed lines in the image (bottom). (b) Photocurrent images at 785 nm for dimeric particles (top; dimers are indicated with white circles; vertical and horizontal polarizations for left and right images, respectively) and line profiles along the white lines in the image for respective polarization directions (bottom).

pressions at the plasmon resonance wavelengths (~530 nm). In assemblies (dimers and trimers) of the spheres, a remarkable decrease of photocurrent at the gap site between the spheres was observed at ~800 nm near the dimer plasmon resonance, despite anticipated field enhancements in the gap sites. From the results, it is concluded that the enhanced optical fields induced by the plasmons do not improve the photovoltaic efficiency. The far-field forward scattering of photons by the gold nanoparticles may be more important than the enhanced field effects for the GaAs photovoltaic device studied.

### 3. Construction of Apparatuses for Sub-20-fs Ultrafast Near-Field Spectroscopy

Surface plasmons of noble metal nanoparticles have very short lifetimes in the range of ~2–20 fs. To achieve a SNOM observation of such fast dynamics in the individual nanoparticles, we have to overcome serious dispersion effects arising from the optical components involved, especially from the optical fiber. We achieved that by combining the conventional dispersion compensation devices composed of prism and grating pairs with adaptive pulse shaping technique, and succeeded in delivering 17-fs pulses to the the near-field aperture with a spatial resolution of ~100 nm.<sup>3</sup>)

The layout of the experimental setup is shown in Figure 2. A pair of transmission gratings and an SF14 prism pair were adopted to precompensate for the second- and third-order group velocity dispersion effects. The beam was then sent to adaptive pulse shaping system equipped with a deformable mirror, whose surface shape was adjusted to yield a shortest pulse duration at the probe tip using a genetic algorithm. The



Figure 2. Experimental setup of sub-20-fs ultrafast near-field measurement.

laser beam was then coupled to a 150-mm-long optical fiber to introduce the light into the apertured near-field probe tip. We used time-correlated two-photon-induced photoluminescence (TPI-PL) measurements to demonstrate the performance of the system. The light pulses were incident on a gold nanostructure sample through the near-field probe to excite the TPI-PL from Au, whose intensity was recorded as a function of the pump– probe delay time. With this setup we succeeded in measuring dephasing of ~10 fs in Au nanostructures. Position dependent dephasing measurements for Au nanostructures are now under way.

### 4. Near-Field Circular Dichroism Microscopy of Nanomaterials

Circular dichroism (CD) spectroscopy is widely used in the studies of chiral materials and magnetism. Some of nanomaterials composed of achiral molecules are reported to show CD activities arising from the nanoscale chirality. Two-dimensionally chiral metal nanostructures also show CD activities. Investigation of nanoscale local CD may provide valuable information on the origins of CD activities of such materials. For this purpose, we are developing an apparatus for near-field CD microscopy and measuring CD images of nanoscale chiral materials.

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# Design and Reconstruction of Molecular Quantum States of Motion

# Department of Photo-Molecular Science Division of Photo-Molecular Science I



OHSHIMA, Yasuhiro MIZUSE, Kenta FUJIWARA, Masanori HAYASHI, Masato MIYAKE, Shinichiro INAGAKI, Itsuko Professor Assistant Professor IMS Research Assistant Professor Post-Doctoral Fellow Graduate Student Secretary

Molecules are vital existence. In a gas-phase ensemble at room temperature, they are, in an average, flying away by a few hundred meters, making turns almost reaching to  $10^{11}$ times, and shaking themselves more than  $10^{13}$  times within the duration of only one second. The ultimate goal this research group has been aiming to is to capture the lively figures of molecules moving in such a dynamic manner and to have a perfect command over the molecular motions. Here lasers with ultimate resolution in time and energy domains are employed complementally and cooperatively for this purpose.

### 1. Nonadiabatic Excitation of Molecular Rotation Induced by Intense Ultrashort Laser Fields

When a gaseous molecular sample is irradiated by an intense nonresonant ultrashort laser pulse, the laser field exerts a torque that aligns the molecular axis along the laser polarization vector, due to the interaction with the molecular anisotropic polarizability. Here the field-matter interaction only remains in much shorter duration than the characteristic time for molecular rotation, and thus the rotation of the molecules is coherently excited to create a rotational quantum wave packet (WP). We have developed a method to explore the nonadiabatic excitation in a quantum-state resolved manner and applied it to diatomic and symmetric-top molecules.<sup>1)</sup> It has been shown that the state distribution is a useful experimental source for verifying the excitation process.<sup>2,3)</sup> When a pair of excitation pulses is implemented with appropriate time delay between them, partial control of rotational-state distribution has been achieved.<sup>1,4)</sup> In a favorable case, the double-pulse excitation coupled with the state-selective probe has enabled us to reconstruct experimentally a rotational WP thus created.<sup>5)</sup> If the mutual polarization direction and time delay between the two pulses are adjusted, the sense of rotation around the laser propagation direction can also be controlled, yielding to a rotational WP exhibiting angular-momentum orientation.<sup>6,7)</sup>

# 2. Coherent Excitation of Intermolecular Vibrations in Molecular Clusters by Nonresonant Intense Ultrashort Laser Fields

Nonadiabatic interaction with a nonresonant intense ultrashort laser field can also coherently excite vibration of molecule through the structural dependence of the molecular polarizability. We have recently succeeded in creating and observing WPs pertinent to intermolecular vibrations of several molecular clusters in their electronic ground states. So far, NO-Ar, benzene dimer and trimer have been studied. Here, vibrational distribution after nonadiabatic vibrational excitation (NAVEX) is probed in a quantum-state resolved manner. By monitoring the resonant two-photon ionization (R2PI) spectra, substantial decrease of the transition from the vibrationally ground state has been clearly observed when the femtosecond (fs) pump pulse was applied, and in the case of NO-Ar, the observation was accompanied with the emergence of several hot bands, assigned to those from vibrationally excited states pertinent to intermolecular modes. The double fs pulse excitation has also been implemented for examining the real-time quantum interference of the WPs. The observed time-domain signals for NO-Ar were directly compared with the calculation on the WP propagation by numerically solving the time-dependent Schrödinger equation on the intermolecular potential energy surface. For more detailed examination of the NAVEX process and further application to various molecular cluster systems, we are now renovating the experimental setup, in particular, the molecular beam time-of-flight mass spectrometry (TOF-MS) apparatus. The modification includes: Improvement of mass resolution by extending the TOF flight tube, rejection of the interference of strong monomer signals by implementing a mass gate, and realization of more extensive cooling by enlarging the vacuum pumping capabilities.

## 3. High-Resolution Laser Spectroscopy of Benzene Clusters with Atoms and Small Molecules

Molecular clusters containing benzene are prototypical systems for elucidating the intermolecular interaction pertinent to aromatic rings. The information on the precise cluster geometry and energy-level structure pertinent to the intermolecular vibration are useful experimental input to reconstruct the intermolecular potential energy surface. We are now focusing on clusters of benzene attached by small numbers of atoms and molecules. So far, electronic spectra of C<sub>6</sub>H<sub>6</sub> complexed with one and two He atom(s), two Ar atoms, up to three H<sub>2</sub> molecules, and one H<sub>2</sub>O molecule have been examined via two-color R2PI in the vicinity of the monomer  $S_1 - S_0 6_0^{-1}$ band. We employed a tripled output from a ns pulsed dye amplifier, which was injection-seeded by the CW output from a Ti:Sapphire laser, as an excitation source. Owing to the narrow band width (~250 MHz) of the laser system and the efficient rotational cooling down to 0.3 K by implementing a high-pressure pulsed valve, rotational structures have been greatly simplified.

Structural parameters of the clusters with He have been substantially refined than those reported previously.<sup>8)</sup> The distances of He above the plane are set to be: 3.602 (+0.063) Å and 3.596 (+0.057) Å, respectively, for the clusters with one and two He atoms, where values in parentheses represent the change by the excitation from  $S_0$  to  $S_1$ . Several vibronic bands with excitation of intermolecular vibrations have also been observed. The vibrational frequencies of  $benzene(S_1)$ -He are derived as: 17 and 13 cm<sup>-1</sup> for the intermolecular stretch and bend modes, respectively. The intermolecular stretch of the cluster with two He is 15.72 cm<sup>-1</sup>, while no experimental information has been obtained for the bend modes. The vibronic bands of benzene-He exhibit tunneling splitting due to a largeamplitude migration of He above and below the benzene molecular plane. This finding is matched with the prediction based on a high-level ab initio calculation. We have also recorded the R2PI spectrum of mono <sup>13</sup>C substituted species of benzene-He.

For the clusters with H<sub>2</sub>, two distinguished isomers, correlating to para and ortho H<sub>2</sub>, are identified. This finding shows that the permutation of the two H atoms in H<sub>2</sub> is feasible in the clusters. When the normal H<sub>2</sub> was used as a sample gas, only the clusters of ortho H2 were observed, and thus we had to use a gas sample of pure para H<sub>2</sub> to record the clusters with para H<sub>2</sub> (see Figure 1). This observation indicates the smaller effective binding energy for benzene-para H<sub>2</sub> than that of the ortho H<sub>2</sub> cluster. This energy relation has been well known for the clusters with smaller molecules (e.g., HF, HCl, and OCS) attached by H<sub>2</sub>. It is noted the present study is the first experimental report on the coexistence of para and ortho H<sub>2</sub> isomers in the clusters of aromatic molecules. Rotationally resolved spectra allowed us to fix the cluster geometry unambiguously, as shown in Figure 2. It has shown that the effective intermolecular distances are substantially different from each other between the clusters with para and ortho H<sub>2</sub>, indicating the change in the average H<sub>2</sub> orientation relative to the benzene plane in the two isomers.







**Figure 2.** High-resolution excitation spectra and the experimentally derived cluster geometry of  $C_6H_6$ -(ortho  $H_2)_n$ .

We also recorded for the first time rotationally resolved excitation spectrum of the cluster with two Ar atoms lying on the same side of the benzene plane.

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# Development of High-Precision Coherent Control and Its Applications

# Department of Photo-Molecular Science Division of Photo-Molecular Science II



OHMORI, Kenji KATSUKI, Hiroyuki TAKEI, Nobuyuki GOTO, Haruka SOMMER, Christian NAKAGAWA, Yoshihiro KOYASU, Kuniaki INAGAKI, Itsuko YAMAGAMI, Yukiko Professor Assistant Professor Assistant Professor Post-Doctoral Fellow Post-Doctoral Fellow Graduate Student Graduate Student Secretary Secretary

Coherent control is based on manipulation of quantum phases of wave functions. It is a basic scheme of controlling a variety of quantum systems from simple atoms to nanostructures with possible applications to novel quantum technologies such as bond-selective chemistry and quantum computation. Coherent control is thus currently one of the principal subjects of various fields of science and technology such as atomic and molecular physics, solid-state physics, quantum electronics, and information science and technology. One promising strategy to carry out coherent control is to use coherent light to modulate a matter wave with its optical phase. We have so far developed a high-precision wave-packet interferometry by stabilizing the relative quantum phase of the two molecular wave packets generated by a pair of femtosecond laser pulses on the attosecond time scale. We will apply our high-precision quantum interferometry to gas, liquid, solid, and surface systems to explore and control various quantum phenomena.

#### Award

OHMORI, Kenji; Humboldt Research Award (Germany).

# Molecular Inner-Shell Spectroscopy: Local Electronic Structure and Intermolecular Interaction

# Department of Photo-Molecular Science Division of Photo-Molecular Science III



KOSUGI, Nobuhiro YAMANE, Hiroyuki NAGASAKA, Masanari YUZAWA, Hayato MOCHIZUKI, Kenji LELOUP, Valentin NAKANE, Junko Professor Assistant Professor Assistant Professor IMS Fellow Graduate Student Graduate Student\* Secretary

In order to reveal local electronic structures and weak intermolecular interactions in molecular systems such as organic solids, liquids, aqueous solutions, and molecular clusters, we are developing and improving several kinds of soft X-ray spectroscopic techniques, such as X-ray photoelectron spectroscopy (X-ray PES, XPS), X-ray absorption spectroscopy (XAS), resonant Auger electron spectroscopy (RAS), X-ray emission spectroscopy (XES), resonant XES (RXES), and resonant inelastic X-ray scattering (RIXS), at UVSOR in-vacuum undulator beamlines BL-3U and BL-6U with some international collaboration programs, and also an original *ab initio* quantum chemical program package GSCF, which is optimized to calculation of molecular inner-shell processes.

# 1. Concentration Dependence of Local Structure in Methanol-Water Binary Liquid

A water molecule has two H(hydrogen)-accepting ('acceptor') and H-donating ('donor') sites, and the water liquid forms three dimensional (3D) hydrogen bonding (HB) networks. On the other hand, a methanol molecule with a hydrophobic methyl group has one H-donor and one or two H-acceptor sites, and the methanol liquid forms one and two dimensional (1D, 2D) HB networks, such as chains and rings of 6–8 methanol molecules. It is known that the methanol-water binary liquid has 3D cluster structures. However, the local structures of the methanol-water mixtures are still unknown. In this work, the local structure of the methanol-water binary liquid at different concentrations is studied by using carbon K-edge XAS.

The experiments were performed at soft X-ray undulator

beam line BL3U at UVSOR-II facility. The liquid thin layer was sandwiched between two 100 nm-thick  $Si_3N_4$  membranes. The photon energy was calibrated by the first peak (287.96 eV) of methanol gas filled in a He buffer space between the  $Si_3N_4$  membrane and the detector in the liquid cell.

Figure 1 shows the carbon K-edge XAS spectra of the methanol–water binary liquids. The peak around 288.5 eV arises from the molecular orbital with a component of the hydroxyl group in the methanol molecule, and has a HB effect. On the other hand, the peak around 290 eV is mainly distributed at the methyl group. This peak is shifted to the higher photon energy by increasing the mixing ratio of water. It means that the methyl groups approach with each other in a dilute aqueous solution and the interaction between the methyl groups are apart from each other.



**Figure 1.** Carbon K-edge XAS spectra of the methanol-water binary liquid at different concentrations at 25 °C. The mixing ratio of water is increased along indicated arrows.

Figure 2 shows the intensities of the 289 eV region at different mixing ratios. The intensities are changed nonlinearly, and indicate three different regions. The intensities are not changed so much in the methanol-rich region (I), indicating the local structure of pure methanol chain and ring clusters is preserved. On the other hand, in the region (II), the intensity is suddenly decreased at the methanol molar fraction  $X \approx 0.7$ , indicating a 1D/2D methanol–water mixed cluster is suddenly formed and this local structure is preserved down to  $X \approx 0.3$ . When  $X < \sim 0.3$  (III), the intensity is going down, indicating a 3D methanol–water mixed cluster is almost proportionally formed.



**Figure 2.** Intensities of the 289 eV region as a function of molar fraction of methanol in the methanol–water mixtures. The intensities of methanol and water are normalized to one and zero, respectively

### 2. Intramolecular and Intermolecular Interactions in Crystalline Films of Zinc Phthalocyanine (ZnPc)

It has been theoretically predicted that local and nonlocal electron-phonon (e-ph) interactions play a crucial role in the charge transport mechanism in organic solids. Recently, we have succeeded in observation of a small energy *versus* wave-vector  $[E(\mathbf{k})]$  relation for the crystalline film of an archetypal organic semiconductor of ZnPc on Au(111) by using the high-resolution angle-resolved UV photoemission spectroscopy (ARPES). Based on the precise  $E(\mathbf{k})$  data, we have evaluated the e-ph interaction on the charge transport mechanism of the ZnPc crystalline film.

Figure 3(a) shows the  $E(\mathbf{k})$  relation of the ZnPc crystalline film on Au(111), which is obtained from the photon energy dependence of the normal-emission ARPES spectra at 15 K. The highest occupied molecular orbital [HOMO ( $\pi$ ,  $a_{1u}$ )], Zn 3d ( $b_{1g}$ ), and  $\pi$ 's derived peaks show a dispersive shift with the same periodicity. This periodicity corresponds well with the highly symmetric  $\Gamma$  and Y points, estimated from the lattice constant of 3.32 Å determined by the specular X-ray diffraction [Figure 3(b)].

In order to examine the e-ph coupling in crystalline ZnPc film, we measured the temperature dependence of ARPES at the  $\Gamma$  and Y points with the 60 and 90 eV incident photons, respectively. As seen in Figure 3(c), the  $\Gamma$ -point HOMO energy shifts depending on the temperature, while the Y-point HOMO energy seems independent on the temperature. From the lineshape analysis, we found that both the HOMO bandwidth (*i.e.*, energy difference between the  $\Gamma$ -HOMO and Y-HOMO) and the  $\Gamma$ -HOMO peakwidth are narrowed with increasing the temperature, which is an opposite behavior to the remanent Au(111) Fermi edge. This is distinct indication of the local/nonlocal e-ph interaction; that is, band narrowing due to the interaction with local phonons (intramolecular vibrations) and band widening due to the interaction with nonlocal phonons (intermolecular lattice vibrations). Judging from the temperature dependence of the HOMO bandwidth and peakwidth, it is considered that the temperature-dependent transition between the incoherent hopping transport associated with the local e-ph coupling and the coherent band transport associated with the nonlocal e-ph coupling occurs within the energy scale of 100 meV in the ZnPc crystalline film.



**Figure 3.** (a) Intermolecular  $E(\mathbf{k})$  relation of the crystalline ZnPc film on Au(111) at 15 K, wherein the second derivative of the ARPES spectra were used for mapping out. White curves for the HOMO and Zn 3d derived bands are the best-fit results in the tight-binding model. (b) CuK $\alpha$  X-ray diffraction of the crystalline ZnPc film on Au(111) at 300 K. (c) Temperature dependence of ARPES in the HOMO-band region at the  $\Gamma$  and Y points, measured with 60 and 90 eV incident photons, respectively.

# Photoabsorption and Photoionization Studies of Fullerenes and Development of High-Efficiency Organic Solar Cells

# Department of Photo-Molecular Science Division of Photo-Molecular Science III



MITSUKE, Koichiro KATAYANAGI, Hideki PRAJONGTAT, Pongthep MORENOS, Lei Angeli S. VAILIKHIT, Veeramol BASHYAL, Deepak ASARI, Chie SHIMIZU, Atsuko Associate Professor ( –March, 2012)\* Assistant Professor Research Fellow Visiting Scientist Visiting Scientist Graduate Student Technical Fellow Secretary

We have observed the formation of multiply-charged photoions from gaseous fullerenes or aromatic hydrocarbons irradiated with synchrotron radiation at hv = 25 to 200 eV. We thus studied the mechanisms and kinetics of consecutive C<sub>2</sub>-release reactions on the basis of (i) the yield curves for the fragments  $C_{60(70)-2n}^{z+}$  ( $n \ge 1$ , z = 1-3) as a function of the primary internal energy and (ii) the three dimensional velocity distributions of the fragments. The velocity distributions of  $C_{60-2n}^{z+}$  and  $C_{70-2n}^{z+}$  were measured for the first time. Concepts of the microcanonical temperature and Arrhenius-type rate constants for individual C<sub>2</sub> ejection steps allowed us to compare the experimental total average kinetic energy with theoretical kinetic energy release predicted from the "model free approach" developed by Klots.

In the second topic, we have fabricated dye-sensitized solar cells (DSSCs) containing ruthenium dye, iodide electrolyte, and platinum or carbon nanotube catalysts. The incidence photon-to-current conversion efficiencies (IPCE) and photoabsorbance (ABS) were measured in the range of 300 nm to 1  $\mu$ m. We also evaluated the quantum yield (APCE) of DSSCs for the electron injection from the excited orbital of Ru dye to the conduction band of mesoporous TiO<sub>2</sub> nanoparticles. Our final goal is to develop DSSCs with high performance by improving ABS and APCE.

## 1. Photodissociation Dynamics of Fullerenes by Velocity Map Imaging with Improved Mass Resolution

We observed images of three-dimensional scattering velocity distributions of fragments produced by the photodissociation of the fullerenes,  $C_{60}$  and  $C_{70}$ , using a massresolved velocity map imaging spectrometer.<sup>1,2)</sup> From the images we successfully obtained translational temperatures of the fragments and kinetic energy releases. Still the peaks of the fragments were found to be partially overlapped in their time-of-flight (TOF) mass spectra. We tried to improve the

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mass resolution of the imaging spectrometer by adding a focusing electrode to the Eppink-Parker type imaging electrode assembly.<sup>3)</sup> This focusing electrode was designed after the ion lens device presented by Peši.

After introducing the focusing electrode, vertical stripes involving the fragment ions are completely resolved in the velocity vs. TOF maps, or *y*–*t* maps, and their masses are unquestionably identified. The *y*–*t* maps with improved mass resolution were obtained for C<sub>70</sub> with the excitation photon energy *h*v range of 70–125 eV. The translational temperature  $T_{\text{trans}}$  of each fragment was extracted from the map. The *h*v dependence of  $T_{\text{trans}}$  exhibits the same propensity as that obtained previously<sup>2)</sup> at low mass resolution. For C<sub>60</sub><sup>2+</sup> from C<sub>70</sub> singularity in  $T_{\text{trans}}$  remains also in the present measurements. Probably the C<sub>60</sub><sup>2+</sup> formation mechanism is quite different from the stepwise C<sub>2</sub> emission from C<sub>70</sub> that has widely been accepted for the other fragments.



**Figure 1.** *Y*–*t* map of parent and fragment ions produced from  $C_{70}$  with the excitation *hv* of 110 eV. The *y* coordinate is proportional to the *y* component of the ion velocity.

# 2. Estimation of the Translational Temperatures of High-Speed Fragment Ions Released from Sumanene and Coronene

Dissociative photoionization processes of polycyclic aromatic hydrocarbons (PAHs), sumanene C<sub>21</sub>H<sub>12</sub> and coronene

C<sub>24</sub>H<sub>12</sub>, were studied by a mass-resolved velocity map imaging (VMI) technique. Since PAHs constitute the substructures of  $C_{60}$ , their behaviors in fragmentation are expected to have some analogies with those of fullerenes. Nevertheless, the y-tmaps revealed a large difference in fragmentation mechanisms between PAHs and C<sub>60</sub>. Main species produced from C<sub>21</sub>H<sub>12</sub> and C<sub>24</sub>H<sub>12</sub> were singly-charged ions with very high speed: CH<sup>+</sup>, C<sub>2</sub>H<sub>2</sub><sup>+</sup> and C<sub>n</sub>H<sub>n/2</sub><sup>+</sup> (2 < n < 13). The translational energies of these fragments are so large that the scattering distributions extend beyond the coverage of our imaging apparatus. The peak profiles of the time-of-flight distributions were computer simulated on the basis of the ion trajectories calculated by the Simion-3D software. The translational temperatures of fragments from C<sub>21</sub>H<sub>12</sub> were thus estimated to be 3000 to 30000 K (see Figure 2), depending on the number of carbon atoms in the fragments. A similar trend was observed for dissociation of C24H12. Our results suggest that dissociation of PAHs proceeds through the fission mechanism, which is not dominant in the fullerene cases.



**Figure 2.** Translational temperatures of the fragment ions from coronene at hv = 100 eV. Error bars show the standard deviations of repeated experimental runs.

### 3. Enhancement of the Solubility, Thermal Stability, and Electronic Properties of Carbon Nanotubes Functionalized by MEH-PPV

Multi-walled carbon nanotubes functionalized with poly [2-methoxy-5-(2'-ethylhexyloxy)-1,4-phenylenevinylene] (MWCNT-f-MEH-PPV) nanocomposites were successfully prepared by employing a "grafting from" approach.<sup>4)</sup> The content of the functionalizing MEH-PPV in the composites was observed as 76% wt. Compared with pristine MWCNT, p-MWCNT, the aqueous solubility and thermal stability of the former are significantly enhanced. The effect of covalently and non-covalently functionalized nanotubes on the performances of the DSSCs has been investigated. The cells having the counter electrodes coated with isolated MEH-PPV, p-MWCNT/ MEH-PPV, and MWCNT-f-MEH-PPV/MEH- PPV were fabricated. The cells based on a MWCNT- f-MEH-PPV/MEH-PPV counter electrode demonstrate the best photovoltaic performance as observed by higher  $J_{SC}$ ,  $V_{OC}$ , and *FF* values. The experimental phenomena can be explained by quantum chemical calculations: Charge transfer from MEH-PPV oligomers to nanotubes is greater when covalently functionalized than when non-covalently functionalized. These suggest that the improvement in the photovoltaic parameters of the cells containing covalently functionalized nanotubes results not only from the higher concentration present in the nanotube films of the counter electrode, but also from the greater electron delocalization between the oligomers and nanotubes. The obtained results are very useful for enhancement of functionalized nanotubes applied to DSSCs.

### 4. Gas Phase Spectroscopy of Carbon Nanotubes

Our efforts have been concentrated on elucidation of the electronic and optical properties of fullerenes as well as their reaction dynamics in the gas phase. Such knowledge is critical in applying fullerenes to molecular functional devices. Recently, we started to investigate the properties of gaseous carbon nanotubes (CNTs) which of course form another group of promising carbon material.

CNTs are vaporized by matrix assisted laser desorption ionization (MALDI) and analyzed by time-of-flight (TOF) mass spectrometry. Commercially available multiwall CNT powder and single-wall CNT aqueous dispersion were used as samples. We have checked several solvents and matrices to optimize the MALDI conditions. The images of electron microscope prove that the laser irradiation certainly gives rise to vaporization of the sample. A temporal increase in the pressure of the sample chamber also supports the vaporization during the laser irradiation. Nevertheless, the signal of whole CNTs has not been observed yet. Smaller fragment ions  $C_n$  (n< 200) were only detected. Interaction between matrices and dispersants may prevent effective ionization.

We will continue to seek for the suitable vaporization method and construct a new vacuum chamber for the gas phase CNT experiment. The chamber will be connected to the beamline 4B in the UVSOR facility. Photoabsorption and photoionization experiments will be performed using the same chamber at the outset.

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# Light Source Developments by Using Relativistic Electron Beams

# UVSOR Facility Division of Advanced Accelerator Research



KATOH, Masahiro ADACHI, Masahiro KONOMI, Taro OHIGASHI, Takuji TANAKA, Seiichi ARAI, Hidemi TAIRA, Yoshitaka GOTO, Yoshiaki WASA, Naoki HIDA, Yohei UEMATSU, Youhei NIWA, Takahiro

This project involves researches and developments on synchrotron light source, free electron laser, beam physics and their related technologies, including application of the light sources.

### 1. Developments on UVSOR Accelerators

The magnetic lattice of UVSOR was modified in 2012. This was the second major upgrade of the lattice, following the first one in 2003. This time all the bending (dipole) magnets were replaced with newly designed combined function ones. They have specially designed pole shapes and edge shapes as shown in Figure 1, to produce dipole, quadrupole and sextupole fields at the same time. By this modification, the emittance was reduced from 27 nm-rad to around 15 nm-rad, which would results in higher brightness of the synchrotron radiation. The magnetic lattice function changed, as shown in Figure 2. The storage ring was successfully commissioned in July with the new lattice. Fine machine tuning is in progress. After this upgrade, the ring is called UVSOR-III.



Figure 1. New combined function bending magnet (a lower half is shown).

Professor Assistant Professor Assistant Professor Assistant Professor Post-Doctoral Fellow Post-Doctoral Fellow Graduate Student\* Graduate Student\* Graduate Student\* Graduate Student\* Graduate Student\*



**Figure 2.** Lattice functions of UVSOR-II (upper) and UVSOR-III (lower). One quadrant of the ring is shown. The emittance is 27nm-rad and ~15nm-rad, respectively.

### 2. Construction of STXM Beam-Line

A new in-vacuum undulator was constructed and, in May 2012, it was installed at the last straight section in the ring reserved for insertion devices. This is the sixth undulator at UVSOR. It would provide soft X-rays to a scanning transmission X-ray microscope (STXM) beam-line.
Since April in 2012, the STXM beam-line, BL4U, has been constructed. Recently, its construction and installation of the STXM system were finished and their commissioning is due to start in October, 2012. We have been preparing the facilities of the beam-line, such as microscopes for preobservations, a glove-box, an experimental hutch and infrastructures, for stating the operation in next April.

#### 3. Light Source Developments

We have demonstrated that coherent synchrotron radiation of various properties could be generated in an electron storage ring by using an external laser source.<sup>1–3)</sup> This research is supported by the Quantum Beam Technology Program of JST/ MEXT. Under this support, a new experimental station has been constructed.<sup>4)</sup> A new optical klystron/undulator was installed as shown in Figure 3. The magnetic field properties were verified by observing the spontaneous synchrotron radiation. The upgrade of the laser system was completed. The laser transport line, which was carefully designed as considering the effects of the air fluctuation, was also constructed. Two new beam-lines dedicated to the coherent lights in the VUV range and in the THz range has been almost completed.<sup>5)</sup> The generation of coherent synchrotron radiation at the new site was successfully demonstrated. Some applications will be demonstrated soon in this fiscal year.



Figure 3. Experimental ser-up of Coherent Harmonic Generation at UVSOR-II.

A tunable, quasi-monochromatic ultra-short pulse and polarization-variable gamma-ray source is under development, based on a technology called Laser Compton scattering. The laser photons are Compton back-scattered by the high energy electrons and are converted to gamma-rays. The electron beam circulating in the storage ring is very thin in the vertical direction, typically in the order of 10 microns. By injecting laser light from the vertical direction to the beam, it is possible to produce ultra-short, quasi-monochromatic, energy tunable, polarization variable gamma-ray pulses. The energy tunability was successfully demonstrated.<sup>6)</sup> An application utilizing the short pulse property was successfully demonstrated.

#### 4. Accelerator Technology Developments

A novel beam injection scheme using a pulse sextupole magnet has been studied, in which the perturbation to the stored beam could be minimized during the injection. This is beneficial to the users experiments in the top-up operation mode. A new pulse sextupole magnet was designed and has been constructed. The field measurement is in progress. This device will be installed in the ring in September 2012.

A turn-by-turn beam position measurement system has been constructed, collaborating with the Equipment Development Center. This system enables to measure the electron beam orbit turn-by-turn just after the injection. It was proved that this device was a powerful tool for the commissioning of the storage ring.

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### Synchrotron Radiation Spectroscopy on Strongly Correlated Electron Systems

#### UVSOR Facility Division of Advanced Solid State Physics



KIMURA, Shin-ichi MATSUNAMI, Masaharu IMURA, Keiichiro IIZUKA, Takuya HAJIRI, Tetsuya NIWA, Ryosuke SHIMURA, Yusuke HIRATE, Satoshi Associate Professor Assistant Professor Post-Doctoral Fellow<sup>\*</sup> Graduate Student<sup>†</sup> Graduate Student<sup>‡</sup> Graduate Student<sup>‡</sup> Graduate Student<sup>‡</sup> Graduate Student<sup>‡</sup>

Solids with strong electron–electron interaction, namely strongly correlated electron systems (SCES), have various physical properties, such as non-BCS superconducting, colossal magneto-resistance, heavy fermion and so on, which cannot be predicted by first-principle band structure calculation. Due to the physical properties, the materials are the candidates of the next generation functional materials. We investigate the mechanism of the physical properties as well as the electronic structure of SCES, especially rare-earth compounds, organic superconductors and transition-metal compounds, by infrared/THz spectroscopy and angle-resolved photoemission spectroscopy based on synchrotron radiation. Since experimental techniques using synchrotron radiation are evolved rapidly, the development of the synchrotron radiation instruments is also one of our research subjects.

### 1. Existence of Heavy Fermions in the Antiferromagnetic Phase of Celn<sub>3</sub><sup>1)</sup>

Recently, physics at the quantum critical point (QCP), which is the border between local magnetism and itinerant paramagnetism at zero temperature, has become one of the main topics in the condensed-matter field because new quantum properties such as non-BCS superconductivity appear in the vicinity of the QCP. The ground state of rare-earth intermetallic compounds, namely, heavy-fermion (HF) materials, changes between the local magnetic and itinerant nonmagnetic states through external perturbation by such factors as pressure and magnetic field. The QCP appears owing to the energy balance between the local magnetic state based on the Ruderman-Kittel-Kasuya-Yoshida (RKKY) interaction and the itinerant HF state due to the Kondo effect. In the itinerant HF regime, the conduction band hybridizes with the nearly local 4f state, so that a large Fermi surface as well as the hybridization band between them, namely, the c-f hybridization band, is realized. In the case of a magnetic regime, on the other hand, two theoretical scenarios have been proposed.

One is the spin-density wave (SDW) scenario based on spin fluctuation, in which large Fermi surfaces due to c-f hybridization remain even in magnetically ordered states. The other is the Kondo breakdown (KBD) scenario, in which the c-fhybridization state disappears in the magnetic state and only small Fermi surfaces due to conduction electrons appear. Many controversies for these scenarios have been performed so far, but the conclusion has not been obtained yet.

Here, we report the pressure-dependent electronic structure as well as the *c*-*f* hybridization state obtained by farinfrared reflectivity [ $R(\omega)$ ] and optical conductivity [ $\sigma(\omega)$ ] measurements of CeIn<sub>3</sub> under pressure. CeIn<sub>3</sub> has an AFM ground state with a Néel temperature  $T_N$  of 10 K. With the



**Figure 1.** (a) Pressure dependence of the reflectivity  $[R(\omega)]$  spectrum (solid circles) of CeIn<sub>3</sub> and Drude-Lorentz fitting results (dotted-dashed lines) in the photon energy range of 14–27 meV at 5.6 K. The spectra are shifted by 0.1 for clarity. (b) The edges of  $R(\omega)$  spectra as functions of pressure. The size of the marks denotes the intensity of the corresponding edge in the  $R(\omega)$  spectra. The pressure-dependent Néel temperature (solid and open triangles,  $T_N$ ) and valence transition temperature (solid squares,  $T^*$ ) are also plotted at the bottom.

application of pressure,  $T_N$  monotonically decreases and disappears at a critical pressure of approximately 2.6 GPa. We observed that the *c*-*f* hybridization gap appears not only in the HF state but also in the AFM state, and both the energy and intensity of the  $\sigma(\omega)$  peak due to the *c*-*f* hybridization band continuously increase with the application of pressure as shown in Figure 1. Our observations suggest that the electronic structure of CeIn<sub>3</sub> in the AFM phase can be explained by the SDW scenario because the *c*-*f* hybridization state exists even in the AFM phase.

#### 2. Three-Dimensional Electronic Structure and Interband Nesting in the Stoichiometric Superconductor LiFeAs<sup>2)</sup>

Recently discovered iron pnictide superconductors have two-dimensional (2D) Fe–As layers that are similar to the Cu–O planes in high- $T_c$  superconducting cuprates. In high- $T_c$ cuprates, 2D magnetic interaction is important for the origin of the high  $T_c$  because of their 2D electronic structure. Such 2D interaction as well as the 2D nesting condition in iron pnictides has long been a focus of discussions. However, the crystal structures of iron pnictides are more three dimensional (3D) than those of cuprates. It is therefore important to clarify the 3D electronic structure of iron pnictides in order to understand the effective interaction of Cooper pair formation.

We reported the electronic structure as well as the orbital characters of a stoichiometric iron-based superconductor LiFeAs ( $T_c = 18$  K) using polarization-dependent 3D angleresolved photoemission spectroscopy (ARPES) as shown in Figure 2. The obtained band dispersions and orbital characters are qualitatively in good agreement with those derived from local density approximation (LDA) band calculations. Considering a 3D nesting condition, we find that each 2D hole and electron Fermi surface (FS) of  $d_{xy}$  orbital character is weakly nested. This weak nesting suggests that ( $\pi$ , $\pi$ , $\pi$ ) interband scattering is important for the superconducting behavior of LiFeAs.

## 3. Momentum-Dependent Hybridization Gap and Dispersive In-Gap State of the Kondo Semiconductor SmB<sub>6</sub><sup>3)</sup>

Materials with strong electron correlation have exotic physical properties that cannot be predicted from first-principle band calculations. One example may be seen in a semiconductor with a very small energy gap, which appears in rare-earth compounds such as the Kondo semiconductor or insulator (KI). At high temperatures, KI behaves as a dense Kondo metal, while an energy gap with activation energy of several 10 meV appears at low temperature. The energy gap is believed to originate from hybridization between the nearly



**Figure 2.** (a)ARPES spectra at the Z point. The solid circles are values that are expected and the open squares are those that are not expected by the band calculation. (b) ARPES image with the band calculation at the Z point. (c) Schematic band dispersions obtained from the experiments. The open squares used in (c) have the same meaning as in (a) and (b). (d) Schematic figure of the 3D FS nesting conditions. The solid and dashed lines indicate the original and nested FSs, respectively. The bold arrow indicates the expected nesting wave vector of  $(\pi,\pi,\pi)$ .

localized 4*f* state near the Fermi level ( $E_F$ ) and the conduction band (*c*–*f* hybridization).

We reported the temperature dependence of the dispersion curve of the hybridization state using temperature-dependent 3D-ARPES, in order to determine the electronic structure and the reason for the different temperature dependences of the valence transition and magnetic excitation. We found that the hybridization band with a peak at a binding energy of 15 meV near the X point gradually appears on cooling from 150 to 40 K, which has the same temperature dependence as the valence transition. At the  $\Gamma$  point, on the other hand, the peak at  $E_{\rm B} \sim$ 20 meV has the same temperature dependence as the magnetic excitation at Q = (0.5, 0.5, 0.5), which differs from the 15-meV peak at the X point. This suggests that the magnetic excitation originates from the hybridization band at the  $\Gamma$  point.

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<sup>\*</sup> Present Address; Graduate School of Science, Nagoya University

<sup>†</sup> Present Address; Physikalisches Institut, Universität Stuttgart, Germany

<sup>‡</sup> carrying out graduate research on Cooperative Education Program of IMS with Nagoya University

### Electronic Structure and Decay Dynamics in Atoms and Molecules Following Core Hole Creation

UVSOR Facility Division of Advanced Photochemistry



SHIGEMASA, Eiji IWAYAMA, Hiroshi ISHIKAWA, Lisa Associate Professor Assistant Professor Post-Doctoral Fellow

The dynamics of the inner-shell photoexcitation, photoionization, and subsequent decay processes is much more complex, in comparison to outer-shell photo-processes. For instance, the inner-shell photoionization is concomitant with the excitation and ionization of valence electrons, which reveal themselves as shake-up and shake-off satellite structures in the corresponding photoelectron spectrum. The one-photon multielectron processes, which are entirely due to the electron correlation in the system, are known to happen not only in the primary inner-shell hole creation processes, but also in their relaxation processes. Our research project is focused on elucidating the electronic structures and decay dynamics in core-excited atoms and molecules, by utilizing various spectroscopic techniques together with monochromatized synchrotron radiation in the soft x-ray region.

#### 1. Ultrafast Molecular Dissociation of Core-Excited HBr Studied by High-Resolution Electron Spectroscopy

Soft X-ray absorption spectra of molecules exhibit rich structures in the region below the ionization thresholds, which are due to the excitations of a core electron to unoccupied valence or Rydberg orbitals. The core excited states are predominantly relaxed via Auger electron emission, in the case of the molecules composed of light elements, and subsequently fragmentation follows. As demonstrated by Morin and Nenner, however, a fast neutral dissociation could precede the resonant Auger decay.<sup>1)</sup> In other words, the electronic decay of the core hole takes place after the constituent atoms come apart. Since then, many research works have been conducted to identify such ultrafast dissociation processes in various different systems. Recent works on high-resolution resonant Auger electron spectroscopy have revealed that the nuclear motion of the molecular core-excited states is promoted in competition with the Auger decay. Here, we revisit the first discovery of ultrafast dissociation following the Br 3d core excitation in HBr. High-resolution electron spectroscopy for the subsequent Auger decay has been applied.

The experiments were carried out on the soft X-ray beamline BL6U at UVSOR. The radiation from an undulator was monochromatized by a variable included angle varied linespacing plane grazing monochromator. The exit slit opening was set to 300 µm, which corresponds to the photon energy resolution  $E/\Delta E$  of ~1500 at 70 eV. The monochromatized radiation was introduced into a gas cell with sample gases. Kinetic energies of the emitted electrons were measured by a hemispherical electron energy analyzer (MBS-A1) placed at a right angle relative to the photon beam direction. The degree of the linear polarization of the incident light was essentially 100%, and the direction of the electric vector was set to be parallel to the axis of the electrostatic lens of the analyzer. The energy resolution of the analyzer was set to ~12 meV. Under these experimental conditions, the full width at half maximum of the vibrational fine structure for the X state of HBr<sup>+</sup> was measured to be ~50 meV.

Figure 1 represents a resonant Auger spectrum taken at 70.4 eV, where the  $3d_{5/2}$  core-hole states are mainly populated. The vertical scale of the red line spectrum is magnified by the five times. The most marked feature in Figure 1 is the numerous sharp peaks, due to the vibrational structures for the HBr<sup>+</sup> states and the atomic Auger lines from the Br fragments with a 3d core-hole. In the previous work,<sup>1)</sup> only five atomic lines indicated by the blue arrows in Figure 1 have been identified, owing to the limited resolution. In contrast, at least ten more

atomic Auger peaks and some molecular peaks with vibrational structures are clearly resolved in the present work. The detailed analyses for the spectral features observed are just beginning to be performed.



Figure 1. Resonant Auger spectrum taken at the Br  $3d_{5/2} \rightarrow \sigma^*$  resonance of HBr.

#### 2. Angle-Resolved Metastable Fragment Yields Spectra

Multi-electron processes, where more than two electrons are excited simultaneously, have usually been observed in absorption and photoelectron spectra of inner-shell excited molecules. Double excitations, one of the typical multielectron processes involving two electrons, are often observed above the ionization threshold energy region in photoabsorption spectra. Shake-up excitations in ions, where one of two electrons is excited into the continuum, are often observed in the energy region close to resonance features in photoelectron spectra. These processes occur mainly due to correlation effects among the electrons in a system and are fundamental aspects of atomic and molecular physics.

Recently, we demonstrated that "metastable" fragments spectroscopy, in which highly excited neutral fragments are observed, is quite useful as a spectroscopic technique for the investigation of multi-electron processes.<sup>2)</sup> Metastable photofragment spectroscopy can also be used for investigating shake-up satellites at the threshold excitation energy, because neutral photofragments can be detected wherever the scanning photon energies match the threshold energies of the ionic states. In addition, if one measures the angular distribution of the fragment emission with respect to the polarization direction of the synchrotron radiation, the symmetries of the innershell excited states can be deduced from this information. The angular distribution of fragments emitted after core-hole creation is related to the molecular orientation upon photoabsorption, because the lifetime of the core-hole  $(\tau \sim 10^{-14} \text{ s})$ is dominated by Auger decay leading to dissociative states and is much shorter than the molecular rotational period ( $\tau \sim 10^{-10}$ s). Thus, symmetry-resolved, namely, angle-resolved metastable spectroscopy provides complete symmetry resolution between the  $\Delta \Lambda = 0$  (parallel) and  $\Delta \Lambda = \pm 1$  (perpendicular) transitions in the *K*-shell photoabsorption of diatomic molecules. As a result, the  $\Sigma$ - or  $\Pi$ -symmetry character of transition states can be detected.<sup>3)</sup>

Figure 2 demonstrates a zero kinetic energy electron (ZEKE) spectrum, and angle-resolved metastable yield spectra of CO in the C K-edge region, measured at  $0^{\circ}$  ( $\Sigma$ -symmetry) and 90° (II-symmetry), respectively, as an example of successful measurements. The photon energy range shown in Figure 2 corresponds to double and triple excitation regions, where shake-up satellite states also lie. Several peaks are clearly separated by  $\Pi$ -symmetry and  $\Sigma$ -symmetry spectra. We observed new peaks at 310.2 eV, 313.2 eV, and 313.9 eV. These should be double/triple excitation peaks since angleresolved photoion yield spectra show small peaks at the same energy positions and these peaks have not been observed by conventional photoelectron spectroscopy. We successfully showed that all peaks in angle-resolved metastable fragment yields spectra of CO in the C 1s ionization threshold region can be assigned as either satellite states or double/triple excitation states.



**Figure 2.** ZEKE spectrum, and angle-resolved metastable yield spectra, measured at  $0^{\circ}$  ( $\Sigma$ -symmetry) and  $90^{\circ}$  ( $\Pi$ -symmetry), in the satellite and double/triple excitation regions of CO in the C *K*-edge region.

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### **Micro Solid-State Photonics**

#### Laser Research Center for Molecular Science Division of Advanced Laser Development



TAIRA, Takunori ISHIZUKI, Hideki AKIYAMA, Jun TSUNEKANE, Masaki SATO, Yoichi KAUSAS, Arvydas KONG, Weipeng ONO, Yoko INAGAKI, Yayoi Associate Professor Assistant Professor IMS Research Assistant Professor Post-Doctoral Fellow Post-Doctoral Fellow Graduate Student Secretary Secretary

The artistic optical devices should be compact, reliable, efficient and high power light sources. With the approaches of domain structures and boundaries engineering, it is possible to bring the new interaction in their coherent radiation. The high-brightness nature of Yb or Nd doped single crystal or ceramic microchip lasers can realize efficient nonlinear wavelength conversion. In addition, designed nonlinear polarization under coherent length level allows us new function, such as the quasi phase matching (QPM). The development of "*Micro Solid-State Photonics*," which is based on the micro domain structure and boundary controlled materials, opens new horizon in the laser science.

### 1. High Brightness, Passively Q-Switched Yb:YAG/Cr:YAG Micro-Laser

A QCW diode end-pumped, high peak power passively Q-switched Yb:YAG/Cr:YAG micro-laser was demonstrated as shown in Figure 1. An output pulse energy of 3.6 mJ was obtained with a duration of 1.3 ns. The peak power is 2.8 MW. An  $M^2$  value was nearly 1.



(a) Beam profile (b) Pulse duration Figure 1. Beam profile and pulse duration at 3.6 mJ.

#### 2. 60% FHG Efficiency from Fluxless-Grown BBO Using Nd:YAG/Cr<sup>4+</sup>:YAG Microchip Laser

High-efficiency, compact ultra-violet (UV) sources are

desirable for many applications, such as, ultrafast UV spectroscopy, photolithography, micromachining, *etc*.

We have developed a Nd:YAG/Cr<sup>4+</sup>:YAG microchip laser which gives > 8 MW peak power at 100 Hz. We used LBO crystal to obtain 85% second harmonic generation (SHG) and a BBO crystal, grown by a new fluxless method, to obtain 60% fourth harmonic generation (FHG) as shown in Figure 2. A BBO crystal grown by the conventional flux technique can give only 40% FHG under identical conditions. We can obtain a stable pulse train having 3.4 MW peak power with 250 ps pulse width and 100 Hz repetition rate at 266 nm wavelength. We believe that this is the highest conversion efficiency reported so far for the BBO crystal in the generation of deep ultraviolet light.



Figure 2. FHG conversion characteristics by different BBO crystal.

#### 3. Lens-Less Edge-Pumped High Power Microchip Laser

A lens-less edge-pumped microchip laser was realized, which is directly pumped by single-emitter diode chips from multiple directions. Figure 3 shows the scheme of this laser head. The lens-less design makes the microchip laser more compact; the multi-direction designable pump schemes can realize desired pump shape. As a preliminary result, continuous wave output power of 32.5 W and slope efficiency of 45% was obtained by 9-direction pumped Yb:YAG ceramic microchip. Further, 27.2 W single peak Gaussian beam was obtained by a small adjustment of the output mirror. Besides, watt-level high-order Hermite–Gaussian mode, doughnutshape mode and vortex arrays were demonstrated as an evidence of designable pump scheme. Power scalability is easy by increasing the number of pump diodes to achieve hundredwatt-level high power microchip laser.



Figure 3. The scheme of lens-less edge-pumped microchip laser.

#### 4. Temperature Dependencies of Stimulated Emission Cross Section for Nd-Doped Solid-State Laser Materials

Temperature dependencies of stimulated emission cross section for Nd:YAG, Nd:YVO<sub>4</sub>, and Nd:GdVO<sub>4</sub> were carefully evaluated. Our spectral evaluations with fine spectral resolution were carried out under the condition that the population inversion was induced into samples by a weak pumping field. Within the temperature range from 15 °C to 65 °C, the variation of emission cross section at 1.06 µm in Nd:YAG was -0.20%/°C, while those in Nd:YVO<sub>4</sub> and Nd:GdVO<sub>4</sub> for  $\pi$ -polarization were -0.50%/°C and -0.48%/°C, respectively. Consideration of measured temperature dependence gave the numerical model for temperature dependent emission cross sections of Nd-doped solid-state laser materials. We have also presented numerical approximations of this model for our samples by a simple polynomial, which can be applicable as shown in Figure 4.



Figure 4. Temperature dependence of emission intensity of various Nd-doped laser media.

#### 5. Fabrication of 10-mm-Thick Periodically Poled Mg-Doped Congruent LiNbO<sub>3</sub> Device for High-Energy Wavelength Conversion

Periodically poled Mg-doped congruent  $LiNbO_3$  device with 10mm thickness of 32.2 µm period was fabricated by 33 kV pulse application. High-energy optical-parametric oscillation with half-joule-class output energy can be expected by joule-class pumping source in 10 ns pulse region.



Figure 5. 10-mm-thick PPMgLN device with QPM period  $\Lambda = 32.2 \ \mu m$ .

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### **Ultrafast Laser Science**

#### Laser Research Center for Molecular Science Division of Advanced Laser Development



FUJI, Takao NOMURA, Yutaka KAWAI, Shigeko Associate Professor Assistant Professor Secretary

Speed of ultrafast energy transfer from light to molecules (*i.e.* primary processes of photosynthesis, photoisomerization in visual pigments, *etc.*) is on the order of femtosecond  $(10^{-15} \text{ s})$ . In our laboratory, we develop cutting edge lasers for such ultrafast molecular science, namely, femtosecond or attosecond  $(10^{-18} \text{ s})$  ultrashort pulse lasers.

For example, arbitrary waveform synthesis can be performed with simultaneous generation of femtosecond light pulses in various wavelength regions and superimposition of them with precisely controlled phases.

We would like to develop such advanced light control technology, which can push forward the research on ultrafast photochemical reactions.

## **1.** Generation of Phase-Stable Half-Cycle Mid-Infrared Pulses through Filamentation in Gases<sup>1)</sup>

Filamentation of powerful ultrashort laser pulses in gases is one of the most interesting phenomena in nonlinear optics. The balance between self-focusing and plasma self-defocusing makes the pulse propagate much longer than the Rayleigh range with a very high intensity. It results in a dramatic enhancement of nonlinear processes occurring in the filamentation zone. This phenomenon enables high intensity pulse compression and efficient nonlinear wavelength conversion with gas media.

Enhanced nonlinear-optical processes in laser-induced filaments suggest a new strategy for the generation of ultrashort pulses of long-wavelength radiation. Ultrabroadband mid-infrared (MIR, 3–20  $\mu$ m) pulse generation through filamentation in air was firstly demonstrated in 2007,<sup>2)</sup> and the technique was followed by several groups. Such MIR pulses with more than one octave at full width at half maximum are very attractive to be applied for molecular spectroscopy, *e.g.* two-dimensional infrared spectroscopy.

Here we report the latest progress of the ultrabroadband MIR pulse generation through filamentation. By using argon

gas as the nonlinear medium and purging the optical path for the MIR pulse with the argon gas, a long tail of the pulse due to free induction decay of atmospheric carbon dioxide and water vapor disappeared, and it was possible to achieve generation of high contrast MIR pulses. Full characterization of the pulse shape of the MIR field indicates that its pulse duration was 7.4 fs, which is about half-cycle period of the center wavelength ( $3.9 \mu m$ ) of the pulse.

The experimental setup is shown in Figure 1. The light source was based on a Ti:Sapphire multi-pass amplifier system (Femtolasers, 800 nm, 25 fs, 0.9 mJ at 1 kHz). The second harmonic ( $\omega_2$ , 25 µJ) and fundamental ( $\omega_1$ , 675 µJ) pulses were spatially and temporally overlapped and focused into argon by a concave mirror (r = -1000 mm), generating a bright filament with a length of ~3 cm around the beam focus. This filament generated an MIR pulse ( $\omega_0$ ) through an ionization-assisted wave mixing process ( $\omega_1 + \omega_1 - \omega_2 \rightarrow \omega_0$ ). The energy of this MIR pulse was measured as ~250 nJ by using a pyroelectric detector (J-10MB-LE, Coherent). With this energy level, it is possible to apply the pulses for the nonlinear spectroscopy of condensed matter. The pulse-topulse intensity fluctuation was about 2.5% rms.



**Figure 1.** Schematic of the system. The shaded area was purged with argon. BS: beam splitter (5% reflection), BBO:  $\beta$ -BaB<sub>2</sub>O<sub>4</sub> crystal (Type 1,  $\theta$  = 29 deg, *t* = 0.1 mm), D: Dichroic mirror, P: Periscope, CM1: *r* = 1 m concave mirror, CM2: *r* = 0.5 m concave mirror, MH: Aluminium-coated mirror with a hole ( $\phi$  = 7 mm), PM: Aluminium-coated parabolic mirror, BF: Bandpass filter for 335–610 nm (FGB37, Thorlabs), OMA: Spectrometer for ultraviolet region (USB2000+, OceanOptics).



**Figure 2.** (a) Experimental and (b) simulated radial intensity distributions of the mid-infrared pulses.

The beam profile of the MIR beam after ZnSe and Si filters measured with a pyroelectric camera (Pyrocam III, Spiricon) is shown in Figure 2(a). The shape of the beam was ring and the angle of the cone was estimated to be about 3 deg. Some asymmetric shape and distortion from ideal ring pattern comes from residual pulse front tilt and/or astigmatism of the light source. The generated MIR pulse has basically pure one-direction linear polarization (>40:1) in the entire cross-section of the beam as the input pulses, which fact was confirmed with a wire grid polarizer (NT62-774, Edmund).

We compare the experimental result with FWM-beam analysis based on a straightforward integration of the FWM response over the beam overlap region. As can be seen from Figure 2(b), the simple approach provides an accurate agreement with the experimental result. It confirms that the ringshaped beam profile originates from a dramatic confocalparameter mismatch between the MIR field and the laser beams.

Additionally, the ~12 mm diameter beam was focused down to 1.0 mm with a r = 2 m concave mirror, indicating a reasonable focusability for a ring shaped spatial mode. The beam profile at the focal point. Although the beam may contain some angular dispersion, the dispersion is basically radially symmetric, and thus does not significantly deteriorate the good focusability of our MIR beam.

In order to quantitatively evaluate the temporal shape of the generated MIR pulse, we measured cross-correlation frequency-resolved optical gating (XFROG) We used argon again as a nonlinear medium and used four-wave mixing process ( $\omega_1 + \omega_1 - \omega_0 \rightarrow \omega_2$ ) as a nonlinear interaction between the test pulse (MIR pulse) and the reference pulse. The scheme is free from spectral filtering caused by phase matching condition in the nonlinear interaction.

The system for the XFROG measurement is also shown in Figure 1. Small portion (~2  $\mu$ J) of the fundamental 25-fs pulse was used as a reference pulse. The reference pulse and the MIR pulse (test pulse) were combined through a mirror with a hole and focused into argon with an aluminium-coated parabolic mirror (f = 150 mm). Generated blue spectra (centered around 440 nm) were measured with a spectrometer (USB 2000+, OceanOptics) by scanning the delay time ( $\tau$  in Figure 1) between the reference pulse and the MIR test pulse.

The measured and the retrieved XFROG traces are shown



Figure 3. (a) Experimental and (b) retrieved XFROG traces. The retrieved pulse in (c) time and (d) frequency domain. The temporal profile from the 3D numerical simulation and the spectrum measured with Fourier transform spectrometer are also shown. Measured XFROG trace.

in Figure 2(a) and (b). The main feature of the trace indicates that the residual chirp of the test pulse is very small. The FROG error was 0.0009 with  $256 \times 256$  grid. The retrieved time and frequency domain pictures are shown in Figure 2(c) and (d), respectively. The pulse width is estimated to be 7.4 fs, which is 0.57 cycles for 3.9 µm carrier wavelength. The retrieved spectrum is spread over whole MIR region (500–5000 cm<sup>-1</sup>). The broadness of the spectrum was due to the weak dispersion of the medium, with the phase-matching length exceeding the length of the filament for all the MIR spectral components observed in the experiments.

In conclusion, ultrabroadband coherent MIR spectrum which covers the entire MIR region was generated through two-color filamentation. Due to the spatial and temporal quality, ultrashort MIR pulses generated in two-color filaments are ideal for numerous applications. In our experiments, MIR pulses as short as 7.4 fs were generated, which corresponds to about a half-cycle of 3.9 µm center wavelength.

The light source has a potential to change the situation of traditional MIR spectroscopy dramatically. For example, the coherent broadband MIR light source enables us to obtain absorption spectra through entire MIR region by single-shot with chirped pulse up-conversion technique. The reasonable quality of the spatial mode can be useful for efficient MIR microscope imaging combined with the up-conversion technique. Multi-dimensional spectroscopy for entire MIR region to monitor vibrational coupling among very different vibrational modes can be realized with the light source.

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### **Visiting Professors**



#### Visiting Professor KONDOH, Hiroshi (from Keio University)

#### Surface Dynamic Processes Studied by Soft X-Ray Spectroscopy

Surface physics and surface chemistry have made significant progress in understanding of surface static structures but much less progress has been made with respect to understanding of surface dynamic processes. We have been interested in surface dynamic processes such as catalytic reactions at surfaces of metals and oxides. Recently we have been working on oxide photocatalysts using synchrotron radiation in a

soft-x-ray region. In particular, we studied the geometric and electronic structures of N-doped TiO<sub>2</sub> at BL3U and BL6U in the UVSORII and determined the local structure of the doped nitrogen that induces the visible-light response of the photocatalyst. Another research subject concerns understanding of catalytic reactions on metal nanoparticles which proceed under practical working conditions. A soft x-ray absorption spectroscopy with the transmission mode has been developed at BL3U in such a way that real-time observations of catalytic processes under gas flow with ambient-pressure conditions can be conducted. Catalytic reactions over metal nanoparticles that are enhanced by absorption of the gas molecules inside the particles will be studied with this technique.



#### Visiting Professor NODA, Susumu (from Kyoto University)

#### Strong Coupling of Single Atoms to Photonic Crystal Cavity Field

We have investigated photonic crystal structures which enable modification of propagation properties of an electromagnetic field and also tight confinement of the field to a tiny resonator. Accordingly the field strength inside the resonator is much enhanced and therefore the field can be strongly coupled to a quantum emitter such as a quantum dot even at a single photon level. Such a nanostructure device would be suitable for applications in optical communication and future quantum information processing in terms of its

scalability. We have studied the strong coupling of the cavity field with a quantum dot and also the Purcell effect. Recently we have been interested in adopting a single cold atom as a quantum emitter, which shows much longer coherence time and therefore would be desirable for future application. Cold atoms are first loaded into a magneto-optical trap and then one of them is captured in tightly-focused optical tweezers. A movable lens-positioner can translate the position of the focal point, thereby transferring the trapped atom to the vicinity of the photonic crystal cavity. With this technique, the strong coupling of the single atom with the cavity field will be studied.



#### Visiting Professor ITO, Atsushi (from Tokai University)

#### X-Ray Spectromicroscopy of Biomedical Specimens

Soft X-ray microscopy has a great advantage over other microscopies in the mapping of light elements or molecules containing such elements at high resolution. The mapping is realized by X-ray spectromicroscopy which utilizes unique spectral features of elements and molecules, that is, absorption edges and XANES profiles observed in the vicinity of the absorption edge. Our effort to apply spectromicroscopy to biomedical specimens has been focused on the mapping of sulfur oxidation state in

human hair, because hair consists of mainly cystine, a sulfur-rich amino acid. Cystine is oxidized to cysteic acid by oxidative damage. Both products were found to have significantly different XANES peaks at the S-K absorption edge. Spectromicroscopy in combination with an electronic zooming tube as a two dimensional detector revealed that the oxidative damage was preferentially generated in the outer part of hair called cuticle. Furthermore a bleaching treatment of hair also increased the content of cysteic acid.



#### Visiting Associate Professor **TSUBOUCHI, Masaaki** (from Japan Atomic Energy Agency)

#### Development of THz Tomography of Photo-Induced Carrier

We are developing a new technique for the THz tomography of photo-induced carriers in a semiconductor based on the optical-pump THz-probe reflection spectroscopy. Since the photo-induced carrier strongly interacts with the THz light, the measurement and control of the carrier distribution and dynamics are significantly important to design THz optics using semi-conductors. Our THz tomography is a noncontact technique which determines both the dynamics and the spatial distribution of photo-induced

carrier, simultaneously. We are demonstrating the THz tomography using photo-excited silicon (Si), and comparing the experimental results with the exact solutions of Maxwell's equations.

### **RESEARCH ACTIVITIES** Materials Molecular Science

Extensive developments of new molecules, molecular systems, and their higher-order assemblies are being conducted in three Divisions of Electronic Structures, Electronic Properties, and Molecular Functions, one division for visiting professors, and Research Center for Molecular Scale Nanoscience, in an attempt to discover new phenomena and useful functions. The physical properties of electronic, optical and magnetic properties on new functional materials are investigated, and moreover, the chemical properties like catalysis and photochemistry and technological applications like solar cells are also examined in this department.

### Novel Properties of Magnetic Ultrathin Films Studied by *In Situ* Spectroscopic Methods

#### Department of Materials Molecular Science Division of Electronic Structure



YOKOYAMA, Toshihiko NAKAGAWA, Takeshi TAKAGI, Yasumasa EGUCHI, Keitaro NAKANO, Hirohito FUNAKI, Yumiko TOYAMA, Yu YASUI, Yuco Professor Assistant Professor Assistant Professor Graduate Student Graduate Student\* Secretary Secretary (Nanotechnology Platform project) Secretary (Graduate school affairs)

Novel properties of magnetic metal ultrathin films have been attractive both from fundamental interest and from technological requirements. We are especially interested in drastic modification of metal thin films through the film– substrate interaction and/or a surface chemical treatment. The magnetic properties are characterized by means of several kinds of *in situ* spectroscopic methods like MOKE (Magneto-Optical Kerr Effect) using UV-visible lasers and XMCD (X-ray Magnetic Circular Dichroism) using synchrotron radiation soft X-rays, and UV magnetic circular dichroism photoelectron emission microscopy (UV MCD PEEM) using such ultrashort pulsed UV lasers.

## 1. Giant Magnetic Anisotropy and Coercivity in Fe Island and Atomic Wire on $W(110)^{1)}$

Hard magnets is an important industrial material both for data storage and power applications like electromotive actuators and power generators. Although the magnetic energy product  $(BH_{max})$  that characterizes hard magnets has improved in the last century, rare earth (RE) magnets like Nd<sub>2</sub>Fe<sub>14</sub>B have marked a saturation of the maximum available value of  $BH_{\text{max}} = 500 \text{ kJ/m}^3$  since 1990. Moreover, the replacement of the RE has been demanded due to the scarcity of RE, and novel approaches to improve hard magnets have attracted interests. Magnetic nanostructures with single atomic layers have a possibility to manifest prominent magnetic properties such as magnetocrystalline anisotropy energy (MAE) and coercivity  $(H_c)$  that remarkably differ from the properties of bulk materials. The Fe monolayer on W(110) has been extensively studied as an ideal two-dimensional ferromagnetic system with a pseudomorphic growth mode and an extremely large strain due to a large difference in the lattice constants between Fe and W. Owing to its extreme hard magnet character, the anisotropy energy or the coercivity has not been directly determined from the hard axis magnetization curve measurements. In this work, we directly measure the giant MAE as well as large coercivity in Fe nanodots and nanowires



Figure 1. STM images of nanowires for Fe 0.3 ML (~6 atom width) (left) and nanodots for Fe 0.4 ML. The top and bottom panels correspond to the spin-polarized dI/dV and topographic images, respectively. The right-top panel clearly shows magnetically single domain dots (arrows denote the magnetization directions along [1-10]).

grown on W(110) by using our ultrahigh vacuum superconducting magnet XMCD system installed in UVSOR-II. A remarkable difference between the dots and the stripes is found in the coercivity, although the MAE is similar for both systems.

Figure 1 shows the spin-polarized and topographic STM (scanning tunnel microscopy) images of Fe nanowires and nanodots. The measurements were performed in Johannes Gutenberg Mainz University through the collaboration with Prof. H. J. Elmers group. In the spin-polarized STM image of the nanodots, a clear contrast (darker and brighter islands) can be seen, implying that the nanodots are magnetically ordered as single domains with the magnetization easy axis of the [1-10] crystal axis.

Figure 2 shows the anisotropic magnetization curves of Fe 0.25 ML nanodots and nanowires at 5 K taken by measuring the Fe  $L_3$ -edge circularly polarized X-ray absorption intensities using Beamline 4B at UVSOR-II. The G.I.[1-10] and G.I.[001] directions imply that the magnetization direction is tilted by 35° toward the surface normal with respect to the given crystal axis direction. The apparent coercivities are found to be larger than ~4 T for both the nanodots and nanowires. The anisotropic magnetic field, which is obtained by the extrapolation of the hard-axis magnetization curve (green lines in Figure 2) up to the saturated magnetization, can be esti-



mated to be >15 T. Huge magnetic anisotropy can be concluded both for the nanodots and nanowires. By fitting the experimental data with a simple anisotropic magnetic energy model, the magnetic anisotropic constants were determined to be  $\sim$ 1 meV for both the in-plane and out-of-plane directions, which is a few hundred times larger than that of bulk *bcc* Fe.

Figure 3 shows the Fe *L*-edge XMCD spectra of the nanodots at 5 K. The spectra were normalized with the  $L_2$ -edge peak tops. As the Fe coverage (nanodot size) decreases, the  $L_3$ -edge negative peak intensity increases drastically, implying a larger orbital magnetic moment as shown in the right panel of Figure 3. The orbital magnetic moment of the edge Fe atoms may be enhanced.

## 2. Magnetic Properties of Fe Nanostructures on Si<sub>3</sub>N<sub>4</sub>/Si(111)-(8×8) and Ag/Si(111)-( $\sqrt{3}$ × $\sqrt{3}$ )R30°<sup>2)</sup>

The magnetic properties of ferromagnetic transition metals on Si substrates have been widely investigated for the exploitation of new magnetic devices. Since clean Si surfaces react with transition metals very easily to form usually non-magnetic transition-metal silicides, it is essential to insert some inert film between transition metals and Si substrate. No reports have been however published for epitaxially ordered substrates on Si(111). In the present study, we investigated growth processes and magnetic properties of Fe deposited on well-

**Figure 2.** Anisotropic magnetization curves of 0.25 ML Fe nanodots (a) and nanowires (b) at 5 K. The data were recorded by measuring the Fe  $L_3$ -edge circularly polarized X-ray absorption intensities. Huge coercivities along the [1-10] easy axis (black lines) and anisotropic magnetic fields along the hard axes (green) were observed for both nanodots and nanowires.

**Figure 3.** (left) Fe *L*-edge XMCD spectra of Fe nanodots on W(110) at 5 K along the G.I.[1-10] axis. (right) Ratio of the orbital magnetic moment ( $m_1$ ) relative to the effective spin magnetic moment ( $m_s^{eff}$ ). Smaller nanodots (lower Fe coverage) show larger orbital magnetic moments.

**Figure 4.** (left) Fe *L*-edge XMCD of Fe on clean Si(111) (1.6 ML), Ag/Si(111) (2.0 ML) and Si<sub>3</sub>N<sub>4</sub> (1.6 ML) at  $H = \pm 5$  T and T = 5 K. The X-ray incidence angles  $\theta$  are 0° (normal incidence) and 55° (grazing). (right) Magnetization curves of Fe on clean Si(111) (1.6 ML, black), Ag/Si(111) (2.0 ML, red) and Si<sub>3</sub>N<sub>4</sub> (1.6 ML, blue), taken at T = 5 K and  $\theta = 0^\circ$ .

defined Si<sub>3</sub>N<sub>4</sub>/Si(111)-(8×8) and Ag/Si(111)-( $\sqrt{3}\times\sqrt{3}$ ) R30° by using STM and XMCD.

Figure 4 show the Fe L-edge XMCD spectra and the corresponding magnetization curves of Fe deposited on Si<sub>3</sub>N<sub>4</sub>/ Si(111)-(8×8), Ag/Si(111)-( $\sqrt{3}\times\sqrt{3}$ )R30° and clean Si(111)- $(7 \times 7)$ . It is clearly found that the Fe XMCD signals is enhanced in the sequence of  $Si(111) < Ag/Si(111) < Si_3N_4/Si(111)$ . The XMCD sum-rule analysis yields the spin magnetic moments as  $m_s = 0.17 \ \mu_B$  (clean), 1.25  $\mu_B$  (Ag), and 2.62  $\mu_B$  (Si<sub>3</sub>N<sub>4</sub>). On clean Si(111)-(7 $\times$ 7), the spin magnetic moment is almost quenched. On the contrary, Fe/Si<sub>3</sub>N<sub>4</sub> has a much larger spin magnetic moment, which is even larger than that of bcc bulk Fe (2.2  $\mu_B$ ). Fe/Ag/Si(111) shows a intermediate property. Such a drastic difference among the three substrates is caused by the fact that the Si<sub>3</sub>N<sub>4</sub> substrate most effectively suppresses the silicide formation. Moreover, it is also found that the magnetic moment of Fe on Ag/Si(111) is lost quite easily by annealing at 500 K ( $m_s = 0.32 \mu_B$ ), while that of Fe on Si<sub>3</sub>N<sub>4</sub> is sufficiently kept even when the sample is annealed at 500 K  $(m_{\rm s} = 2.19 \ \mu_{\rm B}).$ 

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### Design and In-Situ Characterization of Catalyst Surfaces

#### Department of Materials Molecular Science Division of Electronic Structure



TADA, Mizuki MURATSUGU, Satoshi MAITY, Niladri ZHANG, Shenghong SAIDA, Takahiro WAKI, Minoru SODE, Aya THUMRONGPATANARAKS, Wipavee GAN, Raymond JIANG, Lu LIM, Min Hwee WANG, Fei ISHIGURO, Nozomu FUNAKI, Yukino USUI, Chika GONDO, Makiko FUKUTÓMI, Yukiyo

Associate Professor Assistant Professor Post-Doctoral Fellow Post-Doctoral Fellow Post-Doctoral Fellow Post-Doctoral Fellow Post-Doctoral Fellow JENESYS Program Visiting Scientist Visiting Scientist\* Visiting Scientist\* Visiting Scientist\* Graduate Student Graduate Student<sup>†</sup> Technical Fellow **Technical Fellow Technical Fellow** Secretary

#### 1. Preparation and Catalytic Performances of a Molecularly Imprinted Ru-Complex Catalyst with an NH<sub>2</sub> Binding Site on SiO<sub>2</sub>

The potential of immobilized metal-complex catalysts remarkably interplays with the nature of support surfaces, resulting in significant rate enhancements and unique catalytic performances that their homogeneous counterparts do not exhibit. We proposed the design of molecularly imprinted metal-complex catalysts on SiO<sub>2</sub> surfaces, whose ligand was used as a template, and molecularly imprinted metal-complex catalysts have been applied for selective catalysis.

Natural enzymes, which are quite selective catalysts for particular reactants, possess highly sophisticated catalytic systems with active metal species, shape-selective reaction space, and molecular recognition sites spatially arranged for the particular reactants. Artificial design of such a sophisticated catalytic system is still difficult and the preparation of both catalytically active site and shape-selective reaction space with spatially arranged molecular binding sites. We have prepared a molecularly imprinted Ru-complex catalyst with NH<sub>2</sub> binding site on the wall of a molecularly imprinted cavity toward the shape-selective transfer hydrogenation catalysis.<sup>1)</sup>

A template alcohol, which was the template for *o*-fluorobenzophenone hydrogenation, was coordinated to a SiO<sub>2</sub>supported Ru complex (Ru-*N-p*-styrenesulfonyl- 1,2-diphenylethylenediamine) and the hydrolysis polymerization of tetramethoxysilane produced SiO<sub>2</sub>-matrix overlayers surrounding the supported Ru complex on the surface. A carbamate group (-NHCOO-) connected to a silane-coupling moiety (-Si (OC<sub>2</sub>H<sub>5</sub>)<sub>3</sub>) was tethered to the alcohol template and the cleavage of the carbamate moiety after hanging the silanecoupling branch on the wall of a molecularly imprinted cavity produced a spatially arranged  $NH_2$  binding site for the hydrogenation of *o*-fluorobenzophenone as shown in Figure 1.



Figure 1. Preparation of molecularly imprinted Ru catalysts with an imprinted cavity with an NH<sub>2</sub> binding site on the wall of the cavity, with an imprinted cavity without an NH<sub>2</sub> binding site, and without an imprinted cavity.

We prepared three catalysts with a molecularly imprinted cavity with an  $NH_2$  binding site, with a molecularly imprinted cavity without an  $NH_2$  binding site, and without a cavity of the template illustrated in Figure 1. The catalytic performances of the transfer hydrogenation of *o*-fluorobenzophenone were significantly different between these three catalysts under identical reaction conditions. The transfer hydrogenation did not proceed at all on the catalyst prepared without the template. On the other hand, other two catalysts prepared with the templates exhibited activity for the hydrogenation, but we found differences in the hydrogenation rates of *o*-fluorobenzophenone and *o*-methylbenzophenone. These two reactants have similar shapes to each other, but *o*-fluorobenzophenone with F group, which can interact to the NH<sub>2</sub> binding site by hydrogen bonding, was preferably hydrogenated on the molecularly imprinted catalyst with the NH<sub>2</sub> binding site. Such differences in the hydrogenation activity were not observed on the imprinted catalyst without the NH<sub>2</sub> binding site. The NH<sub>2</sub> binding site promoted the preferable adsorption of the specific molecule that can interact with the molecular binding site, resulting in increase in the transfer hydrogenation activity. The strategy is promising to create an artificial enzymatic catalyst surfaces in the future.

#### 2. Operando Time-Resolved XAFS for Surface Events on Pt<sub>3</sub>Co/C and Pt/C Cathode Catalysts in Practical PEFCs during Voltage-Operating Processes

Polymer electrolyte fuel cells (PEFCs) are among the most efficient clean energy technologies, but practical application in automobiles remains challenging because of the high cost and insufficient durability of cathode catalysts. To improve fuelcell performance and cathode catalyst durability, alloying of Pt with 3d transition metal elements is a promising approach. Pt-Co alloy is a representative of the Pt-alloy cathode catalysts and has been reported to be more active and durable than Pt/C catalysts. Although electrochemical surface events have been extensively investigated, the structural kinetics of the transformations of catalyst themselves and the ORR reaction mechanism on practical Pt-Co catalysts have not been established. In situ time-resolved QXAFS spectra at the Pt LIII-edge and Co K-edge were measured at 500 ms intervals, and the analysis of the series of the operando QXAFS spectra for the PEFC bimetal cathode catalyst revealed the rate constants of electron transfer, changes in the charge density of Pt<sub>3</sub>Co nanoparticles, and changes in the local coordination structures with Pt-Pt, Pt-Co, Co-Co, and Pt-O bonds for the first time (Table 1).2)

We compared the reaction mechanism and structural kinetics on the Pt<sub>3</sub>Co/C and Pt/C catalysts under similar fuel-cell operating conditions. It is to be noted that all the rate constants on the Pt<sub>3</sub>Co/C catalyst were larger than the corresponding rate constants on the Pt/C catalysts. During the  $1.0 \text{ V} \rightarrow 0.4 \text{ V}$ process, for example, Pt–Pt bond breaking on the Pt<sub>3</sub>Co/C catalyst was 1.5-fold to that on the Pt/C catalyst. The rate constant of Pt–O bond breaking ( $k'_{Pt-O}$ ), which is highly related to the ORR activity of the cathode catalyst, was about 4 times higher on Pt-Co/C ( $0.4 \text{ s}^{-1}$ ) than on Pt/C ( $0.11 \text{ s}^{-1}$ ).

#### Award

ISHIGURO, Nozomu; Student Presentation Award (JXAFS 14).

There were also a significant difference in the rate constant of Pt–Pt bond re-formation ( $k'_{Pt-Pt}$ ), on Pt<sub>3</sub>Co/C (0.3 s<sup>-1</sup>) and Pt/C (0.078 s<sup>-1</sup>), which may be relevant for realizing better catalyst durability of the Pt<sub>3</sub>Co/C catalyst.

The irreversible oxidation of the Pt catalysts depends on the difference in rate between oxidation and reduction of the Pt catalysts during voltage cycling. The repeated cycling on oxidized Pt species which could not be recovered causes extensive dissolution of the Pt cathode catalyst. Pt<sub>3</sub>Co alloy systems facilitate reversible oxidation/reduction in the voltagecycling processes. The structural kinetics determined from the in situ time-resolved XAFS revealed that all the rate constants of the surface events on the Pt<sub>3</sub>Co/C cathode catalyst were higher than those on Pt/C, particularly the reduction steps involving Pt-O bond breaking. Reversible redox cycles on Pt<sub>3</sub>Co alloys are suggested as a key factor in the superior performance of Pt<sub>3</sub>Co alloy cathode catalysts in PEFCs. Understanding of the fundamental issues of the structural kinetics of the catalyst surface events by in situ time-resolved XAFS establishes new boundaries for the regulation and operation of fuel cells.

**Table 1.** Rate constants for Pt<sub>3</sub>Co/C and Pt/C MEAs in voltagecycling processes by operando time-resolved XAFS.

Process		Rate constant /s <sup>-1</sup>	
		Pt <sub>3</sub> Co/C	Pt/C
$0.4 \rightarrow$ 1.0 V	XANES white-line height	$0.12\pm0.02$	$0.073 \pm 0.001$
	CN (Pt-Pt)	$0.13\pm0.03$	$0.088 \pm 0.008$
	CN (Pt–Co)	No change	-
	CN (Pt–O)	$0.10\pm0.03$	$0.076\pm0.009$
	charge in the fuel	$2.86\pm0.04$	$1.84\pm0.02$
	cell	$0.258\pm0.003$	$0.167\pm0.001$
$1.0 \rightarrow 0.4 \text{ V}$	XANES white-line height	$0.24\pm0.05$	$0.14\pm0.03$
	CN (Pt-Pt)	$0.3\pm0.1$	$0.078\pm0.009$
	CN (Pt–Co)	No change	-
	CN (Pt–O)	$0.4\pm0.2$	$0.11\pm0.02$
	charge in the fuel	$3.68\pm0.03$	$2.16\pm0.01$
	cell	$0.484\pm0.002$	$0.259 \pm 0.001$

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† carrying out graduate research on Cooperative Education Program of IMS with The University of Tokyo

<sup>\*</sup> IMS International Collaboration Program

### Studies of Field-Effect-Transistor Based on Highly-Correlated Molecular Materials

#### Department of Materials Molecular Science Division of Electronic Properties



YAMAMOTO, Hiroshi URUICHI, Mikio SHIMIZU, Atsuko Professor (April, 2012–) Technical Associate Secretary

Field effect transistor (FET) with organic molecules as a channel material is under intense studies because of its application possibilities such as flexible, printable, and large-area electronic devices. Despite such a thorough investigation on single-component neutral molecules to pursuit high performance (say, high mobility etc.) in FET uses, few studies are known for compound organic semiconductor- based transistors. We have been exploiting FET devices with charged molecular materials, namely cation-radical salts of electrondonating molecules (donors) such as BEDT-TTF and anionradical salts of electron accepting molecules (acceptors) such as Ni(dmit)<sub>2</sub>. Among these materials, our focus concentrates in Mott-insulators in which Coulomb interaction among carriers blocks metallic transport. In this highly correlated situation of charge carriers, Mott insulator stays in a fragile semiconducting state, where carrier injection, chemical and/or physical pressure (or strain), and thermal fluctuation can drive it into a metallic state by phase transition (i.e. Mott transition: Figure 1). Since the carrier density (band-filling) of a FET interface can be finely modulated by an electrostatic field from the gate electrode, it is anticipated that the Mott-insulating state can be switched to a metallic state by a field effect.

Indeed, our previous work showed such an insulator-to-metal transition in terms of the band structure, and these results provide not only a possibility of application use of organic Mott-FET but also a significant insight into the mechanism of Mott transition itself. We are now expanding our research target both to the superconducting transition of Mott-FET and to its room-temperature operation.

(BEDT-TTF = bis(ethylenedithio)tetrathiafulvalene, dmit = 1,3-dithiole-2-thione-4,5-dithiolate)

### 1. Field-Induced Superconductivity in an Organic Mott-FET

 $\kappa$ -(BEDT-TTF)Cu[N(CN)<sub>2</sub>]Br (κ-Br) is an organic superconductor whose electronic state is Mott-insulating at roomtemperature but turns into metallic at low temperature through a crossover around 50–100 K, possibly because of an increase of bandwidth upon thermal contraction. In our previous works, a tensile strain altered its ground state into a Mott-insulating state, when its thin (100–300 nm) crystal is laminated on top of SiO<sub>2</sub>/Si<sup>++</sup> substrate and cooled down to low temperature.



Figure 1. Concept of a Mott-insulator (center) along with its bandfilling-controlled (left) and bandwidth-controlled (right) Mott transitions. When the band is half-filled and band energy is small, the carriers are localized because of on-site Coulomb interaction (U). This situation is broken down either by changing the carrier concentration or by compressing the lattice. In many cases, the Mott transitions are associated with superconductivity at the boundary between metallic and insulating phases.

Although the electronic state at low temperature became completely insulating in this experiment because of the very small thermal expansion coefficient of Si substrate (2 ppm/K), one can anticipate from the T-P (temperature vs. pressure) phase diagram that somehow mixed electronic state between superconducting and Mott-insulating states can be realized when the tensile strain is much weaker. To achieve such a mixed state in the device, where phase-separation occurs between superconducting and Mott-insulating states, we have chosen Nb-doped SrTiO<sub>3</sub> as a back-gate substrate because of its larger thermal expansion coefficient (ca. 10 ppm/K) than Si. An aluminum oxide layer was grown by atomic layer deposition technique to form a gate dielectric on the substrate. After lamination process of  $\kappa$ -Br on the substrate, the Mott-FET device which showed a weakly insulating behavior at low temperature was fabricated (Figure 2). Upon applying a positive gate voltage, however, the resistivity goes down and weakly metallic behavior was observed at  $V_{\rm G} > 2$  V. By further increasing the gate voltage up to 8 V, the device showed a sudden drop of resistivity around 5 K, which can be attributed to superconductivity. Taking account of bistable IV characteristics observed in the low resistance region, the above transition can be understood as a percolation transition of superconducting islands that is induced by the electrostatic doping of electrons. The transition temperature increases as the gate voltage rises and saturates around  $V_{\rm G} = 11$  V. The above result is the first example of field-induced superconductivity in organic materials, and can be utilized for uncovering a phase diagram of organic Mott system in the band filling controlled regime.



**Figure 2.** The device structure and temperature dependency of device resistance under various gate voltages. The device comprises  $\kappa$ -Br thin-layer single crystal and Al<sub>2</sub>O<sub>3</sub>/SrTiO<sub>3</sub> substrate. By applying a positive gate voltage, the resistance of the device starts decreasing, while a negative gate voltage results in an increase of the resistance. Although superconducting (SC) and Mott-insulating (MI) phases are separated in the device, the four-terminal resistance drops to zero when the Josephson junction (JJ) network forms a pathway to shunt the voltage terminals. This kind of percolation transition is clearly observed in the right panel, where the resistance drops at 5 K with a gate voltage of 9 V.

#### 2. Operation of Organic Mott-FET at Higher Temperature

In order to operate the organic Mott-FET at higher temperature, hopefully at room temperature, it is necessary to realize much larger Mott gap energy and thinner crystal thickness. For this purpose, we are examining several strategies in parallel.

 $\kappa$ -(BEDT-TTF)Cu[N(CN)<sub>2</sub>]Cl ( $\kappa$ -Cl) is a Mott-insulator with larger Mott gap than  $\kappa$ -Br and is therefore suitable for examining the influence of an increase in Mott-gap energy on the device performance. For example, it has higher resistance than  $\kappa$ -Br and exhibit better ON/OFF ratio. In addition, it always shows ambipolar behavior, probably because of its clean surface (Figure 3). This situation allows us to analyze the critical exponents of filling-controlled Mott transition, which will give us important information about ON/OFF efficiency of the device. Because it is also possible to control the Mott gap energy by applying a mechanical strain,  $\kappa$ -Cl device provides an ideal platform to expand Mott-FET strategy to higher temperature.

Another candidate of the channel material for Mott-FET that can be used at room temperature is (BEDT-TTF)(TCNQ). We are exploiting this material both in a crystalline and amorphous form to find a better FET setup for simple but highly efficient device at room temperature. (TCNQ = tetracyanoquinonedimethane)



**Figure 3.** Gate voltage dependence of the resistance for  $\kappa$ -Cl based Mott-FET. Because this device clearly shows ambipolar behavior due to a smaller number of mid-gap trap states, it is possible to exactly define the gate voltage for charge neutrality point ( $V_{\text{Mott}}$ ). This fact helps determining the critical exponent for Mott transition, which is important both for theoretical investigation and for analyzing device parameters such as subthreshold swing.

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### Magnetic Resonance Studies for Molecular-Based Conductors

#### Department of Materials Molecular Science Division of Electronic Properties



NAKAMURA, Toshikazu FURUKAWA, Ko TAKAHASHI, Seiya ABE, Hitomi Associate Professor Assistant Professor Graduate Student\* Secretary

Magnetic resonance measurements are advantageous for studying fundamental electronic properties and for understanding the detailed electronic structures of molecular based compounds. Developing an understanding of the electronic phases and functionality of these materials enables us to perform systematic investigations of low-dimensional, highlycorrelated electron systems and functional materials. Competition between the electronic phases in molecular-based conductors has attracted much attention. The investigations of such electronic phases by magnetic resonance measurements are important to understanding unsolved fundamental problems in the field of solid state physics, and to explore novel functionalities in the field of material science.

In this study, we performed broad-line NMR and ESR measurements on molecular-based conductors to understand electron spin dynamics and functionality in low-temperature electronic phases.

#### 1. <sup>13</sup>C NMR Study of the Magnetic Properties of the Quasi-One-Dimensional Conductor, (TMTTF)<sub>2</sub>SbF<sub>6</sub>

Magnetic properties in the quasi-one-dimensional organic salt, (TMTTF)<sub>2</sub>SbF<sub>6</sub>, where TMTTF is tetramethyltetrathiafulvalene, are investigated by <sup>13</sup>C NMR under pressures. Antiferromagnetic phase transition at ambient pressure (AFI) is confirmed. Charge ordering is suppressed by pressure and is not observed under 8 kbar. For 5 < P < 20 kbar, a sharp spectrum and the rapid decrease of the spin-lattice relaxation rate  $1/T_1$  were observed below about 4 K, which is attributed to a spin-gap transition. Above 20 kbar, an extremely broadened spectrum and a critical increase of  $1/T_1$  were observed. This indicates that the system enters into another antiferromagnetic phase (AFII) under pressure. The slope of the antiferromagnetic phase transition temperature  $T_{AFII}$ ,  $dT_{AFII}/dP$ , is positive, while  $T_{AFI}$  decreases with pressure. The magnetic moment is weakly incommensurate with the lattice at 30 kbar.



**Figure 1**. Temperature dependence of the <sup>13</sup>C-NMR spin-lattice relaxation rate,  $1/T_1$ , at six different pressures.

#### 2. Time-Resolved ESR Study for Photoconductive Mechanism of Photoconductive TTF-Derivatives

The spin dynamics of photoconductive ethylene-dithiolotetrathiafulvalene (EDT-TTF) containing 1,3-benzothiazole (BTA) was examined for froze solution and powdered samples using time-resolved electron spin resonance (TR-ESR) spectroscopy. While the TR-ESR signal of a frozen solution sample under visible excitation were attributed to the excited triplet state  $T_1$ , that of the powdered indicates a charge-separated state. It is a characteristic feature of molecular assembled system.



**Figure 2.** (a) Molecular structure of single-molecule-type photoconductive materials, TTF and EDT-TTF derivatives. (b) 2D pulsed TR-ESR spectra for powder sample of photoconductive EDT-TTF containing BTA. The normal axis represents the ESR signal intensity. The color scale denotes the signal intensity. Positive and negative values indicate the absorption and emission of microwaves, respectively.

#### 3. Physical Properties of a Molecular Conductor (BEDT-TTF)<sub>2</sub>I<sub>3</sub> Nanohybridized with Silica Nanoparticles by Dry Grinding

The dry grinding of a mixture of bis(ethylenedithio)tetrathiafulvalene (BEDT-TTF) and silica nanoparticles has produced powdery (BEDT-TTF)-silica nanocomposites. The (BEDT-TTF)-silica nanocomposites are readily doped with iodine in hexane dispersion to give powdery nanocomposites of (BEDT-TTF)<sub>2</sub>I<sub>3</sub>-silica. XRD and TEM measurements suggest that  $(BEDT-TTF)_2I_3$  in the nanocomposite exists as shell layers of core-shell-type nanoparticles and as nanometre-sized crystals incorporated into hollow sites of aggregated silica nanoparticles. Magnetic susceptibility measurements reveal that the nanocomposites accompanied a large number of Curie spins attributable to surface molecules of the core-shell-type nanoparticles. The nanocomposites show a magnetic susceptibility change corresponding to the metal-insulator transition of  $\alpha$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> in a broad temperature range of 110–140 K, which is attributed to the properties of the nanocrystalline components. Doping in diethyl ether dispersion leads to higher amounts of the nanocrystalline component being obtained. The doping of (BEDT-TTF)-silica nanocomposites by dry grinding produces a paramagnetic powder containing amorphous (BEDT-TTF)<sub>2</sub>I<sub>3</sub>, which possesses a Curie spin concentration of 50%. The effects of annealing on these nanocomposites are investigated. The electrical conductivity of the compaction pellets of (BEDT-TTF)–silica nanocomposites is enhanced by iodine doping to reach approximately  $10^{-6}$  S cm<sup>-1</sup>, but the value is much lower than that of the bulk crystals  $(10^1 \text{ S cm}^{-1})$ .

#### 4. Electronic Properties and <sup>1</sup>H Dynamics in Self-Doped Organic Conductor (TTFCOO)[(NH<sub>4</sub><sup>+1</sup>)<sub>1-x</sub>(NH<sub>3</sub>)<sub>x</sub>]

<sup>1</sup>H-NMR and High-Field ESR measurements were carried out for self-doped type organic conductor, (TTFCOO)  $[(NH_4^{+1})_{1-x}(NH_3)_x]$  and its deuterated salt. (TTFCOO)  $[(NH_4^{+1})_{1-x}(NH_3)_x]$  and related salts are TTF-based self-doped hydrogen-bonding conductor developed by NIMS group. The TTF-skeleton stacks to form a one-dimensional column. The pristine TTFCOONH<sub>4</sub> molecule is closed-shell. Self-doped type carrier is generated by substitution of the end group of  $(NH_3^0)$  with  $(NH_4^{+1})$ , which is regarded as a charge-reservoir. Actually, a considerable concentration of NH<sub>3</sub> molecules has been confirmed by X-ray photoelectron spectroscopy data. According to the static magnetic susceptibility and high-field ESR measurements, we found that the observed spin was distributed on quasi-hole-like TTF skeletons, and that the TTFCOO mainframe partially becomes a neutral radical. The spin susceptibility was well fitted with a Curie-Weiss term and an activation-type term. The activation-type term was dominant at high temperatures.

In order to clarify the <sup>1</sup>H dynamics, we carried out <sup>1</sup>H-NMR measurements for pristine and deuterated salts. The <sup>1</sup>H-NMR spin-lattice relaxation rate, <sup>1</sup>H- $T_1^{-1}$ , of indicates a characteristic temperature and frequency dependence. A pronounced peak observed at around 200 K is possibly due to molecular motion. Since the relaxation curve is not single component below 200 K, we assumed the <sup>1</sup>H- $T_1^{-1}$  from initial slope (weighted average). Below 100 K <sup>1</sup>H- $T_1^{-1}$  follows  $T^{0.5}$ , and remains frequency dependence, indicating a 1D spin-diffusion type relaxation. However, the Fermi surface of (TTFCOO)[ (NH<sub>4</sub><sup>+1</sup>)<sub>1-x</sub>(NH<sub>3</sub>)<sub>x</sub>] is considered to be 2D, according to the tight-binding band calculation. This discrepancy is an open question at present.

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### **Design of Porous Polymer Frameworks**

#### Department of Materials Molecular Science Division of Molecular Functions



JIANG, Donglin NAGAI, Atsushi LIU, Xiaoming GUO, Zhaoqi LIN, Jianbin XU, Yanhong DING, Xuesong JIN, Shangbin CHEN, Xiong FENG, Xiao CHEN, Long XU, Hong WU, Yang SUZUKI, Hiroko

Covalent organic frameworks (COFs) are crystalline porous polymers that allow organization of organic units into longrange-ordered two and three-dimensional periodicities.<sup>1)</sup> COFs are intriguing frameworks because they allow a new degree of control of porosity, composition and component positions. In our group, we have pioneered the synthesis of  $\pi$ -electronic 2D COFs with various  $\pi$ -units.<sup>1–6</sup>

Conjugated microporous polymers (CMPs) are a class of porous organic frameworks consisting of extended  $\pi$ -conjugation and inherent nanopores. The  $\pi$ -conjugation skeletons together with the well-defined nanopores enable the development of novel materials that are completely different from traditional linear conjugated polymers and conventional porous polymeric materials lacking of strong  $\pi$  correlations among the building blocks. We have developed CMPs as a new platform to achieve light-emitting, light-harvesting, catalysis, energy storage, and sensing functions.<sup>7,8)</sup>

#### 1. Pore Surface Engineering of COFs

COFs are a class of important porous materials that enable pre-designable pore size and geometry; however, pore surface engineering in COFs remains challenging. COFs, which offers precise control over the composition, density, and functionalities of organic groups anchored onto the pore walls.<sup>2)</sup> For this purpose, an azide-appended phenylene block was synthesized and utilized as the wall component in a two-component condensation system for the construction of COFs with azideintegrated walls. The two-component system was further developed to a three-component system that consists of azideappended phenylene and 'naked' phenylene walls, and corner blocks, thereby enabling the synthesis of COFs bearing a tunable content of azide units on the walls. The azide units on the COF walls were found to undergo quantitative click reaction with alkynes to form triazole-linked groups on the wall surfaces. Consequently, the combination of a three-component condensation system with click reaction led to the establishment of a protocol that allowed control of both the composiAssociate Professor Assistant Professor Post-Doctoral Fellow Post-Doctoral Fellow Graduate Student Secretary



**Figure 1.** A general strategy for the surface engineering of COFs through the combination of condensation reaction and click chemistry. COFs bearing azide units on the walls are synthesized by the condensation reaction of hexahydroxytriphenylene (HHTP) with azide-appended benzene diboronic acid (N<sub>3</sub>-BDBA) and benzene diboronic acid (BDBA) in a designated molar ratio (X = 0-100%). The azide groups on the COF walls are clicked with alkynes to anchor various organic groups onto the walls of COFs (X%RTrz-COF-5). The density of R surface groups on the walls is determined by the azide content in X%N<sub>3</sub>-COF-5.

tion and density of organic groups in the pores (Figure 1). This methodology was widely useful for the creation of various specific surfaces in COF pores.<sup>2)</sup>

COF-5 is a typical mesoporous COF that consists of triphenylene at the corners and phenylene on the walls of a boronate ester-linked 2D hexagonal skeleton.1 We introduced azide groups to the 2 and 5 positions of the phenylene unit of 1,4-benzenediboronic acid and utilized the azide-appended benzenediboronic acid (N<sub>3</sub>-BDBA) as wall blocks for the condensation reaction with hexahydroxytriphenylene (HHTP) (Figure 1). The resulting hexagonal skeleton is identical to that of COF-5, with triphenylene corners but azide-appended phenylene walls. Such a strategy allows the growth of various specific groups on the surface because the azide units in the skeletons are reactive groups that can click with various alkynes for pre-designable functionalization. The content of azide units on the pore walls can be tuned using a three-component condensation system using HHTP as the corners and a mixture of N<sub>3</sub>-BDAB and BDBA as the walls, because the azide units in N<sub>3</sub>-BDAB did not affect the reactivity of BDBA in the condensation.

Azide units are well established for click reactions with alkynes to introduce functional groups via 1,2,3-triazole rings (Trz). Reaction of azide units on the walls of  $X\%N_3$ -COF-5 (X = 5, 25, 50, 75, and 100) in a typical click reaction with 2-propynyl acetate in anhydrous N,N-dimethylacetamide (DMAc) at 50 °C with CuI present as a catalyst gives the corresponding X%AcTrz-COF-5 in quantitative yields. The present strategy can be further extended to the introduction of various specific groups onto the walls of COFs. We selected a paraffinic unit, aromatic hydrocarbon, and ester groups as typical examples to react with 50%N<sub>3</sub>-COF-5. The click reactions of 50%N<sub>3</sub>-COF-5 with 1-hexyne, 3-phenyl-1-propyne, and methyl propiolate were quantitatively achieved and led to the integration of *n*-butyl, benzyl, and methyl ester units to the walls of COF-5.

The present method for the pore surface engineering of COFs is simple and universal. This approach is high throughput due to the high efficiency condensation reaction and the quantitative click reaction. The most significant feature is that the combination of condensation and click reactions enables the molecular design of pore surface with controlled composition, component, and density. Therefore, this approach provides a new aspect in the chemistry of these well-defined porous materials, and provides a means to tailor pores for specific purposes and applications.

#### 2. Highly Luminescent CMPs with Facilitated Exciton Migration and Improved Light-Emitting Activity

Conjugated polymers play a vital role in lasing, lightemitting diodes, flexible transistors, and solar cells. Owing to their rigid conformation, they have a high tendency to aggregate in solution and the solid state. Such aggregation leads to the dissipation of excitation energy and ultimately limits their utility as light-emitting motifs.

Herein we report a new strategy for the construction of light-emitting conjugated polymers based on conjugated microporous architectures.<sup>7)</sup> We have developed a highly luminescent CMP with tetrakisphenylethene(TPE) as building block (Figure 2, TPE-CMP). Owing to the crosslinking nature of CMP, TPE-CMP can suppress the rotation of TPE units, thus allowing the high luminescence in both solution and solid states. As a result, a positive "CMP effect" was observed, *i.e.* the interweaving CMP architecture can promote  $\pi$  conjugation, facilitates excitation migration, and improves luminescence activity. Changing the reaction time allows for the synthesis of a series of CMPs with different size and absorption band. As the reaction time increased, the TPE-CMP

particles become larger to give an increased surface area and display red-shifted electronic absorption band accordingly.



Figure 2. Structure of TPE-CMP.

These results the propagation of network structure takes place with the reaction time. The fluorescence quantum yield could reach 40%. In contrast, linear TPE polymer showed a low quantum yield of only 0.65% at otherwise identical condition. Interesting, owing to the crosslinking nature of the network, TPE-CMP is highly luminescent in various solvents such as methanol, dioxane, tetrohydorfuran, dichloromethane, chloroform, hexane, dimethylforamide, benzene, and water while the linear polymer analogue is almost nonemissive in solvents such as dichloromethane, chloroform, and tetrehydrofuran. Furthermore, TPE-CMP emits strong luminescence at solid state with a similar emission band as that in solutions. Fluorescence anisotropy measurements suggest that the exciton can migrate over the three-dimensional CMP skeletons. These positive CMP effects suggest that the CMP architecture provides a new platform for the design of highly luminescent materials

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### Solid State NMR for Molecular Science

Department of Materials Molecular Science Division of Molecular Functions



NISHIMURA, Katsuyuki IIJIMA, Takahiro TANIO, Michikazu Associate Professor Assistant Professor IMS Research Assistant Professor

We are working on methodology and hardware developments of solid state NMR and structural biology and materials molecular science. In the following, we show studies of peripheral membrane proteins and inorganic compounds based on NMR.

## 1. NMR and Native-PAGE Analyses of the Phospholipase C- $\delta 1$ Pleckstrin Homology Domain

Phospholipase C (PLC) binds to phosphatidylinositol 4,5-bisphosphate (PIP<sub>2</sub>) in the cell membrane through the pleckstrin homology (PH) domain, and hydrolyzes PIP<sub>2</sub> to produce two second messengers, diacylglycerol and inositol 1,4,5-triphosphate (IP<sub>3</sub>), by the catalytic domain. The PLC- $\delta$ 1 PH domain has a characteristic short  $\alpha$ -helix ( $\alpha$ 2) from residues 82-87 (Figure 1). Solid-state <sup>13</sup>C NMR studies of the PLC- $\delta$ 1 PH domain suggested that the  $\alpha$ 2-helix non-specifically interacts with the hydrophobic layer of the membrane due to the membrane localization of the protein.<sup>1</sup>

In this study, the contributions of the  $\alpha$ 2-helix toward the IP<sub>3</sub> binding activity and thermal stability of the PLC- $\delta$ 1 PH domain were therefore investigated by using NMR and Native-PAGE methods.

Native-PAGE analyses indicated that IP<sub>3</sub> binding to the PLC-81 PH domain results in a drastic migration shift on the gel and in an increased thermal stability. In addition, we found that disruption of the  $\alpha$ -helical conformation by replacement of Lys-86 with proline resulted in reduced affinity for IP<sub>3</sub> and in thermal destabilization of the IP<sub>3</sub>-binding state. Although the mutant protein with replacement of Lys-86 with alanine showed a slight reduction in thermal stability, the IP<sub>3</sub>-binding affinity was similar to that of the wild-type protein. Replacement of Phe-87 with alanine, but not with tyrosine, also resulted in reduced affinity for IP<sub>3</sub> and in thermal instability. These results indicated that the helical conformation of the  $\alpha$ 2-helix and the phenyl ring of Phe-87 play important roles in the IP<sub>3</sub>-binding activity and thermal stability of the PLC-δ1 PH domain.<sup>2)</sup> The <sup>1</sup>H-<sup>15</sup>N HSQC NMR study of the selectively  $[\alpha^{-15}N]$ Lys-labeled PLC- $\delta$ 1 PH domain indicated that



**Figure 1.** Crystal structure of the PLC- $\delta$ 1 PH domain complexed with IP<sub>3</sub> (PDB ID, 1MAI).

IP<sub>3</sub> binding induces chemical shift displacement of all lysine signals (Figure 2). Interestingly, among those signals, only line width of  $\alpha$ -<sup>15</sup>N signal from Lys-86 located on the  $\alpha$ 2-helix markedly changed due to IP<sub>3</sub> binding (Figure 2, inset): in the ligand-free form, the  $\alpha$ -<sup>15</sup>N signal line width of Lys-86 was about two-fold broader than those of the other lysine signals in the one-dimensional <sup>15</sup>N projection of the <sup>1</sup>H-<sup>15</sup>N HSQC NMR spectrum, and the addition of IP3 resulted in line narrowing, and that the broad line width of the  $[\alpha^{-15}N]$ Lys-86 signal converts to a narrow line width with addition of IP<sub>3</sub> (Figure 2). These findings provide evidence that the conformation and/or dynamics of the  $\alpha$ 2-helix couple with the ligandbinding activity of the PLC-81 PH domain. These findings suggested that the conformational changes of the  $\alpha$ 2-helix induced by membrane binding result in conversion of the stereospecific ligand-binding site to a weak-affinity site and in protein instability. Based on these results, we propose a new affinity regulation mechanism in which the ligand stereospecificity of the PLC-81 PH domain is significantly reduced due to protein structural changes by the membrane binding, and non-specific membrane binding or insertion is therefore important for stable anchoring to the lipid membrane. These findings also suggest that the ligand stereospecificity of the protein mainly contributes to searching the membrane containing PIP<sub>2</sub> and that stable anchoring to the lipid membrane is mainly achieved by non-specific membrane binding or insertion rather than PIP<sub>2</sub> binding. This mechanism also explains the lower affinity of the protein to lipid bilayers than IP<sub>3</sub>, which is an essential property for feedback control of catalytic reaction of PLC- $\delta$ 1 with respect to PIP<sub>2</sub>.



**Figure 2.** Solution NMR analyses of the  $[\alpha^{-15}N]$ Lys-labeled PLC- $\delta$ 1 PH domain. The <sup>1</sup>H-<sup>15</sup>N HSQC NMR spectra without (black) and with (red) IP<sub>3</sub>. (inset) The one-dimensional <sup>15</sup>N projections of the HSQC NMR spectra around the Lys-86 signal under the absence (black) and presence (red) of IP<sub>3</sub>.

### 2. Solid State <sup>95</sup>Mo NMR of Paramagnetic Crystals of Polyoxomolybdates

For a group 6 element of molybdenum, there exist stable compounds with all valences from Mo(0) to Mo(VI). Among them, Mo(V) species are included in materials such as E-Keggin anion and nanosized oxides with ring-, tube- and ball-structure. Localization of the d<sup>1</sup> electron in Mo(V) has attracted attention in terms of molecular design and solid state physics like optic, electric and magnetic properties. Since <sup>95</sup>Mo is a quadrupole nuclei of spin I = 5/2 with the small gyromagnetic ratio  $\gamma = -1.743 \times 10^7$  rad s<sup>-1</sup> T<sup>-1</sup>, it is difficult to measure <sup>95</sup>Mo NMR spectra in solid state in general. Recently, we have reported solid-state <sup>95</sup>Mo NMR spectra of mixed balanced polyoxomolybdates(V, VI) with the localized or delocalized d<sup>1</sup> electrons measured by enhancing sensitivity and resolution using a high-field magnet.<sup>3)</sup> Furthermore, we clarified a disorder structure of [PMo12O36(OH)4{La(H2O)2.75Cl1.25}4]27H2O with the  $\varepsilon$ -Keggin {Mo<sub>12</sub>} core by <sup>95</sup>Mo NMR in solids.

In this study, we investigated a paramagnetic polyoxomolybdate of  $[Mo_{12}O_{30}(OH)_{10}H_2\{Ni(H_2O)_3\}_4]$  (hereafter, abbreviated as  $\{Mo_{12}\}(Ni)$ ). Although the core of this crystal is also  $\varepsilon$ -Keggin  $\{Mo_{12}\}$ , it is caped with four paramagnetic Ni(II)(H<sub>2</sub>O)<sub>3</sub>. In order to examine coupling constants appearing in solid state NMR, we measured high-field <sup>95</sup>Mo NMR in solids, simulated the obtained spectra numerically, and performed DFT calculation.

Figures 3(i-a) and 3(ii-a) show the <sup>95</sup>Mo NMR static

spectra of {Mo<sub>12</sub>}(Ni) measured under 21.8 and 11.7 T, respectively. Considerably broadened spectra with breadth of several thousands of ppm were obtained under both magnetic fields. In the spectral simulation, in addition to the quadrupole and chemical shift interactions, the hyperfine interaction between <sup>95</sup>Mo nuclei and unpaired electron spins in paramagnetic Ni(II) ions were considered as internal interactions. The simulated spectra shown in Figures 3(ii-a) and 3(ii-b) represented experimental spectra well. Since the broadening due to the quadrupole and anisotropic chemical shift interactions were quite large, the effect of the anisotropic hyperfine interaction on the spectral lineshape was small. In order to investigate the contribution from the isotropic hyperfine interaction, the <sup>95</sup>Mo NMR spectrum was measured at 173 K under 11.7 T (Figure 3(i-c)). Notable difference of the spectra between 301 and 173 K was not found. By lowering of temperature, although the line width was slightly broadened owing to the increase of magnetization of the electron spins, we needed not consider temperature dependence of the isotropic shift in the simulation.

The coupling constants for the internal spin interactions were also obtained from DFT calculation. Although the chemical shift estimated by DFT was larger than that by the spectral simulation, the trend such as a considerably large anisotropic chemical shift agreed between DFT and spectral simulation. Also, a spin density localized around the Ni(II) ions was obtained by DFT calculation, which agreed with small temperature-dependence of the whole shift of the spectra.



**Figure 3.**  $^{95}$ Mo NMR static spectra of {Mo<sub>12</sub>}(Ni) under the magnetic field of (a) 21.8 and (b, c) 11.7 T. (b) and (c) show the spectra at 301 and 173 K, respectively. (i) and (ii) show the observed and simulated spectra, respectively. The solid and broken lines in (ii) show the theoretical curves with and without considering the hyperfine interaction, respectively.

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### **Organic Solar Cells**

Research Center for Molecular Scale Nanoscience Division of Molecular Nanoscience



HIRAMOTO, Masahiro KAJI, Toshihiko NAKAO, Satoru SHINMURA, Yusuke KUBO, Masayuki YOKOYAMA, Kazuya YOSHIOKA, Tadashi ISHIYAMA, Norihiro SUGIHARA, Hidemi Professor Assistant Professor Post-Doctoral Fellow Research Fellow Research Fellow Research Fellow Graduate Student Secretary

Organic solar cell is recognized as a next generation solar cell. In 2009, we started CREST Project; "Bandgap Science for Organic Solar Cells." Target of this project is 15% efficiency of organic solar cells by establishing bandgap science for organic semiconductors, which is equivalent to that for silicon semiconductor.

Recently, we have established *pn*-control technique for organic semiconductors (Topic 1) and applied to the fabrication of the ohmic organic/metal contacts (Topic 2) and of the tandem cell formed in the codeposited film only by doping (Topic 3).

#### 1. *pn*-Control and *pn*-Homojunction Formation in Single $C_{60}$ and $H_2Pc$ Films<sup>1–3)</sup>

Fullerene ( $C_{60}$ ) and phthalocyanines are typical components in small-molecular-type organic photovoltaic cells. Based on experience with inorganic solar cells, to create a built-in potential, *pn*-control of highly purified organic semiconductors by doping is required.

In this study, complete *pn*-control and *pn*-homojunction formation were demonstrated for single  $C_{60}$  films<sup>1,2)</sup> and for single metal-free phthalocyanine (H<sub>2</sub>Pc) films.<sup>3)</sup> Cesium carbonate (Cs<sub>2</sub>CO<sub>3</sub>) and molybdenum oxide (MoO<sub>3</sub>) were used



**Figure 1.** Energy diagram of a H<sub>2</sub>Pc film. The broken lines show the energy levels of  $E_{\rm F}$  for undoped (black), MoO<sub>3</sub>-doped (red) and Cs<sub>2</sub>CO<sub>3</sub>-doped (blue) H<sub>2</sub>Pc films.

as donor and acceptor dopants, respectively. Doping was performed by the coevaporation. Here, results on  $H_2Pc$  films are summarized.

The Fermi level ( $E_F$ ) of H<sub>2</sub>Pc, located at the center of the bandgap (4.4 eV), is shifted to 3.8 eV, close to the conduction band (3.5 eV), by Cs<sub>2</sub>CO<sub>3</sub> doping (5,000 ppm) and shifted to 4.9 eV, close to the valence band (5.1 eV), by MoO<sub>3</sub> doping (5,000 ppm) under oxygen free conditions (Figure 1). Formation of *n*- and *p*-type Schottky junctions and *pn*-homojunctions in single H<sub>2</sub>Pc films, confirmed by their photovoltaic properties, clearly demonstrates the formation of *n*- and *p*-type H<sub>2</sub>Pc.

Moreover, the band bending can be also mapped in real scale based on the Kelvin probe measurements and the widths of depletion regions for *n*-type (Figure 2(a)), *p*-type (Figure 2(b)) Schottky junctions and *pn*-homojunction (Figure 2(c)) were determined to 23, 30, and 35 nm, respectively.

Conventionally,  $H_2Pc$  and  $C_{60}$  have been regarded as inherent *p*- and *n*-type semiconductors, respectively. However, *n*-H<sub>2</sub>Pc and *p*-C<sub>60</sub> can be formed by doping. So, it is reasonable that single organic semiconductors can, in general, be controlled to be *n*- or *p*-type, similar to inorganic semiconductors.



**Figure 2.** Energy structures for  $Cs_2CO_3$ -doped (a), MoO\_3-doped (b), and  $Cs_2CO_3$ -/MoO\_3-doped (c) cells in real scale based on the Kelvin probe measurements. Photocarrier generation occurs in the depletion regions (shaded).

## 2. Invertible Organic Photovoltaic Cells with Heavily Doped Organic/Metal Ohmic Contacts<sup>4)</sup>

In this study, we applied *pn*-control to make the two organic/metal contacts in a photovoltaic cell ohmic. Figure 3 shows the energetic structure for  $ITO/n^+$ - $C_{60}$  contact depicted based on the Kelvin probe measurements. 10,000 ppm  $Cs_2CO_3$  was heavily doped in the vicinity of ITO electrode and made  $C_{60} n^+$ -type. Though there is a distinct barrier to electrons with a height of 0.34 eV from the conduction band of  $C_{60}$  to the ITO, since the band bends down steeply within 5 nm of the interface, photogenerated electrons can tunnel through this barrier. Thus, an ohmic contact is formed.

Heavily doped 10-nm-thick  $p^+$  and  $n^+$ -type regions of H<sub>2</sub>Pc and C<sub>60</sub> were formed to facilitate the formation of ohmic contacts at the organic/metal interfaces of two-layered organic photovoltaic cells [ITO/H<sub>2</sub>Pc/C<sub>60</sub>/Ag]. Formation of the ohmic contacts allowed the cells to be invertible [ITO/C<sub>60</sub>/H<sub>2</sub>Pc/Ag] and independent of the type of electrode material used.



**Figure 3.** Energy structure for  $ITO/n^+$ -C<sub>60</sub> contact in real scale. A tunneling ohmic junction for photogenerated electrons is formed.

#### 3. Tandem Organic Photovoltaic Cells Formed in Codeposited Films by Doping<sup>5)</sup>

*pn*-control technique should be extended to co-deposited films since organic solar cells should use co-deposited films to generate significant photocurrent densities based on the dissociation of excitons by the photoinduced electron transfer process.

In this study, tandem organic solar cells connecting two single homojunctions were formed in C<sub>60</sub>:6T ( $\alpha$ -sexithiophene) co-deposited films only by doping with MoO<sub>3</sub> and Cs<sub>2</sub>CO<sub>3</sub> (Figure 4). Doping were performed by "three component co-evaporation."<sup>6)</sup> The single and tandem cells showed the open-circuit voltage of 0.85 and 1.69 V and conversion efficiency of 1.6 and 2.4%, respectively.

Figure 5 shows the energetic structure of tandem cell. Photocurrent is generated in the intrinsic (i) layers in front and back cells. Two single cells are connected by heavily doped  $n^+p^+$ -homojunction. Since the depletion region width of the  $n^+p^+$ -homojunction is extremely thin, *i.e.*, 20 nm, it can act as ohmic interlayer due to the carrier tunneling. Thus, the photovoltage of tandem cell (1.69 V) is doubled compared to that of single cell (0.85 V).



Figure 4. Structure of a tandem cell connecting two homojunction cells, which is made only by doping to  $C_{60}$ :6T codeposited film.



**Figure 5.** Energy structure of a tandem cell in real scale.  $n^+p^+$ -heavily doped homojunction acts as an ohmic interlayer connecting two single cells.

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#### Award

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### Development of Organic Semiconductors for Molecular Thin-Film Devices

#### Research Center for Molecular Scale Nanoscience Division of Molecular Nanoscience



SUZUKI, Toshiyasu SAKAMOTO, Youichi KURODA, Yasuhiro WATANABE, Yoko Associate Professor Assistant Professor Post-Doctoral Fellow Secretary

Organic light-emitting diodes (OLEDs) and organic fieldeffect transistors (OFETs) based on  $\pi$ -conjugated oligomers have been extensively studied as molecular thin-film devices. Organic semiconductors with low injection barriers and high mobilities are required for highly efficient OLEDs and OFETs. Radical cations or anions of an organic semiconductor have to be generated easily at the interface with an electrode (or a dielectric), and holes or electrons must move fast in the semiconducting layer. Compared with organic p-type semiconductors, organic n-type semiconductors for practical use are few and rather difficult to develop. Recently, we found that perfluorinated aromatic compounds are efficient n-type semiconductors for OLEDs and OFETs.

### 1. Selective Synthesis and Crystal Structure of [10]Cycloparaphenylene<sup>1)</sup>

[10]Cycloparaphenylene ([10]CPP) was selectively synthesized in four steps in 13% overall yield from commercially available 4,4'-diiodobiphenyl by using mono-I-Sn exchange, Sn-Pt transmetalation, I-Pd exchange, and subsequent oxidative coupling reactions. The structure of [10]CPP was determined by using single-crystal X-ray analysis. Suitable crystals were obtained by slow vapor diffusion of *n*-hexane into a solution of [10]CPP in CH<sub>2</sub>Cl<sub>2</sub> at room temperature. In the solid state, [10]CPP is slightly distorted to an ellipsoidal structure with major and minor axes of 13.9 and 13.5 Å, respectively. The cavity of [10]CPP is occupied by a hexane molecule, which was highly disordered. Although the  $D_{5h}$ structure with a dihedral angle between two paraphenylene units of  $32^\circ\mbox{--}33^\circ$  was calculated to be the most stable conformer, the structure is closer to a  $D_{2h}$  conformer with alternating triphenylene and biphenylene units. The dihedral angles

between two paraphenylene units were approximately 20° and 45°, respectively. The average  $C_{ipso}-C_{ipso}$ ,  $C_{ipso}-C_{ortho}$ , and  $C_{ortho}-C_{ortho}$  bond lengths are 1.484(1), 1.399(2), and 1.385(9) Å, respectively. [10]CPP molecules pack in a herringbone manner, and there are no significant  $\pi$ - $\pi$  interactions among [10]CPP molecules. A tilted tubular channel structure was observed along the *b* axis. The crystal packing is similar to [9]- and [12]CPPs but different from [6]CPP. The size of the CPP may be important for determining the packing arrangement.



Figure 1. X-ray structure of [10]CPP.



#### Figure 2. Packing structure of [10]CPP.

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### Building Photosynthesis by Artificial Molecules

#### Research Center for Molecular Scale Nanoscience Division of Molecular Nanoscience



NAGATA, Toshi SAKURAI, Rie YUSA, Masaaki MIURA, Takahiro WANATABE, Yoko Associate Professor IMS Fellow Graduate Student Technical Fellow Secretary

The purpose of this project is to build nanomolecular machinery for photosynthesis by use of artificial molecules. The world's most successful molecular machinery for photosynthesis is that of green plants—the two photosystems and related protein complexes. These are composed almost exclusively from organic molecules, plus a small amount of metal elements playing important roles. Inspired by these natural systems, we are trying to build up multimolecular systems that are capable of light-to-chemical energy conversion. At present, our interest is mainly focused on constructing necessary molecular parts and examining their redox and photochemical behavior.

### 1. Photoreduction of Quinones Sensitized by Soluble Phthalocyanines

Organic pigments are useful molecular parts for photosynthesis because of their ease of modification and high tunability. In particular, pigments for utilizing visible lights with long wavelength (>650 nm) are useful because they can complement the existing series of sensitizers for more efficient use of solar light. Phthalocyanines are good candidates as they have very strong absorption bands around 700 nm. However, studies on photochemistry of phthalocyanines are often hampered by the low solubility. To address this problem, we chose phthalocyanines (Pc) having 2,4-di-*t*-butylphenoxy substituents in the periphery. These phthalocyanines are so soluble in many organic solvents (even in hexanes!) that study of photoreactions in solutions can be easily carried out.

The target photoreaction is shown in Figure 1. The rate of reaction was dependent on the redox potentials of quinones, namely the reaction was faster for quinones with higher potential (*i.e.* more easily reduced). This indicates the reaction proceeds *via* photoinduced electron transfer from Pc to quinones. In particular, 2-*t*-butylanthraquinone and 2,6-dibutoxy-



Figure 1. Photoreduction of quinones sensitized by soluble phthalocyanines with various central metals.

anthraquinone did not react at all; these results suggest that the reaction proceeds *via* the triplet state of Pc (<sup>3</sup>Pc) rather than the excited singlet state (<sup>1</sup>Pc\*), because  $\Delta$ G's for the photo-induced electron transfer to these quinones were positive for <sup>3</sup>Pc but negative for <sup>1</sup>Pc\*.

On the other hand, when different thiols were used as reductant, the rate of reaction was dependent on the acidity of thiols rather than the redox potential. Therefore, it is likely that the thiol works first as the proton donor to the anion radical of quinone, and then gives an electron to the cation radical of Pc (or to the neutral semiquinone radical).

When phthalocyanines with different central metal ions were employed, the rates of photoreaction were in the order Zn > Mg > Pd > 2H. Although the triplet quantum yield is higher for Pd than Zn, the photoreaction was very slow with Pd phthalocyanine. This is attributed to the high redox potential of PdPc, causing the photoinduced electron transfer unfavorable ( $\Delta G = +0.17$  eV for 2,5-di-*t*-butylbenzoquinone).

On the basis of these results, a plausible mechanism of this reaction is proposed as shown in Figure 2. Kinetic studies for important steps are currently underway. This study will lay an important cornerstone in utilization of low-energy photons for direct chemical conversions.



Figure 2. A plausible mechanism of photoreduction of quinones. Pc: phthalocyanine.

### 2. Synthesis of Dinuclear Metal Complexes with Ternary Binucleating Ligands

As we already reported,<sup>1)</sup> combination of photochemistry of organic pigments and redox chemistry of transition metal complexes is a promising approach to develop new photoredox chemistry. In particular, use of first transition metals (3d metals) is worth studying, because they can provide crucial bridge between one-electron (photochemical) and two-electron (common organic chemical) processes. Synthesis of 3d metal complexes, however, is not always straightforward, because many 3d metal ions exchange ligands so easily that there are only limited choices of isolable compounds. In this respect, we can learn from biological systems where 3d metal complexes are utilized with a wide structural variation. This is made possible by the surrounding organic ligands (proteins).

Along this context, we prepared two new "ternary" ligands 1 and 2, which consist of two terpyridines and one "N4 bridge" (1,4-bis(2-pyridyl)phthalazine for 1 and 3,6-bis(2pyridyl) pyridazine for 2). They formed stable dinuclear complexes  $[(1)Co_2(\mu-OH)]^{3+}$ ,  $[(1)Ni_2(\mu-Cl)]^{3+}$  and  $[(2)Co_2$  $(\mu$ -OH)]<sup>3+</sup>. The X-ray structures of the first two complexes revealed subtle but characteristic differences (Figure 4). The steric repulsion between the hydrogen atoms at the peripositions of the phthalazine ring and the 3-positions of the neighboring pyridine rings (indicated by arrows in Figure 4) caused twisting of the pyridine rings resulting in local axial chirality, and the stereochemistry of the axes were different between the Co and Ni complexes. In the Co complex, the two axes have the opposite axial stereochemistry, whereas in the Ni complex, the two axes have the same stereochemistry. Such difference was caused by the different ionic radii of the metal centers and by the different bridging anions (OH- and Cl-).

The cyclic voltammogram of  $[(1)Co_2(\mu-OH)]^{3+}$  showed interesting behavior in the presence of acid (Figure 5). When triflic acid was added, the peaks of the Co<sub>2</sub> parts disappeared and new peaks corresponding to a mononuclear Co complex appeared instead. This change was attributed to the (reversible) release of one Co ion upon addition of acid.



Figure 3. The ternary ligands consisting of two terpyridines and one "N4 bridge."



**Figure 4.** The space-filling drawing of the X-ray structures of (a) [(1)  $Co_2(\mu$ -OH)]^{3+} and (b) [(1)Ni\_2(\mu-Cl)]^{3+}.



**Figure 5.** The changes of the CVs of  $[(1)Co_2(\mu-OH)]^{3+}$  with the addition of CF<sub>3</sub>SO<sub>3</sub>H (a) 0 eq., (b) 1.0 eq., (c) 2.0 eq. and (d) after the addition of 5 eq. *i*PrNEt<sub>2</sub>.

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### Chemistry of Bowl-Shaped Aromatic Compounds and Metal Nanocluster Catalysts

#### Research Center for Molecular Scale Nanoscience Division of Molecular Nanoscience



SAKURAI, Hidehiro HIGASHIBAYASHI Shuhei PANDA, Gautam MURUGADOSS, Arumugam TAN. Qi-Tao MORITA, Yuki SUGIISHI, Tsuyuka REZA, A. F. G. Masud ANWER, Muslih SOPHIPHUN, Onsulang KAMONSATIKUL, Choavarit KOO, Jin Young ONOGI, Satoru DHITAL, Raghu Nath KARANJIT, Sangita SHRESTHA, Binod Babu KAEWMATI, Patcharin HAESUWANNAKIJ, Setsiri OKABE, Yuki SCHMIDT, Bernd TOPOLINSKI, Berit NAKANO, Sachiko KIM. Yukimi SASAKI, Tokiyo TANIWAKE, Mayuko HAZAMA, Kozue

Associate Professor Assistant Professor Visiting Scientist; JSPS Invited Fellow Visiting Scientist; JSPS Post-Doctoral Fellow Visiting Scientist; JSPS Post-Doctoral Fellow Post-Doctoral Fellow Post-Doctoral Fellow Visiting Scientist (JASSO Follow-up Research Fellow) Visiting Scientist Visiting Scientist Visiting Scientist Visiting Scientist Graduate Student Graduate Student' Graduate Student Technical Fellow Technical Fellow Secretary Secretary Secretary

Bowl-shaped  $\pi$ -conjugated compounds including partial structures of the fullerenes, which are called "buckybowls," are of importance not only as model compounds of fullerenes but also as their own chemical and physical properties. For example, in solution they show the characteristic dynamic behavior such as bowl-to-bowl inversion. On the other hand, they sometimes favor stacking structure in a concave-convex fashion in the solid state, giving excellent electron conductivity. Furthermore, some buckybowls are conceivable to possess the bowl-chirality if the racemization process, as equal as bowl-to-bowl inversion, is slow enough to be isolated. Very few buckybowls has been achieved for preparation mainly due to their strained structure, and no report on the preparation of chiral bowls has appeared. In this project, we develop the rational route to the various buckybowls with perfect chirality control using the organic synthesis approach.

We also investigate to develop novel catalytic properties of metal nanoclusters. We focus on the following projects: Preparation of size-selective gold nanoclusters supported by hydrophilic polymers and its application to aerobic oxidation catalysts: Synthetic application using metal nanocluster catalyst: Development of designer metal nanocluster catalyst using the highly-functionalized protective polymers.

### 1. Enantioselective Synthesis of a Chiral Nitrogen-Doped Buckybowl<sup>1)</sup>

Bowl-shaped aromatic compounds, buckybowls, constitute a family of curved polycyclic aromatic carbons along with fullerenes and carbon nanotubes (CNTs). Doping of heteroatoms to the carbon frameworks of such aromatic compounds drastically modulates their physical and chemical properties. In contrast to nitrogen-doped azafullerenes or CNTs, synthesis of azabuckybowls, nitrogen-doped buckybowls, remains an unsolved challenging task. Accordingly, we have achieved the first enantioselective synthesis of a chiral azabuckybowl, triazasumanene. X-ray crystallographic analysis confirmed that the doping of nitrogen induces a more curved and deeper bowl structure than in all-carbon buckybowls. As a result of the deeper bowl structure, the activation energy for the bowl inversion (thermal flipping of the bowl structure) reaches an extraordinarily high value (42.2 kcal/mol). Since the bowl inversion corresponds to the racemization process for chiral buckybowls, this high bowl inversion energy leads to very stable chirality of triazasumanene.



Figure 1. (A) Triazasumanene and sumanene; (B) Enantiomers of triazasumanene.



**Figure 2.** Conditions: (a)  $Pd(OAc)_2$ ,  $PPh_3$ ,  $Bu_4NOAc$ ,  $Na_2CO_3$ , molecular sieve 4 Å, 1,4-dioxane, 100 °C, 2 h, 57%; (b) i) 12 M HCl, AcOH, 60 °C, 3 h; ii)  $C_6F_5OP(=O)Ph_2$ , DIPEA, DMF, 0 °C to 60 °C, 59% (two steps); (c) Lawesson's reagent, dichloroethanene, microwave, 160 °C, 40 min, 92%; (d) Trifluoroacetic acid, microwave, 100 °C, 2 h, 88%; (e) MeI,  $K_2CO_3$ , DMF, 30 °C, 3 h, 79%; (f)  $Ph_3CBF_4$ , DTBMP, CH<sub>2</sub>Cl<sub>2</sub>, 25 °C, 3 d, 73%; (g) *m*-CPBA, CH<sub>2</sub>Cl<sub>2</sub>, 25 °C, 5 h, 90%.

#### 2. Trimethylsumanene: Enantioselective Synthesis, Substituent Effect on Bowl Structure, Inversion Energy, and Electron Conductivity<sup>2)</sup>

 $C_3$  symmetric chiral trimethylsumanene was enantioselectively synthesized through Pd-catalyzed syn-selective cyclotrimerization of an enantiomerically pure iodonorbornenone, ring-opening/closing olefin metathesis, and oxidative aromatization where the  $sp^3$  stereogenic center was transmitted to the bowl chirality. Chiral HPLC analysis/ resolution of the derivatives were also achieved. Based on theoretical calculations, the columnar crystal packing structure of sumanene and trimethylsumanene was interpreted as due to attractive electrostatic or CH- $\pi$  interaction. According to the experimental and theoretical studies, the bowl depth and inversion energy were found to increase on methylation for sumanene in contrast to corannulene. Dissimilarities of the effect of methylation on the bowl structure and inversion energy of sumanene and corannulene were ascribed to differences in steric repulsion. A double-well potential model was fitted to the bowl structure-inversion energy correlation of substituted sumanenes, with a small deviation. The effects of various substituents on the sumanene structure and bowl inversion energy were analyzed by density functional theory calculations, and it was shown that the bowl rigidity is controlled by a combination of electronic and steric effects of the substituents. The electron conductivity of trimethylsumanene was investigated by time-resolved microwave conductivity method, compared with that of sumanene.



**Figure 3.** (A) ORTEP drawings of X-ray crystallographic analysis of  $(\pm)$ -**3**, (B) top view, (C) side view, shown at the 50% probability level.

#### 3. Dual Roles of Polyhydroxy Matrices for Homocoupling of Arylboronic Acid Catalysed by Gold Nanoclusters Under Acidic Conditions<sup>3)</sup>

Polyhydroxy biopolymers including chitosan and starch play the dual roles of stabilization of gold clusters as well as activation of arylboronic acids through reversible binding. In acidic pH of 4.57 the electronic state of boron is changed to the corresponding to that at pH 9 inside the vicinity of Au:chit or Au:starch which favours the transmetallation process *i.e.* why homo-coupling product was observed even under acidic solutions.



#### 4. Anomalous Efficacy of Bimetallic Au/Pd Nanoclusters in C–Cl Bond Activation and Formal Metathesis-Type C–B Bond Activation at Room Temperature<sup>4)</sup>

Au/Pd alloy nanoclusters stabilized by poly (N-vinylpyrrolidone) catalyze two different reactions of phenylboronic acid with 4-chlorobenzoic acid at room temperature in a single reaction cycle, cross coupling and metathesis-type homocoupling that is normally inaccessible through conventional catalysis.



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#### Awards

DHITAL, Raghu Nath; Young Scientist Award, International Conference on Advanced Materials and Nanotechnology (2011). DHITAL, Raghu Nath; Young Oral Presentation Award, International Symposium on Catalysis and Fine Chemicals 2011 (2011). TAN, Qi Tao; Outstanding Poster Presentation, International Symposium on Catalysis and Fine Chemicals 2011 (2011).

### Spectroscopic Studies on Electronic Ferroelectricity in Organic Conductors

Department of Materials Molecular Science Division of Electronic Properties

YAMAMOTO, Kaoru

Assistant Professor

Organic conductors exhibit unusual physical properties due to the critical nature of valence electrons fluctuating under localizing and delocalizing regime. Recently, we discovered that an organic superconductor,  $\alpha$ -(BEDT-TTF)<sub>2</sub>I<sub>3</sub> [BEDT-TTF: bis(ethylenedithio)tetrathiafulvalene], induces optical second-harmonic generation (SHG) along with charge-ordering driven by the repulsive electron-electron interactions.<sup>1</sup>) The SHG activity proves that the organic complex induces macroscopic electric polarization by the electronic transition. To understand the fundamental and functional properties of such "electronic ferroelectrics," we are engaged in the material research and development of experimental technique dedicated for low-temperature SHG microscopy of tiny organic crystals.

#### 1. SHG Microscopy of Ferroelectric Organic Conductor Using Hydrostatic Pressure Apparatus with Argon as a Heat Transfer Medium<sup>2)</sup>

To perform nonlinear microscopy at low-temperatures, one has to dissipate the excessive heat from the photo-irradiated spot of the sample without affecting the physical properties. In the present study, we designed a sapphire-anvil cell, which is usually adopted for high-pressure measurements, as a cooling apparatus used along with liquid-argon as heat-transfer medium. The cryogenic performance of the argon-loaded sapphire cell was examined by the observation of SHG of an organic conductor  $\alpha$ '-(BEDT-TTF)<sub>2</sub>IBr<sub>2</sub> at low temperatures.

This  $\alpha$ -type compound is unusually susceptible to mechanical stress. In a previous study, we have embedded the single crystals of the compound into a polymer resin and sandwiched them by sapphire plates, and observed the SHG at low temperatures. Because of the difference in thermal shrinkage between the samples and the sapphire plates, the crystals were fragmented with decreasing temperature. The observed image of SHG showed distinct inhomogeneity; the nonlinear optical signal was generated only from limited areas of the



**Figure 1.** (a) Thin crystals of  $\alpha$ '-(BEDT-TTF)<sub>2</sub>IBr<sub>2</sub> in a sapphireanvil cell loaded with liquid argon. (b) SHG image in the square region of (a) (T = 140 K).

crystals, suggesting that the single crystals were divided into ferroelectric and antiferroelectric regions, presumably due to the mechanical stress.

Figure 1(a) shows the transmission image of the filmy thin crystals of the complex mounted in the sapphire cell loaded with liquid argon. The crystals cooled via the heat transfer medium were not completely intact, but the pieces of the cleaved crystals were much larger than the above mentioned fragmented crystals. As shown by the SHG image [Figure 1(b)], the cleaved crystals generated SHG from their whole region. This fact proves that the compound is a bulk ferro-electric matter, and at the same time, indicates that the heat at the irradiated spot was efficiently dissipated via the heat transfer medium without applying severe mechanical stress to the sample.

The successful SHG observation demonstrates that the cooling technique is helpful for low-temperature measurements of nonlinear optical microscopy requiring high-density photo-excitations.

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### Multifunction Integrated Macromolecules for Molecular-Scale Electronics

#### Research Center for Molecular Scale Nanoscience Division of Molecular Nanoscience



TANAKA, Shoji

Assistant Professor

Recently a single electron tunnel (SET) device has attracted much attention due to the growing demand for ultra-lowpower device. A SET device manipulates an electron by means of one-by-one electron transfer, resulting in ultimately low power consumption. However, for room temperature operation, the size of SET device must be as small as a few nm to overcome the thermal fluctuation problems. The process size of a few nm is out of the range of conventional micro-technology. In this project, to establish an innovative fabrication process for SET device systems, we have been developing step-wise synthetic protocols for monomolecular single-electron tunnel devices (**MOSET**) and their integrated circuits.

#### 1. Universal Temperature Crossover Behavior of Electrical Conductance in a Single Oligothiophene Molecular Wire<sup>1)</sup>

We have observed and analyzed a universal temperature crossover behavior of electrical conductance in a single oligothiophene molecular wire. The crossover between the Arrhenius type temperature dependence at high temperature and the temperature-invariant behavior at low temperature is found at a critical molecular wire length of 5.6 nm, where we found a change from the exponential length dependence to the lengthinvariant behavior. We have derived a scaling function analysis for the origin of the crossover behavior. After assuring that the analysis fits the explanation of the Keldysh Green's function calculation for the temperature dependence, we have applied it to our experimental results and found successfully that our scaling function gives a universal description of the temperature dependence for all over the temperature range.

#### 2. Fabrication of 3- and 4-Terminal Single-Electron Device Structures

To integrate Coulomb islands and tunnel/capacitive junctions in a single molecular skeleton, we have developed a series of molecular building blocks (1-8). These building blocks make possible to fabricate various type of 3- and 4-terminal monomolecular single-electron tunnel device structures (11, 13, 17).



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### Development of Novel Heterocyclic Compounds and Their Molecular Assemblies for Advanced Materials

#### Safety Office



TOMURA, Masaaki

Assistant Professor

Heterocycles containing sulfur and/or nitrogen atoms are useful as components of functional organic materials since heteroatoms in their rings are helpful to stabilize ions or ionradical species. In addition, intermolecular interactions caused by heteroatom contacts can be expected to form unique molecular assemblies. In this project, novel functional organic materials based on various heterocycles were synthesized and their physical and structural properties were investigated.

# 1. Molecular and Crystal Structure of Bis(tetra-*n*-butylammonium) Bis(5,6-dicyanopyrazine-2,3-dithiolato- $\kappa^2 S, S$ ) palladium(II)<sup>1)</sup>

Metal dithiolene complexes have been widely investigated as molecular conductors and superconductors. Several superconductors and single-component molecular metals involving dithiolene complexes have been discovered to date. We have synthesized the title palladium dithiolene complex derived from 2,3-dicyano-5,6-dimercaptopyrazine ligand. The ligand is expected to extend the  $\pi$ -conjugation of the complex resulting in decreased Coulombic repulsion. Intermolecular interactions caused by S…S and S…N heteroatom contacts may increase the dimensionality in the solid state. In the complex the dianion molecule is located on an inversion center. The dianion is a flat molecule with an r.m.s. deviation of 0.034(8) Å of fitted atoms from the least-squares plane. The central Pd atom has a square-planar coordination geometry and the Pd-S distances and the S-Pd-S angle are 2.276(3), 2.294(3) Å and  $89.39(10)^\circ$ , respectively. These values are comparable to those found in bis(tetra-n-butylammonium) bis(4,5-dicyanobenzene-1,2-di- thiolato-S,S')palladium(II) complex. The dianion molecules form a layered structure with an interlayer distance of 6.5 Å. The tetra-*n*-butylammonium cations are inserted between the layers.



#### 2. Photoinduced Electron-Transfer Reaction of $\alpha$ -Bromomethyl-Substituted Benzocyclic $\beta$ -Keto Esters with Amines: Selective Reaction Pathways Depending on the Nature of the Amine Radical Cations

Photoinduced electron-transfer reaction of  $\alpha$ -bromomethyl-substituted benzocyclic  $\beta$ -keto esters with tertiary amines was investigated. Debrominated  $\beta$ -keto esters and ringexpanded  $\gamma$ -keto esters were obtained as major products. On the basis of mechanistic experiments it was concluded that these products are formed via a reaction sequence of selective carbon-bromine bond cleavage and subsequent competitive hydrogen abstraction and Dowd-Beckwith ring-expansion of the resulting primary alkyl radicals. The characteristic product distribution observed for the type of amine used is rationalized on the basis of selective reaction pathways of generated radical intermediates that depend on the nature of the amine radical cations.

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### **Visiting Professors**



#### Visiting Professor ABE, Manabu (from Hiroshima University)

#### Stretch Effects Induced Molecular Strain in Generating Long-Lived Biradicals

Stretch effects induced by two types of molecular strain were examined by quantum chemical calculations, to design persistent multi-radicals such as localized diradicals and oxyallyls. The cooperative molecular strain (Type-1) induced by the spiro[5.5]undecane and bicyclo[2.1.0]cyclopentane structures was found to significantly destabilize in energy the ring-closed compounds of the diradicals, leading to small

energy differences between the diradicals and the  $\sigma$ -bonded compounds. Another stretch effect (Type-2) induced by macrocyclic systems was also found to energetically destabilize the corresponding ring-closed structures of the 1,3-diradicals. The computational studies predict that the two types of stretch effects are quite effective in lowering the energy barriers of the bond-breaking reaction of the ring-closed compounds and in generating long-lived localized diradicals and oxyallyl derivatives.



#### Visiting Professor KATO, Tatsuhisa (from Kyoto University)

#### Studies of Molecular Magnetization of Super-Molecules

The encapsulation of  $C_{60}$  with  $\gamma$ -cyclodextrin ( $\gamma$ -CD) has been attained by a mechanochemical highspeed vibration (HSV) technique. N@C<sub>60</sub>, which can be a good magnetic probe giving the information of position as well as of chemical environment, should be soluble in water for the purpose of biological application. Then the HSV technique was applied to the powder of  $\gamma$ -CD and C<sub>60</sub> containing N@C<sub>60</sub> at 5%.

The obtained molecular complex exhibited the peculiar electron spin resonance (ESR) spectrum of N@C<sub>60</sub> in water.

We designed the flexible intermolecular communications in a simple molecular architecture by using mechanically interlocked supramolecular motives such as catenanes and rotaxanes, in which two or more molecular components are inseparable but their interactions are flexibly convertible. In a dinuclear  $Cu^{2+}$  complex of the four-fold rotaxane, the  $Cu^{2+}$ -porphyrin and the  $Cu^{2+}$ -phthalocyanine were stacked efficiently on one another to afford spin–spin communication. Spin states of the dinuclear complex were switchable between its protonated form (doublet) and deprotonated form (singlet) reversibly.



#### Visiting Professor ASAKURA, Tetsuo (from Tokyo University of Agriculture and Technology)

#### Determination of Molecular Structure with Ultra Fast MAS under High-Field NMR

In single crystal X-ray diffraction analyses of peptides and proteins, it is well-known that the co-ordinates of carbon, nitrogen and oxygen atoms can be obtained in high accuracy, but enough accuracy cannot be obtained for those of hydrogen. Therefore we are trying to determine the accurate <sup>1</sup>H positions by the combination of NMR observation by ultra fast magic angle spinning under high field magnetic field

and accurate <sup>1</sup>H NMR chemical shift calculation. We are applying this novel analytical technique to determine the structures of silk fibroins before and after spinning together with their model peptides. Since such a <sup>1</sup>H information is sensitive to both the intra- and inter -molecular structures, it is especially useful in molecular design of biomaterials with silks.

### **RESEARCH ACTIVITIES**
### **RESEARCH ACTIVITIES** Life and Coordination-Complex Molecular Science

Department of Life and Coordination-Complex Molecular Science is composed of two divisions of biomolecular science, two divisions of coordination-complex molecular science, and one adjunct division. Biomolecular science divisions cover the studies on functions, dynamic structures, and mechanisms for various biomolecules such as sensor proteins, membrane-anchored proteins, biological-clock proteins, metalloproteins, glycoconjugates, and molecular chaperone. Coordination complex divisions aim to develop molecular catalysts and functional metal complexes for transformation of organic molecules, water oxidation and reduction, and molecular materials such as molecular wires. Interdisciplinary alliances in this department aim to create new basic concepts for the molecular and energy conversion through the fundamental science conducted at each division. Professor Koji Tanaka was retired at the end of March, 2012. Professors Shuji Akiyama and Tetsuro Murahashi were appointed as full professors of the Division of Biomolecular Sensing and the Division of Functional Coordination Chemistry, respectively, both in April, 2012.

### Bioinorganic Chemistry of Metal-Containing Sensor Proteins

Department of Life and Coordination-Complex Molecular Science Division of Biomolecular Functions



AONO, Shigetoshi YOSHIOKA, Shiro SAWAI, Hitomi TANIZAWA, Misako Professor Assistant Professor IMS Research Assistant Professor Secretary

The widely studied biological function of heme is to act as a prosthetic group in hemeproteins that show a variety of functions, including oxygen storage and transport, electron transfer, redox catalysis, and sensing of gas molecules. In addition to these functions, a new function of hemeprotein has been found recently, which is a sensor of diatomic gas molecules or redox change.<sup>1)</sup> In these heme-based sensor proteins, the heme acts as the active site for sensing the external signal such as gas molecules and redox change. A new hemeprotein, aldoxime dehydratase (Oxd), also shows a novel biological function of heme, which catalyzes dehydration reaction of organic substrates. Oxd is the first example of hemeproteins that act as a hydro-lyase. Our research interests are focused on the elucidation of the structure-function relationships of these hemeproteins.

### 1. Bacterial Gas Sensor Proteins Using Transition Metal-Containing Prosthetic Groups as Active Sites<sup>2,3)</sup>

Gas sensor proteins are involved in many biological regulatory systems, including transcription, chemotaxis, and other complex physiological processes. These regulatory systems consist of a sensor and regulatory protein, and, if any, signal transduction proteins that transmit the input signal sensed by the sensor protein to regulator proteins. Sensor proteins are the most upstream component in these regulatory systems, and the sensor and regulator proteins can be distinct molecules, or can sometimes exist in the same molecule as sensor and regulator domains. In both cases, the general scheme is as follows for biological regulatory systems by gas sensor proteins. Once a gas sensor protein/domain senses a gas molecule of its physiological effector, a conformational change of the sensor protein/ domain is induced, and then intra- and/or inter-molecular signal transductions proceed to modulate the activity of the regulator protein/domain that is responsible for the regulation of the above biological functions.

HemAT from *Bacillus subtilis* (HemAT-*Bs*) is a hemecontaining  $O_2$  sensor protein that acts as a chemotactic signal transducer. Binding of  $O_2$  to the heme in the sensor domain of HemAT-*Bs* induces a conformational change in the protein matrix, and this is transmitted to a signaling domain. To characterize the specific mechanism of  $O_2$ -dependent conformational changes in HemAT-*Bs*, we investigated time-resolved resonance Raman spectra of the truncated sensor domain and the full-length HemAT-*Bs* upon  $O_2$  and CO dissociation.

A comparison between the O2 and CO complexes provides insights on O<sub>2</sub>/CO discrimination in HemAT-Bs. While no spectral changes upon CO dissociation were observed in our experimental time window between 10 ns and 100 µs, the band position of the stretching mode between the heme iron and the proximal histidine, v(Fe-His), for the O2-dissociated HemAT-Bs was lower than that for the deoxy form on time-resolved resonance Raman spectra. This spectral change specific to O2 dissociation would be associated with the O2/CO discrimination in HemAT-Bs. We also compared the results obtained for the truncated sensor domain and the full-length HemAT-Bs, which showed that the structural dynamics related to O<sub>2</sub> dissociation for the full-length HemAT-Bs are faster than those for the sensor domain HemAT-Bs. This indicates that the heme proximal structural dynamics upon O2 dissociation are coupled with signal transduction in HemAT-Bs.

We have also studied the protein dynamics of HemAT-*Bs* following CO photodissociation by time-resolved ultraviolet resonance Raman spectroscopy (UVRR). The UVRR spectra indicated two phases of intensity changes for Trp, Tyr, and Phe bands of both full-length and sensor domain proteins. The

W16 and W3 Raman bands of Trp, the F8a band of Phe, and the Y8a band of Tyr increased in intensity at hundreds of nanoseconds after CO photodissociation, and this was followed by recovery in ~50 µs. These changes were assigned to Trp-132 (G-helix), Tyr-70 (B-helix), and Phe-69 (B-helix) and/or Phe-137 (G-helix), suggesting that the change in the heme structure drives the displacement of B- and G-helices. The UVRR difference spectra of the sensor domain displayed a positive peak for amide I in hundreds of nanoseconds after photolysis, which was followed by recovery in ~50 µs. This difference band was absent in the spectra of the full-length protein, suggesting that the isolated sensor domain undergoes conformational changes of the protein backbone upon CO photolysis and that the changes are restrained by the signaling domain. The time-resolved difference spectrum at 200 µs exhibited a pattern similar to that of the static (reduced-CO) difference spectrum, although the peak intensities were much weaker. Thus, the rearrangements of the protein moiety toward the equilibrium ligand-free structure occur in a time range of hundreds of microseconds.

# 2. Spectroscopic Analyses of the Heme Environmental Structure of Aldoxime Dehydratase<sup>4)</sup>

Aldoxime dehydratase (Oxd) is a new heme-containing enzyme that works as a hydro-lyase catalyzing dehydration of aldoximes to nitriles. The enzymatic activity of Oxd is dependent on the oxidation state of the heme iron, though the reaction catalyzed by Oxd is not a redox reaction. Ferrous Oxd containing a Fe<sup>2+</sup>-heme shows the enzymatic activity, but ferric Oxd containing a Fe<sup>3+</sup>-heme does not. Previous spectroscopic analyses reveal a novel mechanism, where the change in the



Figure 1. Heme environmental structure of the substrate-bound OxdRE.

coordination mode of the substrate plays a crucial role for the regulation of the enzymatic activity. Unlike other heme enzymes, the organic substrate is directly bound to the heme iron in OxdRE. While the oxygen atom of aldoxime is coordinated to the ferric heme, the nitrogen atom of aldoxime is coordinated to the ferrous heme. The dehydration reaction proceeds only via N-coordinated substrate in the ferrous heme.

In this study, we have elucidated the heme environmental structure of Oxd from Rhodococcus sp. N-771 (OxdRE) by analyzing Fourier transform infrared (FTIR) spectra and timeresolved step-scan FTIR spectra of CO-bound Oxd. Two C-O modes of heme at 1945 and 1964 cm<sup>-1</sup> have been identified and remained unchanged in H2O/D2O exchange and in the pH 5.6–8.5 range, suggesting the presence of two conformations at the active site. The "light" minus "dark" difference FTIR spectra indicate that the heme propionates are in both the protonated and deprotonated forms, and the photolyzed CO becomes trapped within a ligand docking site (v(CO) = 2138cm<sup>-1</sup>). The time-resolved step-scan FTIR spectra show that the rate of recombination of CO to the heme is  $k_{1945 \text{ cm}^{-1}} = 126 \pm$ 20 s<sup>-1</sup> and  $k_{1964 \text{ cm}^{-1}} = 122 \pm 20 \text{ s}^{-1}$  at pH 5.6, and  $k_{1945 \text{ cm}^{-1}} =$  $148 \pm 30 \text{ s}^{-1}$  and  $k_{1964 \text{ cm}^{-1}} = 158 \pm 32 \text{ s}^{-1}$  at pH 8.5. The rate of decay of the heme propionate vibrations is on a time scale coincident with the rate of rebinding, suggesting that there is a coupling between ligation dynamics in the distal heme environment and the environment sensed by the heme propionates.

The vibrational properties of the CO adduct of OxdRE indicate the formation of a photolabile species in which the proximal histidine and the H-bonding interactions of the negatively charged heme propionates are dominant factors in controlling the strength of the Fe-CO bond. The latter observation indicates that other factors beyond the well-known proximal and distal back-bonding contributions are effective in Oxd. Taken together, these results indicate that the distal residues control the proper orientation of the bound aldoxime and, thus, modulate the heme conformation from inactive to active. Obviously, there is communication linkage between the distal and proximal sites through bond networks suggesting that there is a coupling between ligation dynamics and the environment sensed by the heme propionates. The complexity of the structural implications involved in the transition from oxidized to reduced state in the presence of aldoxime should serve as a basis for uncovering the dynamic processes involved in the reaction mechanism of the enzyme.

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### Elucidation of the Molecular Mechanisms of Protein Folding

Department of Life and Coordination-Complex Molecular Science Division of Biomolecular Functions



KUWAJIMA, Kunihiro MAKABE, Koki NAKAMURA, Takashi CHEN, Jin TAKENAKA, Toshio MIZUKI, Hiroko TANAKA, Kei Professor Assistant Professor IMS Research Assistant Professor OIIB Research Assistant Professor Post-Doctoral Fellow Technical Fellow Secretary

Kuwajima group is studying mechanisms of *in vitro* protein folding and mechanisms of molecular chaperone function. Our goals are to elucidate the physical principles by which a protein organizes its specific native structure from the amino acid sequence. In this year, we studied the fibrillogenic propensity of the GrEL apical domain and the sequential fourstate folding/unfolding of goat  $\alpha$ -lactalbumin and its Nterminal variants.

### 1. Fibrillogenic Propensity of the GroEL Apical Domain: A Janus-Faced Minichaperone

The chaperonin GroEL plays an essential role in promoting protein folding and in protecting against misfolding and aggregation in the cellular environment. In this study, we report that both GroEL and its isolated apical domain form amyloid-like fibrils under physiological conditions, and that the fibrillation of the apical domain is accelerated under acidic conditions. We also found, however, that despite its fibrillation propensity, the apical domain exhibits a pronounced inhibitory effect on the fibril growth of  $\beta_2$ -microglobulin. The analysis of the primary amino-acid sequence by programs, PASTA, TANGO and Zyggregator, which predict aggregation-prone sequences, indicates that the most aggregation-prone sequence is located in residues 260-280, which is coincident with the substrate protein-binding site in the chaperonin GroEL. Therefore, there is a close relationship between the fibrillogenic propensity and the substrate binding of GroEL. Furthermore, the analysis of 1003 sequences of the chaperonin family proteins by the Zyggregator program has shown that the aggregation-prone sequence is present in the substrate-binding site, indicating

that the close relationship between the fibrillogenic propensity and the substrate binding is a general property of the chaperonin family.



**Figure 1.** The aggregation profile of GroEL obtained by the analysis with PASTA (A), and the transmission electron microscopic observations of the isolated apical domain at pH 7.0 and 37 °C at 0.1 M NaCl in the presence of 0.5 mM SDS, where the amyloid-like fibrils of the apical domain were formed when the solution was agitated.

## 2. Sequential Four-State Folding/Unfolding of Goat α-Lactalbumin and Its N-Terminal Variants

Equilibria and kinetics of folding/unfolding of goat  $\alpha$ lactalbumin (GLA) and its two N-terminal variants were studied by circular dichroism spectroscopy. The two variants were wild-type recombinant and Glu1-deletion (E1M) variants expressed in Escherichia coli. The presence of an extra methionine at the N terminus in recombinant GLA destabilized the protein by 2 kcal/mol, while the stability was recovered in the E1M variant in which Glu1 was replaced by Met1. Kinetic folding/unfolding reactions of the proteins, induced by stoppedflow concentration jumps of guanidine hydrochloride, indicated the presence of a burst-phase in refolding, and gave chevron plots with significant curvatures in both the folding and unfolding limbs. The folding-limb curvature was interpreted in terms of accumulation of the burst-phase intermediate (I). However, there was no burst phase observed in the unfolding kinetics to interpret the unfolding-limb curvature. We thus assumed a sequential four-state mechanism, in which the folding from the burst-phase intermediate takes place via two transition states separated by a high-energy intermediate (J). We estimated changes in the free energies of the burstphase intermediate I and two transition states (‡1 and ‡2), caused by the N-terminal variations and also by the presence of stabilizing calcium ions. The  $\Phi$  values at the N terminus and at the Ca<sup>2+</sup>-binding site thus obtained increased successively during folding, demonstrating the validity of the sequential mechanism. The stability and the folding behavior of the E1M variant were essentially identical to those of the authentic protein, allowing us to use this variant as a pseudowild-type GLA in future studies.

Residue No. Protein	0	1	2	3	4	5	6	7	8	9	10
Authentic GLA	T	Е	Q	L	Т	K	С	Е	V	F	Q
Recombinant GLA	М	Е	Q	L	Т	K	С	Е	V	F	Q
E1M variant		М	Q	L	т	K	С	Е	v	F	Q

Figure 2. The amino acid sequences of the N-terminal residues of authentic GLA, recombinant GLA, and the E1M variant.



**Figure 3.** Chevron plots of authentic GLA (a), recombinant GLA (b), and the E1M variant (c) in the holo (filled symbols) and apo (open symbols) forms at pH 7.0 and 25 °C. The solid lines are theoretical curves fitted by the sequential four-state model. The broken lines and dotted lines represent the GdnHCl dependence of logarithms of microscopic rate constants ( $\ln(k_1)$ ,  $\ln(k_{-1}k_{-2}/k_2)$  and  $\ln(k_{-2})$ ) of folding and unfolding kinetics; the first two are rate-limited by the transition state 1 (‡1), and the last one by the transitions state 2 (‡2).



**Figure 4.** Schematic free-energy profiles of unfolding of GLA under a weakly unfolding condition (solid line) and under a strongly unfolding condition (broken line). Under the weakly unfolding condition, the transition state is located at ‡1, while it is located at ‡2 under the strongly unfolding condition at a high GdnHCl concentration (>5.5*M*). The free-energy profiles shown correspond to those for authentic holo GLA at 3.3*M* GdnHCl (solid line) and at 6.5*M* GdnHCl (broken line). A hypothetical intermediate (J) located between ‡1 and ‡2 is metastable under all conditions, that is, it is higher in free energy than U, I, and N.

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### Elucidation of Dynamical Structures of Biomolecules toward Understanding the Mechanisms Underlying Their Functions

Department of Life and Coordination-Complex Molecular Science Division of Biomolecular Functions



KATO, Koichi YAMAGUCHI, Takumi KAMIYA, Yukiko YAGI-UTSUMI, Maho YANAGI, Kotaro UEKUSA, Yoshinori BOONSRI, Pornthip BUI DINH, Long SUGIHARA, Takahiro CHANDAK, Mahesh ZHANG, Ying KUNIHARA, Tomoko UNO, Tsuyoshi YAMAMOTO, Sayoko KUMOI, Kentaro OKAWA, Keisuke INAGAKI, Kouya THAMMAPORN, Ratsupa SUZUKI, Mariko ISONO, Yukiko TANAKA, Kei Professor Assistant Professor IMS Research Assistant Professor\* OIIB Research Assistant Professor IMS Fellow Post-Doctoral Fellow Visiting Scientist /isiting Scientist Research Fellow Graduate Student Graduate Student Graduate Student Graduate Student Graduate Student<sup>†</sup> Graduate Student Graduate Student Graduate Student Graduate Student<sup>‡</sup> Technical Fellow **Technical Fellow** Secretary

Our biomolecular studies are based on detailed analyses of structures and dynamics of various biological macromolecules and their complexes at atomic level, using NMR spectroscopy and X-ray crystallography in conjunction with other biophysical, biochemical and molecular biology techniques. Here we report our recent studies of conformations, dynamics, and interactions of oligosaccharides and glycoconjugates along with proteins involved in the ubiquitin (Ub)-proteasome system.

### 1. Lanthanide-Assisted NMR Analyses of the Conformational Ensemble of Oligosaccharides in Conjunction with Molecular Dynamics Simulations

Conformational flexibility is an important property of biological molecules functioning in living systems, as best exemplified by oligosaccharides. We attempted to combine the lanthanide-assisted NMR method with molecular dynamics (MD) simulations for the evaluation of dynamic conformational ensembles of highly flexible oligosaccharides. In this approach, a metal-chelating tag was covalently attached to the reducing end of the oligosaccharide moieties of gangliosides, which form integral parts of cellular membranes, for observing pseudocontact shifts (PCSs). Upon complexation with paramagnetic lanthanide ions, the tagged GM3 trisaccharide, which is the common core structure shared among the gangliosides, exhibited NMR spectral changes due to PCSs according to the relative positions of the individual atoms with respect to the lanthanide ion coordinated at the tag. The observed PCS values were in excellent agreement with those back-calculated from the vast conformational ensemble of the trisaccharide derived from MD simulations (Figure 1). Thus, the PCS measurements offer a valuable experimental tool for the validation of MD simulation of highly flexible biomolecules.<sup>1)</sup> Furthermore, this approach was successfully applied to the characterization of the conformational dynamics of the branched tetrasaccharide of ganglioside GM2.<sup>2)</sup> The interbranch interactions responsible for the conformational differences between the GM2 tetrasaccahride and the GM3 trisaccharide were identified by the paramagnetic NMR method in conjunction with MD simulations. These results demonstrated the utility of our approach in the evaluation of dynamic conformational ensembles of oligosaccharides, considering their minor conformers in a systematic manner.



**Figure 1.** <sup>1</sup>H–<sup>13</sup>C HSQC spectral change of the GM3 trisaccharide tagged with a paramagnetic ion and snapshots of the sugar from an MD-simulated trajectory.

### 2. Structural Basis for Improved Effector Functions of Antibodies by Engineering of Their Glycans

More than half of proteins in nature are estimated to be modified by sugar chains, which affect the physical and biological properties of proteins. The effector functions of immunoglobulin G (IgG) critically depend on N-glycosylation of its Fc region. Removal of the fucose residue from the N-glycans of IgG-Fc results in a dramatic enhancement of antibody-dependent cellular cytotoxicity (ADCC) through improved affinity for Fcy receptor IIIa (FcyRIIIa). We successfully determined the crystal structure of the complex formed between non-fucosylated IgG1-Fc and a soluble form of FcyRIIIa (sFcyRIIIa) with two N-glycosylation sites (Figure 2a).<sup>3)</sup> The crystal structure demonstrates that one of the two N-glycans of sFcyRIIIa mediates the interaction with the N-glycan of non-fucosylated Fc, thereby stabilizing the complex. However, fucosylation of the Fc N-glycans impairs this interaction because of steric hindrance. On the other hand, our NMR data demonstrated that Tyr296 of the non-fucosylated Fc glycoform exhibits conformational multiplicity in its uncomplexed state, suggesting that conformational selection is governed by the presence or absence of the fucose residue of the Fc N-glycan. These findings offer a structural basis for improvement in ADCC of therapeutic antibodies by defucosylation.

### 3. Conformational Dynamics of Proteins Involved in the Ubiquitin-Proteasome System

While recent progresses have been made in understanding intra-domain conformational fluctuations of proteins, the evaluation of the relative motions of individual domains of multi-domain proteins is still a challenge. We successfully characterized conformational dynamics of Lys-48-linked Ub dimer (diUb) in solution using NMR spectroscopy.<sup>4)</sup> Comparison of a chemical shift of wild-type diUb with that of monomeric Ub and cyclic diUb, which mimic the open and closed states (Figure 2b), respectively, with regard to the exposure of hydrophobic surfaces to the solvent indicates that wild-type Lys-48-linked diUb in solution predominantly exhibits the open conformation (75% at pH 7.0), which becomes more populated upon lowering pH. The intrinsic properties of Lys48-linked Ub chains to adopt the open conformation may be advantageous for interacting with Ub-binding proteins. We also characterized interaction modes of the Ub-like domains of HOIL-1L and HR23 with their specific binding-partners by NMR spectroscopy.<sup>5,6)</sup>

Furthermore, we developed a novel technique for real-time monitoring of subunit exchange in homooligomeric proteins, using deuteration-assisted small-angle neutron scattering, and applied it to the tetradecamer of the proteasome  $\alpha$ 7 subunit.<sup>7)</sup>



**Figure 2.** 3D structures of (a) sFcγRIIIa bound to non-fucosylated Fc and (b) wild-type Lys-48-linked diUb. The open form of diUb (purple) was superposed on the closed state (green).

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#### Awards

YAMAMOTO, Sayoko; Young Scientists Poster Awards, The International Symposium on Nuclear Magnetic Resonance 2011 (2011). KATO, Koichi; The Erwin von Bälz Prize 2011 (First Prize) (2011). YAMAGUCHI, Takumi; CSJ Presentation Award 2012, The 92<sup>nd</sup> Annual Meeting of The Chemical Society of Japan (2012). ZHANG, Ying; FY2012 Sokendai President's Award (2012). KUMOI, Kentaro; Young Poster Award, The 12<sup>th</sup> Annual Meeting of The Protein Science Society of Japan (2012).

<sup>\*</sup> Present Address; Graduate School of Engineering, Nagoya University

<sup>†</sup> carrying out graduate research on Cooperative Education Program of IMS with Nagoya City University

<sup>‡</sup> carrying out graduate research on Cooperative Education Program of IMS with Kasetsart University

### Structure-Function Relationship of Metalloproteins

### Department of Life and Coordination-Complex Molecular Science Division of Biomolecular Functions



FUJII, Hiroshi KURAHASHI, Takuya CONG, Zhiqi OHTSUKI, Akimichi WANG, Chunlan TANIZAWA, Misako Associate Professor Assistant Professor IMS Fellow Post-Doctoral Fellow Graduate Student Secretary

Metalloproteins are a class of biologically important macromolecules, which have various functions such as oxygen transport, electron transfer, oxidation, and oxygenation. These diverse functions of metalloproteins have been thought to depend on the ligands from amino acid, coordination structures, and protein structures in immediate vicinity of metal ions. In this project, we are studying the relationship between the electronic structures of the metal active sites and reactivity of metalloproteins.

### 1. Formation of Iron(III) *Meso*-Chloro-Isoporphyrin as a Reactive Chlorinating Agent from Oxoiron(IV) Porphyrin $\pi$ -Cation Radical<sup>1)</sup>

Oxoiron(IV) porphyrin  $\pi$ -cation radicals are generally known to function as key reactive intermediates in a variety of oxidation reactions catalyzed by heme enzymes such as cytochrome P450. The oxoiron(IV) porphyrin  $\pi$ -cation radical complex has several isoelectronic forms which are two oxidation state equivalents higher than that of the iron(III) porphyrin complex. Isoporphyrins, tautomers of porphyrins with a saturated meso carbon, were originally postulated by Woodward, and its metal complex was first reported by Dolphin et al., who prepared a zinc(II) 5'-methoxy-5,10,15,20-tetraphenylisoporphyrin complex by nucleophilic attack of methanol on zinc(II) 5,10,15,20-tetraphenylporhyrin  $\pi$ -dication complex. Since then, iron(III) isoporphyrin complexes, particularly meso-tetraaryl derivatives, have been synthesized chemically and electrochemically. However, to our surprise, atom transfer reactions of isoporphyrin complexes to substrates, such as oxygenation and halogenation reactions, have not been studied well. In this study, we show that an oxoiron(IV) porphyrin  $\pi$ -cation radical complex can be converted to iron(III) mesochloro-isoporphyrin complex in the presence of trifluoroacetic acid (TFA) and chloride ion. The formation of the isoporphyrin complex would be due to protonations of the oxo ligand of oxoiron(IV) porphyrin  $\pi$ -cation radical species. More importantly, this study shows that the iron(III) *meso*-chloro-iso-porphyrin complex is a reactive reagent for chlorination of aromatic compounds and olefins.



**Figure 1.** Reaction of oxoiron(IV) porphyrin  $\pi$ -cation radical complex with chloride ion in the presence and absence of proton.

### 2. The Effect of the Axial Ligand on the Reactivity of the Oxoiron(IV) Porphyrin $\pi$ -Cation Radical Complex: Higher Stabilization of the Product State Relative to the Reactant State<sup>2</sup>)

The proximal heme axial ligand plays an important role in tuning the reactivity of oxoiron(IV) porphyrin  $\pi$ -cation radical species (compound I) in enzymatic and catalytic oxygenation reactions. To reveal an essence of the axial ligand effect on the reactivity, we investigated from a thermodynamic viewpoint. Compound I model complexes, (TMP<sup>++</sup>)Fe<sup>IV</sup>O(L) (where TMP is 5,10,15,20-tetramesitylporphyrin and TMP<sup>++</sup> is its  $\pi$ -cation radical), can be provided with altered reactivity by changing the identity of the axial ligand, but the reactivity is not correlated with spectroscopic data (v(Fe=O), redox potential, and so on) of (TMP<sup>++</sup>)Fe<sup>IV</sup>O(L). Surprisingly, a clear correlation was found between the reactivity of (TMP<sup>++</sup>)



**Figure 2.** Curve crossing diagram of potential-energy surfaces of the reactant,  $(TMP^{+*})Fe^{IV}O(L)$ , and product,  $(TMP)Fe^{III}(L)$ , states. Blue line; Change of stability of  $(TMP^{+*})Fe^{IV}O(L)$  and red line; change of stability of  $(TMP)Fe^{III}(L)$ .

Fe<sup>IV</sup>O(L) and the Fe<sup>II</sup>/Fe<sup>III</sup> redox potential of (TMP)Fe<sup>III</sup>L, the final reaction product. This suggests that the thermodynamic stability of (TMP)Fe<sup>III</sup>L is involved in the mechanism of the axial ligand effect. Axial ligand-exchange experiments and theoretical calculations demonstrate a linear free-energy relationship, in which the axial ligand modulates the reaction free energy by changing the thermodynamic stability of (TMP) Fe<sup>III</sup>(L) to a greater extent than (TMP<sup>+•</sup>)Fe<sup>IV</sup>O(L). The linear free energy relationship could be found for a wide range of anionic axial ligand and for various types of reactions, such as epoxidation, demethylation, and hydrogen abstraction reactions. An essence of the axial ligand effect is neither the electron donor ability of the axial ligand nor the electron affinity of compound I, but the binding ability of the axial ligand (the stabilization by the axial ligand). An axial ligand that binds more strongly makes (TMP)Fe<sup>III</sup>(L) more stable and (TMP+•)Fe<sup>IV</sup>O(L) more reactive. All results indicate that the axial ligand controls the reactivity of compound I (the stability of the transition state) by the stability of the ground state of the final reaction product and not by compound I itself.

### 3. Oxidation of Chloride Ion and Subsequent Chlorination of Organic Compounds by Oxoiron(IV) Porphyrin π-Cation Radical Complexes<sup>3)</sup>

Enantioselective transition-metal-catalyzed oxygenation

reactions have received much attention because of the demand for organic synthesis strategies and their biological relevance with respect to metalloenzymes. Terminal oxidants such as peroxides, iodosylarenes, and peracids have been utilized as an oxygen source for these oxygenation reactions. Since the terminal oxidants must be stable for easy handling, the primary role of the transition-metal catalyst is to activate a stabilized oxidant and to generate a transient species that remains active enough to transfer an oxygen atom to a substrate. The activation of a terminal oxidant is initiated by binding to the metal complex to form a terminal oxidant adduct of the metal complex. Recently, evidence has been mounting in support of the proposal that the terminal oxidant adduct of a metal complex is not only a precursor to a reactive high valent metal-oxo species, but also itself may serve as a reactive species for an oxygenation reaction. Although terminal oxidant adducts of metal complexes are unstable and reactive compounds in most cases, metal complex adducts with hydrogen peroxide, alkylperoxides, and m-CPBA have been isolated and structurally characterized. In contrast to these successful reports, and much to our surprise, there have been no examples of structural characterization of any iodosylarene adducts of metal complexes, although they have emerged as useful oxidants for various organic reactions. The most intensive spectroscopic study was performed by Hill et al., who thoroughly investigated an iodosylbenzene adduct of a manganese porphyrin complex with <sup>1</sup>H NMR, IR and <sup>127</sup>I Mössbauer spectroscopy. However, the nature of the bonding interaction between iodosylbenzene and the metal ion remains unclear. Here, we report on the preparation and X-ray crystal structure of an iodosylarene adduct of a manganese(IV) salen complex bearing a trans-cyclohexane-1,2-diamine linkage as chiral unit.



**Figure 3.** Synthesis and structural characterization of bis-iodosylmesitylene adduct of manganese(IV) salen complex.

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#### Award

WANG, Chunlan; The best pster presentation, The 44<sup>th</sup> Symposium on Chemical and Biological Oxidation (2011).

### Molecular Origin of 24 Hour Period in Cyanobacterial Protein Clock

Department of Life and Coordination-Complex Molecular Science Division of Biomolecular Sensing



AKIYAMA, Shuji ABE, Hitomi Professor (April, 2012–) Secretary

Circadian (approximately 24 h) clocks are endogenous time-keeping systems encapsulated in living cells, enabling organisms to adapt to daily fluctuation of exogenous environments on the Earth. These time-keeping systems, found ubiquitously from prokaryotes to eukaryotes, share the three characteristics. First, the circadian rhythmicity of the clocks persists even without any external cues (self-sustainability). Second, the period is little dependent on ambient temperature (temperature compensation). Third, the phase of the clock can be reset by external stimuli such as lightning, humidity, or temperature so as to be synchronized to the external phase (synchronization).

KaiC, a core protein of the circadian clock in cyanobacteria, undergoes rhythmic structural changes over approximately 24 h in the presence of KaiA and KaiB (Kai oscillator). This slow dynamics spanning a wide range of both temporal and spatial scales is not well understood, and is central to a fundamental question: what determines the temperature-compensated 24 h period?<sup>1,2</sup>) The Kai oscillator reconstitutable *in vitro* is advantageous for studying its dynamic structure through a complementary usage of both X-ray crystallography and solution scattering, its transient response by using physicochemical techniques, and its molecular motion through a collaborative work with computational groups.

Our mission is to explore the frontier in molecular science of the cyanobacterial circadian clock from many perspectives. This Annual Report summarizes our recent activities from April 1, 2012 through August 31, 2012.

### 1. Tracking the Ticking of Cyanobacterial Clock Protein KaiC in Solution<sup>3)</sup>

The ATPase activity of KaiC *alone* is strongly correlated with the oscillatory period of the Kai oscillator. This correlation suggests that the ATPase activity of KaiC is one of the period-determining factors of the Kai oscillator. Hence, the determination of the structural change of KaiC interlocked with the ATPase activity is of great of importance.



**Figure 1.** Circadian dynamics of cyanobacterial clock protein KaiC. The C1 and C2 domains in each protomer of KaiC are drawn as red and blue spheres, respectively. Expansion and contraction motions of the C2 ring (B, C) in solution serves as a timing cue for assembly/ disassembly of KaiA and KaiB (D), and is interlocked with its C1 ATPase udder a control of negative-feedback regulation (A).

To track the dynamic transition of KaiC in real-time, we recorded the time evolution of intrinsic tryptophan (Trp) fluorescence from KaiC contained in the Kai oscillator. KaiC is a dumbbell-shaped molecule composed of tandemly duplicated N-terminal (C1) and C-terminal (C2) domains. Six protomers are assembled into a hexamer to attain a doubledoughnut shape. Two tryptophan (Trp) residues located in the protomer-protomer interface of the C2 domain can serve as a sensitive probe to monitor the potential structural transition of the C2 ring. The intensity of the Trp fluorescence from KaiC revealed a rhythmic fluctuation with the period of approximately 24 h (Figure 1, panel C). So far as we know, this is the first experimental evidence that demonstrated a dynamic



**Figure 2.** Octuple cuvette. (a) Photograph of the central plate covered by the quartz window, Teflon spacer and external plate, in that order. (b) Front view of the assembly. (c) Side view of the assembly. (d) Filling and removal of solutions using a commercially available eight-channel pipette. The sample volume of each chamber is  $25 \,\mu$ l.

structural transition of the C2 ring of KaiC in solution.

Concomitantly with the Trp-fluorescence dynamics, KaiC underwent a periodic change in its phosphorylation state (Figure 1, panel B). KaiC has the two phosphorylation sites, *i.e.*, Ser431 and Thr432, in the C2 domain, and both residues are phosphorylated and then dephosphorylated in a programmed sequence during the phosphorylation cycle as follows: KaiCS/pT  $\rightarrow$  KaiCpS/pT  $\rightarrow$  KaiCpS/T  $\rightarrow$  KaiCpS/T (where 'S' represents Ser431, 'pS' represents phosphorylated Ser431, 'T' represents Thr432, and 'pT' represents phosphorylated Thr432). Interestingly, the Trp fluorescence was maximized at the timing when the KaiCpS/T state was populated (Figure 1, panels B and C). The result suggests the Trp fluorescence is an excellent measure of the phosphor-coupled transition of the C2 ring in KaiC.

### 2. Visualization of Dynamic Structural Changes of KaiC Using Small-Angle X-Ray Solution Scattering Technique<sup>3,4)</sup>

To visualize the C2-ring dynamics confirmed by tracking Trp fluorescence, we measured the small-angle x-ray scattering (SAXS) from KaiC in solution. To obtain the SAXS data of biological samples in solution, one must first record the scattering intensity of the sample (biomacromolecules in solution) and then that of the matching buffer in the separate



**Figure 3.** Expansion and contraction motions of C2 ring of KaiC interlocked with ATPase activity.

experiment, and finally find the difference between two intensities. The cuvette used for conventional SAXS experiments has only a single observation chamber in order to ensure the qualitative subtraction of the scattering contributed by the solvent molecules. On the other hand, the use of the singlechamber cuvette makes both the experiment and analysis time-consuming.

To record the SAXS pattern of KaiC both efficiently and qualitatively, we designed and constructed an eight-chamber cuvette (octuplet cuvette), each chamber of which was fabricated so uniformly to ensure the inter-chamber subtraction (Figure 2). The developed cuvette enabled us to acquire SAXS dataset of KaiC roughly 10 times faster without any significant degradation of data quality.

On the basis of the obtained SAXS data, we built lowresolution models of the KaiC hexamer as shown in Figure 3. The overall shape is almost unchanged in the transition from KaiCS/pT to KaiCpS/pT, whereas the radius of the C2 ring is dramatically enlarged in the subsequent transition from KaiCpS/ pT to KaiCpS/T. The expanded C2 ring is partly contracted in the transition from KaiCpS/T to KaiCS/T, and is further contracted in the subsequent transition from KaiCS/T to KaiCS/pT. The present model suggests that KaiC ticks through expanding and contracting motions of the C2 ring.

The dynamic motion of the C2 ring uncovered throughout our study is chronobiologically meaningful, we believe, in terms of the elucidation of the key conformational change tightly coupled to the period-determining ATPase of KaiC. Our group is trying to improve spatio-temporal resolution of the experiments so as to draw a more dynamic and detailed picture of KaiC ATPase.

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# Investigation of Molecular Mechanisms of Channels, Transporters and Receptors in Membrane

### Department of Life and Coordination-Complex Molecular Science Division of Biomolecular Sensing



FURUTANI, Yuji KIMURA, Tetsunari TSUKAMOTO, Hisao INAGUMA, Asumi GUO, Hao FUJIWARA, Kuniyo SHIMIZU, Atsuko

Associate Professor Assistant Professor IMS Research Assistant Professor Post-Doctoral Fellow Graduate Student Graduate Student Secretary

Membrane proteins are important for homeostasis of living cells, which work as ion channels, transporters, various types of chemical and biophysical sensors, and so on. These proteins are considered as important targets for biophysical studies.

Our main goal is to clarify molecular mechanisms of channels, transporters and receptors in cell membrane mainly by using stimulus-induced difference infrared (IR) spectroscopy, which is sensitive to the structural and environmental changes of organic- and bio-molecules. Recently, Dr. Kimura started to construct a microfluidic device to monitor biological and chemical reactions by infrared and fluorescent microscopic techniques. Dr. Tsukamoto has established a protein expression system with mammalian cell line.

### 1. Anion-Transport Mechanism of a Chloride Ion Pump *p*HR Studied by Time-Resoved FTIR Spectroscopy

pharaonis Halorhodopsin (pHR) is a light-driven inward chloride ion pump protein. The ion transportation is performed through the sequential formation of several intermediates (K, L<sub>1</sub>, L<sub>2</sub>, N and O) during the photocyclic reaction. The structural details of each intermediate state have been studied by various kinds of physicochemical methods, however it is still open question how the structural changes of protein molecule and water molecules are involved in the translocation of a chloride ion inside protein at physiological temperature. To analyze the structural dynamics, we performed the timeresolved Fourier transform infrared (TR-FTIR) spectroscopy in the whole mid-infrared region under various hydration conditions. Measurements under D<sub>2</sub>O reveal the structural information of the water inaccessible backbone of protein itself, and those with H<sub>2</sub>O or H<sub>2</sub><sup>18</sup>O give an insight on the dynamics of water molecules inside protein. We concluded that the chloride ion release and uptake occurring in the N and O intermediate states are accompanied by the drastic conformational changes in the water-inaccessible transmembranehelices and the hydrogen- bonding network rearrangement of the internal water molecules in pHR (manuscripts in preparation).

### 2. Time-Resolved Difference FTIR Spectroscopy Triggered by Rapid Buffer-Exchange

Attenuated total reflection (ATR) FTIR spectroscopy is a powerful technique to obtain infrared spectra of membrane proteins immersed in aqueous solution.<sup>1)</sup> By exchanging buffer with and without salts, the difference spectra between the two conditions provide the structural information relating to the interaction between protein and ions. In this year, we executed the construction of the kinetic ATR FTIR spectroscopic system, which can follow the chemical and biological reaction triggered by the rapid buffer-exchange, in collaboration with Mr. K. Okamoto in UNISOKU Co. Ltd.

We have developed the phototriggering system for time-resolved FTIR measurements, but the application of this system has been limited mainly to the photo-inductive chemical reactions. Rapid buffer exchange (<1 ms) allows us to monitor the early stage of reactions, such as substrate or ion binding. For further improvement of the time-resolution, we will construct the flowflash system.



**Figure 1.** Picture of the rapid buffer-exchange system.

### 3. Development of a Microfluidic Device to Monitor Biological and Chemical Reactions

Real-time observation is one of the powerful techniques to understand the molecular mechanisms of the self-organization and molecular association. The solution mixing technique realizes many reaction fields for biological and chemical reactions by changing the buffer condition and can be combined with spectroscopic equipment easily. However, conventional mixing techniques limit their targets because of the large consumption of the sample and the stress of the turbulent mixing. We are developing a novel microfluidic device and trying to reduce both the sample consumption to ~1/1000 and the stress by using sheath flow.



**Figure 2.** Pictures of the developed microfluidic mixer. (A) Picture of the mixer on the XY-stage of the fluorescence microscope. (inset) Schematic figure of the designed flow channels. (B) Fluorescence image of the Ca-Green mixed with the Ca<sup>2+</sup> buffer.

Microfluidic devices with the flow channels of micrometer scale are made of a single cast of poly(dimethylsiloxane), PDMS, with the photolithography technique. These devices were combined with the fluorescence microscope and the complete mixing with the time resolution of ~20  $\mu$ s was confirmed by the fluorescent intensity measurements. We are attempting to combine these devices with FTIR or other fluorescence spectroscopy.

The devices were manufactured with the strong supports of Equipment Development Center of IMS (especially Ms. N. Takada and Mr. M. Aoyama).

### 4. Investigation of Membrane Proteins Which Are Hardly Expressed in *E. coli*

Heterologous expression system using *E. coli* is established and easy way to obtain large amount of purified proteins, but in general, expression of membrane proteins, in particular mammalian ones, is very difficult. In this year, we have tried to establish a protein expression system using cultured mammalian cells for IR spectroscopic analyses of various membrane proteins, including mammalian ones.

In order to study structure-function relationship using purified proteins, it is very important to select target proteins that can be highly expressed in cell lines. Traditionally the screening of membrane proteins is very painful and timeconsuming. However, a powerful method named fluorescencedetected size-exclusion chromatography (FSEC) was recently developed (Kawate and Gouaux, 2006). FSEC method evaluates expression level and properties of target proteins using GFP-tag and size-exclusion chromatograph even without purification (Figure 3). We set up an FSEC system and



Figure 3. Scheme of the fluorescence-detected size-exclusion chromatography (FSEC).

expressed various membrane proteins with GFP tag to select appropriate proteins. Also, we tried to purify screened proteins and apply them to IR spectroscopic analyses.

### Channelrhodopsins (collaboration with Prof. Yawo's group in Tohoku Univ.)

Channelrhodopin is a light-gated cation channels, which is originally identified in Chlamydomonas. Currently this protein is used as a tool for "optogenetics" which enables excitation of specific neural cells by irradiation. Prof. Yawo's group has developed several interesting channelrhodopsin mutants that show unique electrophysiological properties. Prof. Yawo kindly provided constructs of the channelrhodopsin mutants to us, and we screened the mutants using FSEC method. Several mutants can be expressed efficiently in mammalian cells, and can be purified. Also, we have already succeeded to measure the difference IR spectra before and after photo-activation of the channelrhodopsins. We are now trying to analyze the spectra and extend to time-resolved measurements with 12.5 µs resolution to reveal real-time analyses of open-close mechanism of the channel protein. (collaboration with Prof. Nureki's group in Univ. of Tokyo)

### Mammalian potassium channels (collaboration with Prof. Kubo's group in National Institute for Physiological Sciences (NIPS))

Mammalian potassium channels play important roles in various organs, including brain, heart, kidney and retina. Prof. Kubo's group has revealed structure-function relationship of these channels mainly using electrophysiological method. Prof. Kubo kindly provided constructs of various mammalian potassium channels to us. We tested properties of these channels expressed in cultured cells using FSEC, and found several inwardly rectifying potassium channels can be expressed efficiently and show appropriate properties. In addition, we have just successfully purified the channels. Currently we are trying to measure IR spectra of the purified channels.

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### Heterogeneous Catalytic Systems for Organic Chemical Transformations in Water

Department of Life and Coordination-Complex Molecular Science Division of Complex Catalysis



UOZUMI, Yasuhiro OSAKO, Takao HAMASAKA, Go ZHOU, Haifeng KOBAYASHI, Noboru SAKURAI, Fumie TSUJI, Hiroaki ITO, Kenichi TORII, Kaoru SASAKI, Tokiyo FUKUSHIMA, Tomoko Professor Assistant Professor Assistant Professor Post-Doctoral Fellow Graduate Student Graduate Student Graduate Student Graduate Student Technical Fellow Secretary Secretary

Various transition metal-catalyzed organic molecular transformations in water were achieved under heterogeneous conditions by use of poly(imidazole-palladium), amphiphilic resinsupported palladium complexes, or a microflow device containing a polymeric palladium nanoparticle membrane which were designed and prepared by this research group. The enantioselective carbenoid insertion into phenolic O-H bonds with a new chiral copper(I) imidazoindolephosphine complex has been also developed. In particular, development of a highly active reusable poly(imidazole-palladium) and a microflow device containing a plolymeric Pd nanoparticle membrane for organic transformations in water and the highly enantioselective O-H insertion using a new chiral copper(I) complex are highlights among the achievements of the 2011-2012 period to approach what may be considered ideal chemical processes of next generation. Representative results are summarized hereunder.

### 1. Self-Assembled Poly(Imidazole-Palladium): A Highly Active, Reusable Catalyst<sup>1,2)</sup>

A polymeric imidazole Pd catalyst (MEPI-Pd) was readily prepared by the coordinative convolution of (NH<sub>4</sub>)<sub>2</sub>PdCl<sub>4</sub> and poly[(*N*-vinylimidazole)-*co*-(*N*-isopropylacrylamide)<sub>5</sub>] in a methanol/water solution at 80 °C for 30 min. The polymeric Pd catalyst was utilized for the allylic arylation/alkenylation/ vinylation of allylic esters and carbonates with aryl/alkenylboronic acids, vinylboronic acid esters, and tetraaryl borates. Even 0.8–40 mol ppm Pd of the catalyst efficiently promoted allylic arylation/alkenylation/vinylation in alcohol and/or water with a catalytic turnover number (TON) of 20,000– 1,250,000. Furthermore, the polymeric Pd catalyst efficiently promoted the Suzuki-Miyaura reaction of a variety of inactivated aryl chlorides, as well as aryl bromides, and iodides in water with a TON of up to 3,570,000.



Scheme 1. Preparation of Self-Assembled Poly(Imidazole-Palladium) (top) and Application to Allylic Arylation and Suzuki-Miyaura Reaction.

### 2. Instantaneous Hydrodehalogenation of Haloarenes by a Microflow Device Containing a Polymeric Pd Nanoparticle Membrane<sup>3)</sup>

An architecture-based system of transition metal catalysis usinWe developed a variety of polymeric Pd nanoparticle membrane-installed microflow devices for the first time. Three types of polymers were convoluted with palladium salts under laminar flow conditions in a microflow reactor to form polymeric Pd membranes at the laminar flow interface. These membranes were reduced with aqueous sodium formate or with heat to create microflow devices containing polymeric palladium nanoparticle membranes. These microflow devices achieved instantaneous hydrodehalogenation of 10–1,000 ppm of aryl chlorides, bromides, iodides, and triflates, within a residence time of 2–8 s at 50–90 °C using a safe, nonexplosive aqueous sodium formate to quantitatively afford the corresponding hydrodehalogenated products. PCB (10–1,000 ppm) and PBB (1,000 ppm) were completely decomposed under similar conditions, yielding biphenyl as a fungicidal compound.



**Scheme 2.** Instantaneous Hydrodehalogenation of Haloarenes by a Microflow Device Containing a Polymeric Pd Nanoparticles Membrane.

### 3. C–N and C–S Bond Forming Cross Coupling in Water with Amphiphilic Resin-Supported Palladium Complexes<sup>4)</sup>

Catalytic C–N and C–S bond forming reactions of haloarenes with secondary amines and thiophenols were achieved in water under heterogeneous conditions by the use of immobilized palladium complexes coordinated with the amphiphilic polystyrene-poly(ethylene glycol) resin-supported di(*tert*butyl)phosphine ligand to afford aryl(dialkyl)amines and diarylsulfides in high yield.



Scheme 3. C–N and C–S Bond Forming Cross Coupling in Water with Amphiphilic Resin-Supported Pd Complexes.

### 4. Enantioselective Carbenoid Insertion into Phenolic O–H Bonds with a Chiral Copper(I) Imidazoindolephosphine Complex<sup>5)</sup>

The enantioselective O–H carbenoid insertion reaction with a new chiral copper(I) imidazoindolephosphine complex has been developed. The chiral copper(I) complex catalyzed the insertion of carbenoids derived from  $\alpha$ -diazopropionates into the O–H bonds of various phenol derivatives to give the corresponding  $\alpha$ -aryloxypropionates with up to 91% ee.



**Scheme 4.** Enantioselective Carbenoid Insertion into Phenolic O–H Bonds with a Chiral Copper(I) Imidazoindolephosphine Complex.

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### Water Oxidation Catalyzed by Dimeric Ru Complexes

Department of Life and Coordination-Complex Molecular Science Division of Functional Coordination Chemistry



TANAKA, Koji OHTSU, Hideki KOBAYASHI, Katsuaki PADHI, Sumanta Kumar NAKANE, Daisuke YAMAGUCHI, Yumiko NOGAWA, Kyoko Professor (-March, 2012)\* Assistant Professor\* IMS Research Assistant Professor\* Post-Doctoral Fellow Post-Doctoral Fellow Secretary Secretary

Artificial photosynthesis aimed at carbon dioxide reduction and water splitting has become a top research theme. Two-electron transfer from or to substrates through redox reactions is requisite for stable molecular transformation. Intermolecular electron transfer, however, always generates free radical species, which often causes undesired side reactions. Success of artificial photosynthesis, therefore, depends on the designing of reaction systems that can provide or take out multi-electrons to or from reaction centers.

Irradiation of  $[Ru(bpy)_2(pbn)]^{2+}$  (pbn = 2-(2-pyridyl)benzo[*b*]-1,5-naphthyridine) with visible light causes proton coupled one-electron reduction, and the subsequent disproportionation affords an NADH model complex,  $[Ru(bpy)_2$  $(pbnH_2)]^{2+}$  (pbnH<sub>2</sub> = 5,10-dihydro-2-(2-pyridyl)benzo[*b*]-1,5naphthyridine) (Scheme 1). Smooth conversion from NAD analog to NADH one under visible light irradiation would lead to a new methodology for utilization of water as the hydrogen source in molecular transformation.



Scheme 1. Photochemical two-electron reduction of [Ru(bpy)<sub>2</sub>(pbn)]<sup>2+</sup>.

Three Ru-dioxolene complexes,  $[Ru^{II}(Q)(trpy)(OH_2)]^{2+}$ ,  $[Ru^{II}(Q)(trpy)(OH)]^+$ , and  $[Ru^{II}(Sq)(trpy)(O^{-*})]^0$  (Q = 3,5-dibutylquinone; Sq = 3,5-dibutylsemiquinone) exist as equilibrium mixtures in water with the p $K_a$  values of 5.5 and 10.5, respectively. A novel oxyl radical complex,  $[Ru^{II}(Sq)(trpy)(trpy)(trpy)]^{1+}$ 



Scheme 2. Two-electron oxidant generation taking advantage of acidbase equilibrium of Ru-aqua(dioxolene) complex.

 $(O^{-\bullet})]^0$ , undergoes reversible one electron oxidation around 0 V (*vs.* Ag/AgCl), and the subsequent further one electron oxidation creates the unique Ru<sup>III</sup>–O<sup>-•</sup> framework in the product (Scheme 2). The oxyl radical and Ru(III) center involved in the product work as simultaneous hydrogen atom and one electron acceptors in two-electron oxidation of alcohols.

### 1. Direct Evidence for O–O Bond Formation in the Four Electron Water Oxidation

The difficulty of water decomposition results from fourelectron oxiation of water rather than that of two-electron reduction. Among various water oxidation catalysts reported so far, much attention has been paid to a dinuclear Ru

complex,  $[Ru_2(OH)_2(Bu_2q)_2$ (btpyan)]<sup>2+</sup> (Bu<sub>2</sub>q = 3,6-di-*tert*butylquinone, btpyan = 1,8-bis (terpyridyl)-anthracene), known as Tanaka Catalyst because of its high activity toward four-electron oxidation of water. However, any direct evidences for O–O bond formation prior to O<sub>2</sub> evolution have not been obtained so far.



Four redox centers  $(2 \times Ru^{II/III} and 2 \times [Q]/[Sq])$  of  $[Ru_2(OH)_2$ 

(Bu<sub>2</sub>q)<sub>2</sub>(btpyan)]<sup>2+</sup> is attributable to the high catalytic capacity. So, we tried to detect the O–O bond formation process by decreasing the number of redox centers of an Ru dinuclear complex. Two-electron oxidation of [Ru<sup>II</sup><sub>2</sub>(Cl)<sub>2</sub>(bpy)<sub>2</sub> (btpyan)]<sup>4+</sup> at +1.0 V in a range of pH 2.0 to 3.0 forms [Ru<sup>IV</sup><sub>2</sub>(=O)<sub>2</sub>(bpy)<sub>2</sub>(btpyan)]<sup>4+</sup>. The controlled potential electrolysis of [Ru<sup>II</sup><sub>2</sub>( $\mu$ -Cl)(bpy)<sub>2</sub>(btpyan)]<sup>3+</sup> at +1.60 V in water at pH 2.6 (H<sub>3</sub>PO<sub>4</sub>/NaH<sub>2</sub>PO<sub>4</sub> buffer) catalytically evolved dioxygen. Addition of a CH<sub>3</sub>CN (100 µl) solution of [Ru<sup>II</sup><sub>2</sub>( $\mu$ -Cl)(bpy)<sub>2</sub> (btpyan)]<sup>3+</sup> (1.0 µmol) into an aqueous solution (10 ml) of Ce(NH<sub>4</sub>)<sub>2</sub>(NO<sub>3</sub>)<sub>6</sub> (2.5 mmol) at pH 1.0 (adjusted with HNO<sub>3</sub>) also caused O<sub>2</sub> evolution (414 µmol).

The electronic absorption spectra of the reaction mixture showed a transient band at 688 nm that emerged only in O<sub>2</sub> evolution. The transient complex prepared by the electrolysis of  $[Ru^{II}_2(\mu$ -Cl)(bpy)<sub>2</sub>(btpyan)]^{3+} in  $H_2^{16}O$  and  $H_2^{18}O$  at +1.40 V displayed two absorption bands at 442 and 824 cm<sup>-1</sup>, and 426 and 780 cm<sup>-1</sup>, respectively, in the resonance raman spectra with irradiation at 633 nm. The isotope shifts ( $\Delta$  16 and 44  $cm^{-1}$ ) between  $H_2^{16}O$  and  $H_2^{18}O$  are quite consistent with the calculated values of the v(Ru-O) and v(O-O) bands of [Ru<sub>2</sub>  $(\mu$ -O<sub>2</sub>)(bpy)<sub>2</sub>(btpyan)]<sup>3+</sup>. If the Ru<sup>IV</sup>=O bond of [Ru<sup>IV</sup><sub>2</sub>(=O)<sub>2</sub> (bpy)<sub>2</sub>(btpyan)]<sup>4+</sup> generated by the oxidation of [Ru<sup>III</sup><sub>2</sub>(OH)  $(bpy)_2(btpyan)]^{4+}$  has more or less Ru<sup>III</sup>–O. oxyl radical character,  $[Ru^{III}_2(\mu-O_2)(bpy)_2(btpyan)]^{4+}$  would be produced by the radical coupling of the two oxo groups of [Ru<sup>IV</sup><sub>2</sub>(=O)<sub>2</sub>  $(bpy)_2(btpyan)]^{4+}$ . The rate of O<sub>2</sub> evolution by chemical oxidation using Ce(IV) is much faster than that of electrochemical reaction at +1.60 V. Nuclear attack of two water to Ru<sup>III</sup> of [Ru<sup>III</sup><sub>2</sub>(µ-O<sub>2</sub>)(bpy)<sub>2</sub>(btpyan)]<sup>4+</sup> will slowly releases O<sub>2</sub> (Scheme 3). On the other hand, Ce(IV) causes further oxidation of  $[Ru^{III}_2(\mu-O_2)(bpy)_2(btpyan)]^{4+}$  to produce  $[Ru^{IV}_2(\mu-O_2)]^{4+}$  $(bpy)_2(btpyan)]^{6+}$ , which smoothly evolves  $O_2$  with regeneration of [Ru<sup>III</sup><sub>2</sub>(OH)(bpy)<sub>2</sub>(btpyan)]<sup>4+</sup> (Scheme 3). In accordance with this,  $[Ru^{III}_{2}(\mu-O_{2})(bpy)_{2}(btpyan)]^{4+}$  was detected only after Ce(IV) was consumed in O2 evolution.



Scheme 3. Four-electron oxidation of water catalyzed by  $[Ru^{III}_2(OH)_2(bpy)_2(btpyan)]^{4+}$ .

### 2. Insight for Activity of Tanaka Catalyst toward Water Oxidation

High catalytic ability of [Ru<sub>2</sub>(btpyan)(3,6-di-Bu<sub>2</sub>Q)<sub>2</sub> (OH<sub>2</sub>)]<sup>2+</sup> (Tanaka catalyst) toward water oxidation produces disputes about the electronic structures in the catalytic cycle. DFT computational works reported so far are not consistent with each other in the viewpoints of the relative stability between the closed-shell (Ru<sup>II</sup>-Q) and open-shell (Ru<sup>III</sup>-SQ) electronic structures, and the pathway to the O-O bond formation. On the other hand, broken-symmetry hybrid density functional computations have provided a rational reaction mechanism for water oxidation (Scheme 4). Deprotonation of waters in  ${}^{1}\mathbf{1}$  affords the key tetraradical intermediate  ${}^{1}\mathbf{3}$  via one-electron transfer (OET) in <sup>1</sup>2. The oxygen-radical pair in <sup>1</sup>**3** is local triplet diradical (LTD), suppressing facile O–O bond formation by the radical coupling (RC) mechanism. The two-electron removal from  $^{1}3$  provides the hexaradical species <sup>3</sup>4. The oxygen radical pair (<sup>3</sup>A) is still LTD-type, indicating the necessity of spin inversion (SI) for generation of local singlet diradical (LSD) pair in <sup>1</sup>5. The RC mechanism in <sup>1</sup>5 is facile, giving the peroxide species  $^{1}6$ . The next step for generation of oxygen dianion may become the rate-determining step as shown in <sup>1</sup>7. The  $\beta$ -spin at the terminal oxygen anion in <sup>1</sup>7 is moved to the Ru<sub>2</sub>(III) site with the  $\alpha$ -spin to form the singlet pair as shown in <sup>1</sup>8. The spin exchange (SE) between  $\downarrow \bullet Ru_1(III)$  and  $SQ_1 \bullet \uparrow$  to generate  $\uparrow \bullet Ru_1(III)$  and  $SQ_1 \bullet \downarrow$  is necessary for one more OET from superoxide anion to  $\uparrow \cdot Ru_1(III)$  to afford triplet molecular oxygen in <sup>1</sup>9. The SE process is easy because the exchange coupling for the •O-O-Ru(III)• is weak. Thus the SQ1• radical plays an important role for spin catalysis. The two-electron removal from <sup>1</sup>9 is necessary for reproduction of <sup>1</sup>**1**. Thus the BS computational results provide the orbital and spin correlation diagram for water splitting reaction.



**Scheme 4.** Proposed mechanism for water oxidation catalyzed by Tanaka catalyst.

### Synthetic Inorganic and Organometallic Chemistry of Transition Metals

Department of Life and Coordination-Complex Molecular Science Division of Functional Coordination Chemistry



MURAHASHI, Tetsuro TACHIBANA, Yuki KIMURA, Seita TANIWAKE, Mayuko NOGAWA, Kyoko Professor (April, 2012–) Graduate Student\* Graduate Student\* Secretary Secretary

Our research focuses on synthesis and structural elucidation of a new class of transition metal complexes. This research leads to development of fundamental concepts of transition metal chemistry as well as applications to catalysis and materials science. Novel synthetic methods are developed to realize transition metal complexes having unique bonding nature. The newly synthesized transition metal complexes are further converted to more reactive forms, and their reaction mechanisms are elucidated. The aspects gained by this research are applied to the understanding and development of molecular catalysis. Furthermore, unique properties of low-dimensional metal-organic hybrid molecules are investigated and developed in our group.

### 1. Synthesis and Chemical Properties of Metal Chain Sandwich Complexes

The molecular sandwich framework is one of the fundamental structures in transition metal chemistry. It had been believed that the structural concept can be applied only to mono- and dinuclear complexes. Our group revealed that the multinuclear sandwich complexes containing a one-dimensional metal chain or a two-dimensional metal sheet exist stably.<sup>1,2</sup> These findings expand the scope of the structural concept of sandwich compounds (Figure 1).



Figure 1. Schematic representation of sandwich compounds: (A) mononuclear metallocenes, (B) one-dimensional metal chain sandwich complexes, (C) two-dimensional metal sheet sandwich complexes. Our group revealed the existence of categories B and C.

For the one-dimensional metal chain sandwich complexes, we successfully developed a synthetic method that enables the size-selective construction of a metal chain sandwich framework. Furthermore, our laboratory revealed that metal chain sandwich complexes show unique chemical properties such as i) dynamic sliding behavior of polyene ligands on a metal chain and ii) photo-induced flipping of polyene ligands on a metal chain.

### Redox-Switchable Metal Assembling and Ligand Coupling in Sandwich Frameworks<sup>3)</sup>

In view of the widely developed redox chemistry of metallocenes and other mononuclear sandwich compounds, it is intriguing to elucidate the redox properties of the multinuclear sandwich compounds. Here we disclosed two novel modes of redox-induced reversible structural changes: i) redox-switchable reversible splitting of a Pd<sub>4</sub> chain via translocation, and ii) redoxswitchable reversible C–C coupling of  $\pi$ -conjugated ligands in tetrapalladium sandwich complexes (Figure 2). These results provide new aspects for the redox function of (sp<sup>2</sup>-carbon)– (multinuclear metal)–(sp<sup>2</sup>-carbon) sandwich frameworks.



Figure 2. Schematic representations of (A) the reversible assembling of dimetal moieties and (B) the reversible coupling of ligands.

Oxidation of 1 or 1-OMe with  $AgBF_4$  or  $AgPF_6$  (2 equiv.) resulted in the formation of 2 or 2-OMe (bpbb = 1,2-bis(4phenyl-1,3-butadienyl)benzene). Reduction of 2 or 2-OMe with  $Fe(C_5H_5)_2$  (2 equiv.) yielded 1 or 1-OMe (eq. 1). During the two-electron oxidation, the Pd–Pd–Pd–Pd chain is cleaved, and the two Pd<sub>2</sub> units undergo translocation to the outer position during the oxidation process. In the reduction process, the two separate Pd<sub>2</sub> units migrate to the inside position to form the Pd–Pd–Pd–Pd chain. The observed redox behavior is highlighted by the coupling of two events, metal–metal bond formation/cleavage and metal translocation, under the redox control, and provides evidence of redox-switchable movement of multiple metal atoms associated with assembling/disassembling behavior in the space between sp<sup>2</sup>-carbon planes.

We also found another mode of the redox-induced structural change featured by the intramolecular reversible C–C bond formation between the two sandwich ligands, by employing 3 (eq. 2).



### 2. Synthesis and Reactivity of Metal Sheet Sandwich Complexes

For the two-dimensional metal sheet sandwich complexes, our group has shown that six-, seven-, eight-, and nine-membered carbocycles, as well as polycyclic arenes behave as the excellent binders for metal sheets. These metal sheet sandwich complexes are stable even in solution. Thus, a reactive form of such metal sheet sandwich complexes may provide a new opportunity to develop sandwich type late transition metal catalysts.

#### Selective Synthesis of Mixed Metal Pd<sub>2</sub>Pt and PdPt<sub>2</sub> Complexes of Tropylium<sup>4)</sup>

Our group has reported a series of homonuclear metal sheet sandwich complexes, which have a metal sheet of a single metal element. Here, we reported selective construction of the triangular PdPt<sub>2</sub> and Pd<sub>2</sub>Pt cores in a common sandwich framework, where a key is to use different carbocyclic ligands for the different composite sandwich; *i.e.*, cycloheptatrienyl for PdPt<sub>2</sub>, and cycloheptatriene for Pd<sub>2</sub>Pt (complexes **5-PdPt<sub>2</sub>** and **6-Pd<sub>2</sub>Pt**).



After construction of the  $Pd_2Pt$  core, the cycloheptatriene ligands can be converted to cycloheptatrienyl ligands where a (carbocyclic ligand)–(metal triangle)–(carbocyclic ligand) sandwich structure is retained (eq. 3). It was confirmed that the



mixed metal core once formed in a cycloheptatrienyl sandwich framework is robust against the intermolecular metal scrambling.

Thus, it has been proven that the mixed-metal triangular trimetal sandwich complexes are synthesizable in a selective manner. The present results expand the scope of the structural variability of multinuclear sandwich complexes from homonuclear to heteronuclear series.

### 3. Reaction Mechanism of Highly Reactive Pd–Pd Complexes

Our group has shown that a homoleptic dinuclear Pd–Pd complexes of nitriles such as  $[Pd_2(CH_3CN)_6][BF_4]_2$  are isolable.<sup>5)</sup> These homoleptic solvento-Pd<sub>2</sub> complexes are highly substitutionally labile, and allow us to investigate reaction mechanism of the Pd–Pd complexes in details. For example, our group has shown that a Pd–Pd moiety adds to various unsaturated hydrocarbons in a syn addition manner.

### Dinuclear Addition of a Pd–Pd Moiety to Arenes<sup>6)</sup>

We found that a Pd–Pd moiety undergoes addition reaction to arenes to afford the bi- $\pi$ -allyl type dipalladium complexes. Previously simple  $\mu - \eta^2 : \eta^2 - (\text{arene}) - \text{Pd}_2$  complexes have been isolated, but now we found that a Pd–Pd moiety undergoes dinuclear addition to several arenes to afford a novel binding mode of arene–Pd<sub>2</sub>, namely  $\mu$ - $\eta^3:\eta^3$ -(arene)–Pd<sup>II</sup><sub>2</sub> (Figure 3), through isolation and characterization of **8'**, **9**, and **10**. The fact that bi- $\eta^3$ -allyl type structure is accessible via dinuclear syn-addition of a Pd–Pd moiety suggest possible involvement of a new activation mode of arenes by a Pd–Pd species in some palladium-catalyzed transformations.



Figure 3. Dinuclear addition of a Pd-Pd moiety to arenes.



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### **Development of Functional Metal Complexes** for Artificial Photosynthesis

Department of Life and Coordination-Complex Molecular Science Division of Functional Coordination Chemistry



MASAOKA, Shigeyuki KONDO, Mio YOSHIDA, Masaki NAKAMURA, Go OKAMURA, Masaya MURASE, Masakazu ITOH, Takahiro KUGA, Reiko LIU, Ke KANAIKE, Mari WAKABAYASHI, Kaori TANIWAKE, Mayuko Associate Professor Assistant Professor Graduate Student\* Graduate Student Graduate Student Graduate Student<sup>†</sup> Graduate Student Technical Fellow Technical Fellow Technical Fellow Secretary

Artificial photosynthesis is a solar energy conversion technology that mimics natural photosynthesis, and poised to be one of the next big breakthroughs in energy. Our group studies chemistry of transition metal complexes for the realization of artificial photosynthesis. Efforts have focused on development of new catalysts for multi-electron transfer reactions and understanding the reaction mechanism. During the last year, we reported (i) a mononuclear ruthenium complex showing multiple proton-coupled electron transfer toward multi-electron transfer reactions,<sup>1)</sup> (ii) kinetics and DFT studies on water oxidation catalyzed by a mononuclear ruthenium complex,<sup>2)</sup> and (iii) two-step photoexcitation of a platinum complex to produce hydrogen from water.<sup>3)</sup> We also demonstrated (iv) the self-assembly of microstructures from dinuclear ruthenium complexes and their structural transformation.<sup>4)</sup> The research projects (i), (ii) and (iv) are introduced in this report.

### 1. A Mononuclear Ruthenium Complex Showing Multiple Proton-Coupled Electron Transfer toward Multi-Electron Transfer Reactions<sup>1)</sup>

Proton-coupled electron transfer (PCET) is an important chemical process that involves the concerted transfer of a proton (H<sup>+</sup>) and an electron (e<sup>-</sup>). It is widely employed to achieve multi-electron transfer reactions such as water oxidation by photosystem II and nitrogen fixation by nitrogenase as well as solar energy conversion in artificial photosynthesis, since high-energy intermediates and/or electrostatic charge buildup during the reactions are generally avoided by going through PCET processes.

In this context, we synthesized and characterized a new ruthenium(II) complex,  $[Ru(trpy)(H_2bim)(OH_2)](PF_6)_2$  (1)

(H<sub>2</sub>bim = 2,2'-biimidazole and trpy = 2,2':6',2"-terpyridine), where the H<sub>2</sub>bim and M-OH<sub>2</sub> moieties in the molecule are expected to serve as proton-dissociation sites. Electrochemical studies in aqueous solutions under various pH conditions afforded the Pourbaix diagram (potential versus pH diagram) of 1, where the  $pK_a$  values found from the diagram agree well with those determined spectrophotometrically. It was also found that 1 demonstrates four-step PCET reactions to give the four-electron oxidized species, [Ru<sup>IV</sup>(trpy)(bim)(O)]<sup>2+</sup>, without electrostatic charge buildup during the reactions (Figure 1). The multiple PCET ability of 1 would be applicable to various multi-electron oxidation reactions. Catalysis of electrochemical water oxidation was indeed evaluated in the initial attempt to demonstrate multi-electron oxidation reactions, revealing that the water oxidation potential for 1 is lower than that for 2,2'-bipyridine analogue,  $[Ru(trpy)(bpy)(OH_2)]^{2+}$  (2) (bpy = 2,2'-bipyridine), which is known as an active catalyst for water oxidation.



Figure 1. A schematic view of the four-step PCET reaction of 1 to give the four-electron oxidized species,  $[Ru^{IV}(trpy)(bim)(O)]^{2+}$ .

# 2. Kinetics and DFT Studies on the Mechanism of Water Oxidation Catalyzed by Mononuclear Ruthenium Complexes<sup>2)</sup>

Visible light-induced water splitting is one promising

approach for artificial photosynthesis. This solar-to-fuels conversion consists of the two half-cell reactions; reduction of water to H<sub>2</sub> (2H<sup>+</sup> + 2e<sup>-</sup>  $\rightarrow$  H<sub>2</sub>) and oxidation of water to O<sub>2</sub> (2H<sub>2</sub>O  $\rightarrow$  O<sub>2</sub> + 4H<sup>+</sup> + 4e<sup>-</sup>). Particularly, development of molecular catalysts for O<sub>2</sub> evolution has been considered more difficult than the hydrogen side, since the O<sub>2</sub> evolution requires removal of four protons and four electrons.

We previously provided two important aspects with regard to this issue as follows. (i) Mononuclear ruthenium catalysts can be classified into two groups based on the rate law observed. One obeys  $d[O_2]/dt = k[catalyst][Ce^{4+}]$  and the other obeys  $d[O_2]/dt = k[catalyst]^2$ , where the rate is first- and secondorder to the catalyst concentration, respectively. (ii) A radical coupling of the oxo atom from Ru catalyst and the radical-like O(hydroxo) atom from hydroxocerium(IV) species plays an important role in the O–O bond formation in the former monoruthenium-catalyzed O<sub>2</sub> evolution reactions. However, our previous studies has not provided the spectroscopic evidence of the highest-valence intermediate for the Ru species prior to the O–O bond formation.

In this work, the reaction mechanism of the Ce<sup>4+</sup>-driven water oxidation catalyzed by  $[Ru(trpy)(bpy)(OH_2)]^{2+}$  (2, Ru<sup>II</sup>-OH<sub>2</sub>). As a result, the Ru<sup>V</sup>=O species, together with other intermediates in the multi-step electron transfer processes, were spectrophotometrically followed by use of the global kinetic analysis using the singular value decomposition (SVD) method. We also demonstrate that each spectral component can be rationally reproduced by TD-DFT (time-dependent density functional theory) calculation. Moreover, it is also found that the Ru<sup>V</sup>=O species can be written as the resonance structure of the Ru<sup>IV</sup>-O<sup>•</sup> species (Ru<sup>V</sup>=O  $\leftrightarrow$  Ru<sup>IV</sup>-O<sup>•</sup>), indicating that the Ru<sup>V</sup>=O species bears a substantial radical character at the O(oxo) atom. This study suggests that a radical-radical coupling of Ru<sup>V</sup>=O and hydroxocerium(IV) species predominates the major path leading to the dioxygen formation.



Figure 2. (Left) The optimized structure of the  $Ru^V=O$  species in the doublet state. The distribution of the Mulliken atomic spin density is also overlaid. The spin density is located on the 3d orbital of the ruthenium ion (+0.40) and the 2p orbital of the O(oxo) atom (+0.61). (Right) A proposed catalytic cycle of the Ce<sup>4+</sup>–driven water oxidation catalyzed by 2.

# 3. Self-Assembly of Tubular Structures from Dinuclear Ruthenium Complexes and Their Structural Transformation<sup>4)</sup>

Controlled self-assembly of metal complexes is of high scientific and technological importance for the development of multi-functional materials and devices. Among various types of metal complexes, mixed-valence complexes have attracted much attention because of their wide range of interesting physical and chemical properties from charge-transfer interactions between metal ions linked via bridging ligands. In particular, lowdimensional assembly of such mixed-valence complexes gives rise to specific electronic, magnetic, and optical properties. Ideally, the characteristics of such systems would be tunable by controlling the spatial arrangement of the mixed-valence complexes, resulting in electric interaction among metal complexes without linkage of covalent or coordination.

In this work, we have demonstrated that the lipid-packaged mixed-valence complex displays morphological changes with aging of the solution in dichloromethane. Formation of a bilayer structure causes morphological evolution from microtapes to microtubes, giving rise to changes in absorption spectral intensities. Moreover, these morphological and spectral changes can be reversed by standing or shaking. The technique of combination of lipid molecules and *discrete* coordination compounds makes it possible to design flexible, reversible and signal responsive supramolecular coordination systems. The concept of lipid packaging could also be expanded of other useful coordination compounds, and should allow us to further develop the nanochemistry of coordination materials.



Figure 3. Schematic illustration of the assembly of discrete metal complexes for constructing nanostructures.

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#### Award

KONDO, Mio; The 5th Shiseido Female Researcher Science Grant (2012).

† carrying out graduate research on Cooperative Education Program of IMS with Nagoya Institute of Technology

<sup>\*</sup> carrying out graduate research on Cooperative Education Program of IMS with Kyushu University

### **Visiting Professors**



#### Visiting Professor SASAI, Hiroaki (from Osaka University)

#### Design and Synthesis of Novel Enantioselective Catalysts and Their Application

Synthesis of optically active complex molecules using catalytic amount of chiral compounds plays an important role in pharmaceutical industrial processes. Our group engages in the development of novel enantioselective catalyses which involve asymmetric domino reaction promoted by an acid-base type organocatalyst, oxidative coupling of 2-naphthol derivatives using dinuclear vanadium(V) catalysts, spiro

bis(isoxazoline) ligand (SPRIX) accelerated Pd catalyses, *etc.* Recently we have realized a highly enantioselective intramolecular Rauhut-Currier reactions catalyzed by an amino acid derived organocatalyst. In addition, an umpolung acetoxylation of Pd enolate derived from alkynyl cyclohexadienones was found to be promoted by Pd-SPRIX catalyst.



### Visiting Associate Professor UEMURA, Takashi (from Kyoto University)

### Polymer Chemistry in Coordination Nanospaces

One of the most outstanding challenges in polymer materials science is the fabrication of systems that allow the controlled arrangement of monomers to be polymerized to materials useful for a desired purpose. We are developing strategies to control polymerizations in nanochannels of Metal–Organic Frameworks (MOFs) composed of metal ions and organic ligands. Use of their regulated and tunable channels for a field

of polymerization can allow multi-level controls of the resulting polymer structures, such as molecular weight, stereoregularity, reaction positions, and monomer sequences. In addition, construction of nanocomposites between MOFs and polymers provides new material platforms to accomplish many optical and electronic functions.



### Visiting Associate Professor SUDO, Yuki (from Nagoya University)

#### Understanding and Controlling the Photoactive Proteins

Light is one of the most important energy sources and signals providing critical information to biological systems. Rhodopsin molecules are photochemically



reactive membrane-embedded proteins, with seven transmembrane  $\alpha$ -helices which bind the chromophore retinal (vitamin A aldehyde). A striking characteristic of these photoactive proteins is their wide range of seemingly dissimilar functions. We are investigating them by using various techniques such as biophysical, molecular biological, biochemical, genetical and spectroscopic methods. In addition, rhodopsin molecules have great potential for controlling cellular activity by light. We are also focusing on the development of novel photocontrollable tools for the life scientists.